

# NLO QCD Corrections to s-channel Single Top Quark Production and Decay at the Tevatron

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# Outline

## 1 Single top production in the SM

- Single top quark production in the SM
- Categorizing s-channel single-top processes
- Methods used in our NLO calculation

## 2 Acceptance

- Jet finding algorithm and kinematics cuts
- Acceptance of inclusive two-jet events

## 3 Kinematical distributions

- Kinematical distributions of the final objects
- Event reconstruction
- Kinematical and spin correlations
- Background to SM Higgs searching

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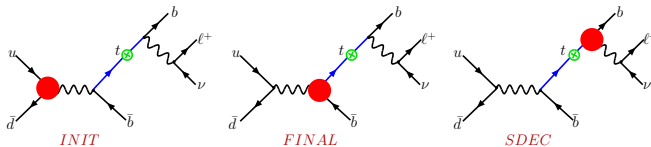
- Kinematical distributions of the final objects
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# Single top quark production in the SM

- Single top quarks are produced at hadron colliders through interactions involving a  $W$  boson and  $b$  quark. Thus,  
 $\sigma \propto |V_{tb}|^2$ .
- At tree level there are three modes:
  - s-channel  $W$  exchange  
Large rate at Tevatron Run II, small rate at LHC
  - t-channel  $W$  exchange  
Dominant rate at Tevatron Run II and LHC
  - $Wt$  associated production  
Very tiny rate at Tevatron Run II, large rate at LHC
- It is very important to understand the acceptances at NLO, including the QCD corrections to both single top quark production and decay.

# Categorizing s-channel single top quark processes

- We separate the single-top processes into smaller gauge invariant sets to organize our calculations.



● includes *soft + virtual* and *real emission corrections*.

- Keeping track on each individual contribution is useful to compare event generators with exact NLO predictions.

# Methods used in our NLO calculation

- To link the top quark production with its decay and also preserve the spin correlation between the final state particles, the **narrow width approximation** is usually adopted.
- Narrow width approximation (NWA) : fixed top quark mass
  - It does NOT work well in s-channel single-top process because the PDF luminosities drop very fast in the relevant  $x \simeq \frac{m_t}{\sqrt{s}} \sim 0.1$  range.
- The “modified” narrow width approximation (MNWA) is adopted in our calculation.
  - Generate Breit-Wigner distribution of the intermediate state top quark mass in phase space.
  - Use the generated top quark mass to calculate the matrix elements.

# Inclusive s-channel single-top cross section

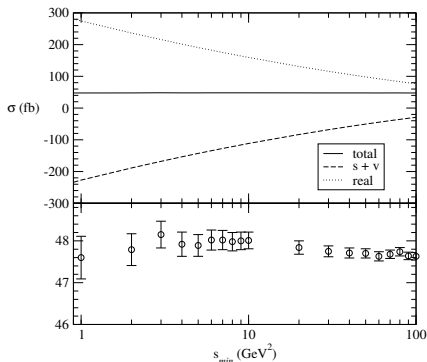
*The branching ratio of  $t \rightarrow bW^+ (\rightarrow e^+ \nu)$  is included.  
 Use CTEQ6M1 PDF*

	MNWA		NWA	
	$\sigma$ (fb)	Fraction of NLO (%)	$\sigma$ (fb)	Fraction of NLO (%)
Born Level	31.2	65.0	23.1	64.5
$O(\alpha_s)$ INIT	10.7	22.3	7.3	20.4
$O(\alpha_s)$ FINAL	5.5	11.5	5.0	14.0
$O(\alpha_s)$ SDEC	0.57	1.19	0.42	1.17
$O(\alpha_s)$ sum	16.8	35.0	12.7	35.5
NLO	47.9	100	35.8	100

The INIT contribution dominates in the  $O(\alpha_s)$  contributions.

# Methods used in our NLO calculation

- Phase space slicing method with one cutoff ( $s_{min}$ ) is used to regularize both **soft** and **collinear** singularities.



Inclusive s-channel single top quark cross section at Tevatron with  $\mu_R = \mu_F = m_t$  for  $m_t = 178$  GeV. The decay branching ratio  $t \rightarrow bW^+ (\rightarrow e^+ \nu)$  is included.

Our calculation does not depend on the theoretical cutoff  $s_{min}$ .

In our calculation,  $s_{min} = 5 \text{ GeV}^2$ .



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# Jet finding algorithm and kinematical cuts

- The cone jet algorithm is used to define the infrared safe observables.
- Smaller  $|p_Z(\nu)|$  of two-fold solutions is used to reconstruct the  $W$  boson ( $\sim 70\%$  efficiency).
- Basic kinematical cuts:

$$P_T^\ell \geq 15 \text{ GeV} \quad , \quad |\eta_\ell| \leq \eta_\ell^{\max},$$

$$\cancel{E}_T \geq 15 \text{ GeV} \quad ,$$

$$E_T^j \geq 15 \text{ GeV} \quad , \quad |\eta_j| \leq \eta_j^{\max},$$

$$\Delta R_{\ell j} \geq R_{\text{cut}} \quad , \quad \Delta R_{jj} \geq R_{\text{cut}}. \quad \Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$$

- To understand the impact of kinematical cuts on the acceptances, two sets of cuts are considered:

loose cuts:

$$\begin{aligned} \eta_\ell^{\max} &= 2.5 \\ \eta_j^{\max} &= 3.0 \\ R_{\text{cut}} &= 0.5 \end{aligned}$$

tight cuts:

$$\begin{aligned} \eta_\ell^{\max} &= 1.0 \\ \eta_j^{\max} &= 2.0 \\ R_{\text{cut}} &= 1.0 \end{aligned}$$

# Acceptances of inclusive two-jet events

$t\bar{b} + t\bar{b}j$	$\sigma$ [fb]		Acceptances (%)	
	LO	NLO	LO	NLO
(a)	22.7	32.3	73	64
(b)	19.0	21.7	61	46
(c)	14.7	21.4	47	45

(a): loose cuts  $\eta_\ell^{\max} = 2.5$ ,  $\eta_j^{\max} = 3.0$   
and  $R_{cut} = 0.5$

(b): loose cuts  $\eta_\ell^{\max} = 2.5$ ,  $\eta_j^{\max} = 3.0$   
and  $R_{cut} = 1.0$

(c): tight cuts  $\eta_\ell^{\max} = 1.0$ ,  $\eta_j^{\max} = 2.0$   
and  $R_{cut} = 0.5$

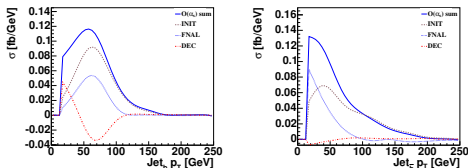
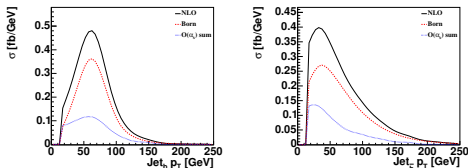
- The acceptances are sensitive to the kinematical cuts.
  - A large  $R_{cut}$  reduces the acceptance significantly because more events fail the lepton-jet separation cut.
  - With tight cuts, LO and NLO acceptances are almost the same.
  - With loose cuts, LO and NLO acceptances are quite different.
    - $\Rightarrow$  cannot use k-factor with LO kinematics
    - $\Rightarrow \Delta\sigma \sim 10\%$ ,  $\Delta|V_{tb}| \sim 5\%$
    - $\Rightarrow$  Important to have full NLO kinematics
- To maximize the acceptance, the loose cuts (a) are used in the following study.

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# Final object distributions

- Lepton and  $\cancel{E}_T$  distributions are not sensitive to NLO QCD corrections.
- $b$  and  $\bar{b}$  distributions



$b$

$\bar{b}$

NLO corrections broaden the LO distributions and shift the peak position to lower value.

$b$  and  $\bar{b}$  are sensitive to DEC and FINAL contributions, respectively.

INIT contribution dominates over FINAL and DEC.

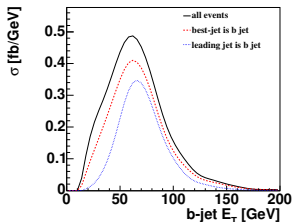
⇒ soft gluon resummation  
⇒ improve the prediction on kinematical acceptance

# Top quark reconstruction

- The two  $b$ -jets in the final state cannot be distinguished experimentally (by detectors).
- Thus, a prescription is needed to identify the correct  $b$  in order to reconstruct top quark event.
- Best-jet algorithm (to identify decay products from  $t$ )
  - Jet (or 2-jet system) which gives an invariant mass closest to 178 GeV, when combined with  $W$  boson.  
Loop over all jets, reconstruct invariant mass of  $W$ -jets system
- Leading-jet algorithm
  - highest  $p_T$  jet in the event as “ $b$ -jet”.

## b identification efficiency

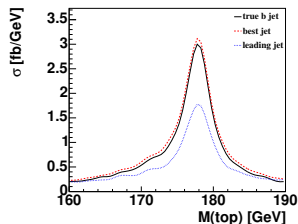
- fraction of picking up correct  $b$



best-jet algorithm: 80%

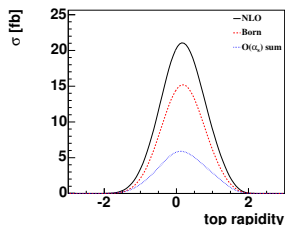
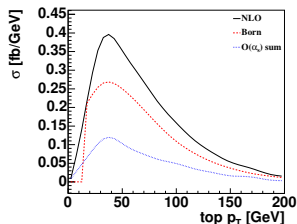
Leading-jet algorithm: 55%

- reconstructed top quark mass



The best-jet algorithm shows higher efficiency than the leading-jet algorithm.

# Top quark ( $W + \text{best-jet}$ ) distributions

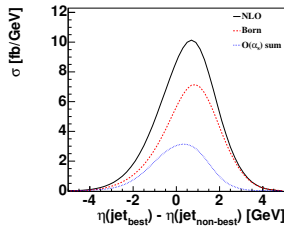
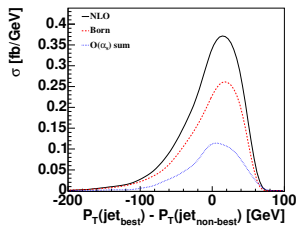


Using the best-jet results in distributions that are very similar to those obtained using the true  $b$  (or  $b + g$ ) from the top quark decay.



# Kinematical correlation between $b$ and $\bar{b}$

- At Born level,  $b$  (best-jet) and  $\bar{b}$  (non-best-jet) are highly correlated.



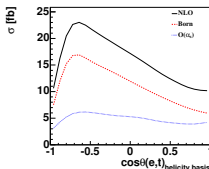
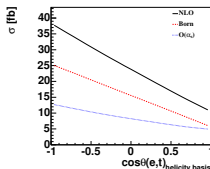
- At NLO
  - Transverse momentum difference is not largely affected.
  - The  $O(\alpha_s)$  corrections have a larger effect on the pseudo-rapidity difference.
  - The additional gluon radiated from FINAL and DEC corrections tends to weaken the correlation between  $b$  and  $\bar{b}$ .

# Top quark polarization

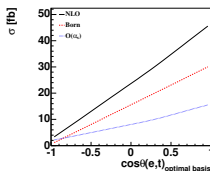
- In SM, the top quark produced in single-top event is highly polarized.
- Two bases for polarization:
  - Helicity basis  
Top quark spin is measured along the top quark direction of motion in the c.m. frame of (reconstructed top + non-best-jet).
  - Optimal basis (maximal spin correlation)  
Taking advantage of the fact that top quark is highly polarized along the direction of  $d$ -type quark (predominantly, from anti-proton).

# Top quark polarization

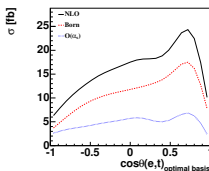
## ● Helicity basis



## ● Optimal basis



without cuts  
parton level



with cuts and  
reconstruction

$$\cos \theta_{hel} = \frac{\vec{p}_t \cdot \vec{p}_\ell^*}{|\vec{p}_t| \cdot |\vec{p}_\ell^*|}$$

$\vec{p}_t$ : three-momentum defined in c.m. frame of the two incoming partons

$\vec{p}_\ell^*$ : charged lepton three-momentum defined in rest frame of top quark.

$$\cos \theta_{opt} = \frac{\vec{p}_p^* \cdot \vec{p}_\ell^*}{|\vec{p}_p^*| \cdot |\vec{p}_\ell^*|}$$

$\vec{p}_p^*$ : anti-proton three-momentum in the rest frame of top quark

$\vec{p}_\ell^*$ : charged lepton three-momentum in the top quark rest frame.

# Top quark polarization

- To quantify the spin correlation, we define the degree of polarization  $\mathcal{D}$  and the fraction of polarization  $\mathcal{F}$ ,

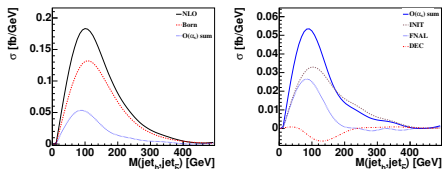
$$\mathcal{D} = \frac{N_- - N_+}{N_- + N_+}, \quad \mathcal{F}_{\mp} = \frac{1 \pm \mathcal{D}}{2}.$$

		$\mathcal{D}$		$\mathcal{F}$	
		LO	NLO	LO	NLO
Helicity	Parton level	0.63	0.58	0.82	0.79
	Reconstructed events	0.46	0.37	0.73	0.68
Optimal	Parton level	-0.96	-0.92	0.98	0.96
	Reconstructed events	-0.48	-0.42	0.74	0.71

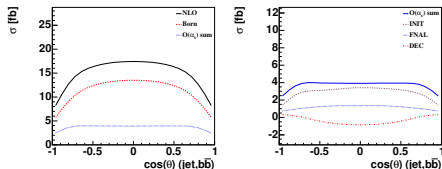
- The strong spin correlation at parton level is smeared out after event reconstruction due to the imperfect  $W$ -reconstruction.
- With  $t$ -reconstruction, both bases give about the same fraction of polarization.

# Intrinsic background to SM Higgs searching

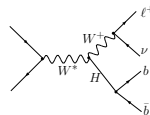
- Invariant mass of ( $b$  jet and  $\bar{b}$  jet)



- Angular distribution of  $\cos\theta_b$  in  $b\bar{b}$



- SM Higgs searching channel



$$115 < m_H < 130 \text{ GeV}$$

- The DEC contribution flattens out the  $O(\alpha_s)$  corrections compared to Born level  
 $\Rightarrow$  closer to Higgs signal distribution  
 $\Rightarrow$  This becomes more important as shape is concerned.

# Summary

- In the s-channel single-top process, the initial state contribution dominates the acceptances in all the single particle inclusive distributions.
- With loose kinematical cuts to maximize the acceptances, the full NLO kinematics needs to be studied. (A constant K-factor with LO kinematics won't work.)
- The best-jet algorithm is better to reconstruct top quark event in the s-channel single-top process.
- Higher order corrections change the kinematical and spin correlations largely. At NLO, the optimal basis give about the same fraction of polarization as the helicity basis.
- NLO QCD corrections to s-channel single-top process is important to SM Higgs boson search via  $WH$  production. In particular, DEC contribution has to be included to make more reliable background prediction.

# Summary

- Ntuple files of this NLO calculation will be made available for those who are interested.
- A similar study at LHC is in progress.
- The phenomenological study of the exclusive three-jet event will be discussed in Reinhard's talk.