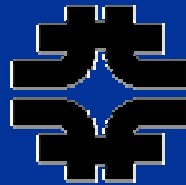


Z-primes at the Tevatron & LHC

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Outline

- Introduction: Why Z's ?
- What We Know Today
 - Theoretical Framework
 - LEP II bounds
- What We Hope To Find Tomorrow
 - Model Independent Studies
 - Bridges between Modelers & Experiment
- Outlook

Introduction: Why Z'?

- The SM is based on local gauge invariance under: $SU(3)_C \times SU(2)_W \times U(1)_Y$.
- It is natural to ask: are these all of the fundamental interactions of nature? What are the limits on more?
 - How well has LEP already explored them?
 - How do hadron machines increase our knowledge?
- Good theoretical reasons point to extra gauge bosons.
 - **GUTs** with “large” gauge groups ($SO(10), E_6, \dots$)
 - **Extra dimensions** with gauge fields in the bulk have copies of γ , Z , which are massive, neutral vector particles.
 - Theories like **Topcolor** use them to drive **EWSB**.
 - The **little Higgs** theories use massive vector particles to cancel quadratic divergences in the Higgs mass induced by W and Z .
 - New **SUSY** theories use them to survive the LEP II bound on m_h .

Theory Framework

- We extend the SM: $SU(3)_C \times SU(2)_W \times U(1)_Y \times U(1)_Z$.
 - There are Z 's coming from larger groups like $SU(2)$. We can usually describe their properties with a $U(1)$, though we don't capture all of the physics.
- Each SM fermion ψ (and the usual Higgs) is assigned a charge z_ψ under $U(1)_Z$. The Z - ψ - ψ coupling is: $g_Z z_\psi$.
- A new scalar Φ is charged under $U(1)_Z$, ($z_\Phi = 1$). Its VEV u determines the Z ' mass: $M_{Z'} = g_Z z_\psi u$.
 - Φ is neutral under the SM.
- Theory can further relate the z_ψ to each other:
 - $SU(2)_W$ invariance: $z_{uL} = z_{dL}$. Yukawas from higher dimensional operators?
 - Yukawa interactions: $z_{uL} + z_{uR} + z_H = 0$.
 - Anomaly cancellation. New Fermions?
 - FCNCs and family universality: $z_{dL} = z_{sL} = z_{bL}$.

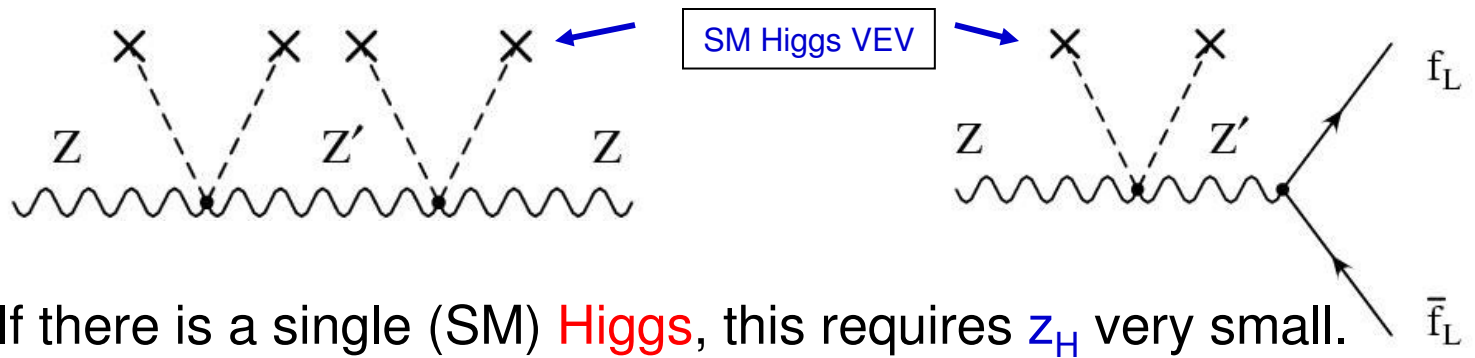
Specific “Model-lines”

	$B-xL$	$Q+xu$	$10+x\bar{5}$	$d-xu$
$q_L=(u_L,d_L)$	$+1/3$	$+1/3$	$+1/3$	0
u_R	$+1/3$	$+x/3$	$-1/3$	$-x/3$
d_R	$+1/3$	$(2-x)/3$	$-x/3$	$+1/3$
$l_L=(e_L,v_L)$	$-x$	-1	$+x/3$	$(x-1)/3$
e_R	$-x$	$-(2+x)/3$	$-1/3$	$+x/3$

- Subject to theoretical motivations, we find a set of 4 “model lines”.
- Each has a free parameter, x , specifying all of the couplings of the Z' .
- The E_6 Z' s are certain values of x for certain model lines.

e^+e^- Limits: Z-Z' Mixing

- e^+e^- colliders constrain Z' 's in two important ways.
- Precision measurements of Z physics
 - SM predictions in agreement at roughly few per mil level.
 - A Z' which mixed with Z would distort these predictions.

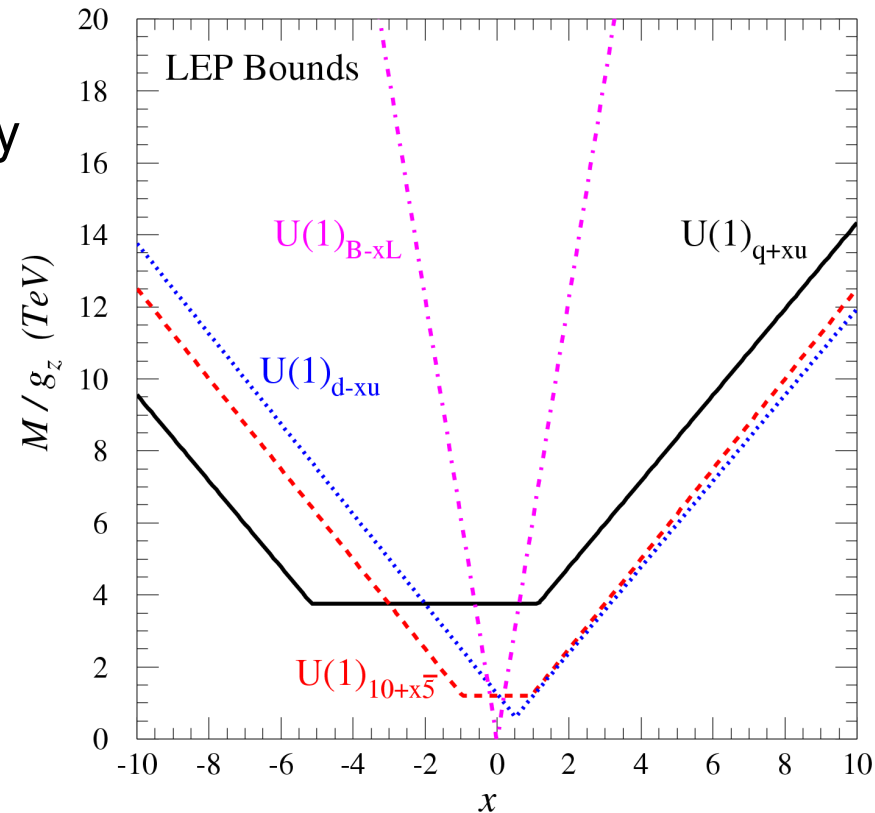
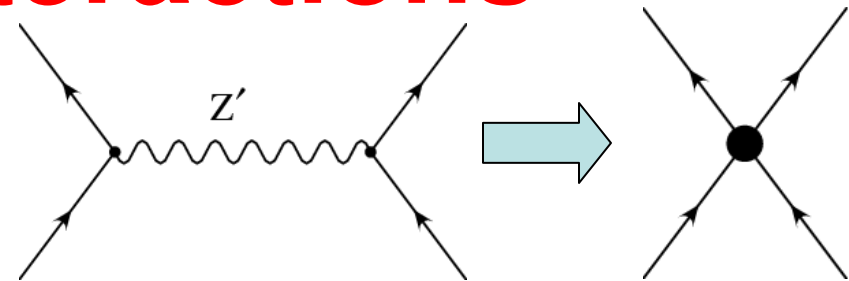


- If there is a single (SM) **Higgs**, this requires z_H very small.
- For **multiple Higgses**, one can cancel the mixing, by fine-tuning the Higgs potential against the z_H s (to the sub-% level).
- This is how the famous E_6 models survive LEP-II limits.
- I will fix $z_H = 0$.

J Hewett, T Rizzo, **Phys.Rept.183, 193 (1989)**

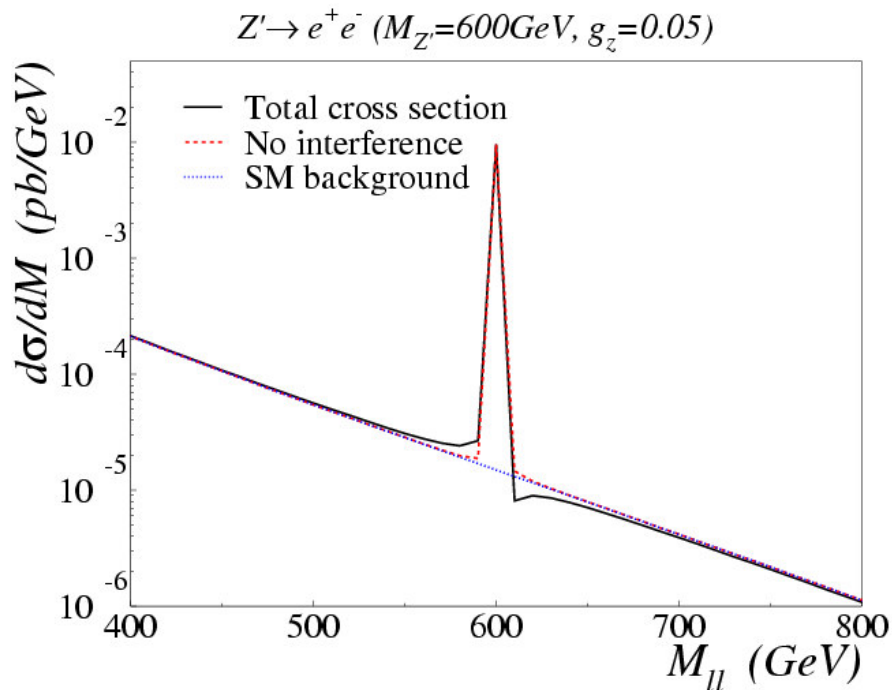
Contact Interactions

- Searching for evidence of physics beyond the SM in $e^+e^-\psi\bar{\psi}$ contact interactions.
 - At $E \ll m_{Z'}$, the Z' effects look like four fermion interactions:
 - These depend quite sensitively on how Z' couples to e , and also to other fermions.
- LEP-II has searched for evidence of contact interactions in almost all conceivable channels.
- Bounds vary from channel to channel, complicating the analysis.

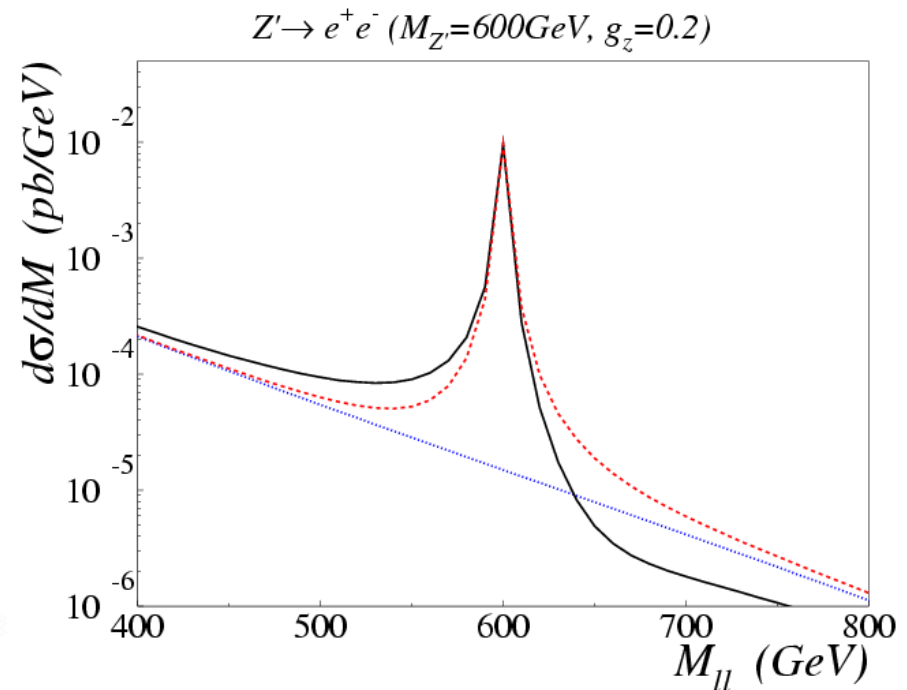


Hadron Colliders

- Hadron colliders look for Z' 's most effectively in the decay into charged leptons (e^\pm, μ^\pm, τ^\pm).
- Unlike LEP, hadron colliders are sensitive to Z' 's which do not couple to e^\pm : unexplored territory!
- The signal appears as a resonance above the (smooth) lepton pair backgrounds. Interference effects are tiny.



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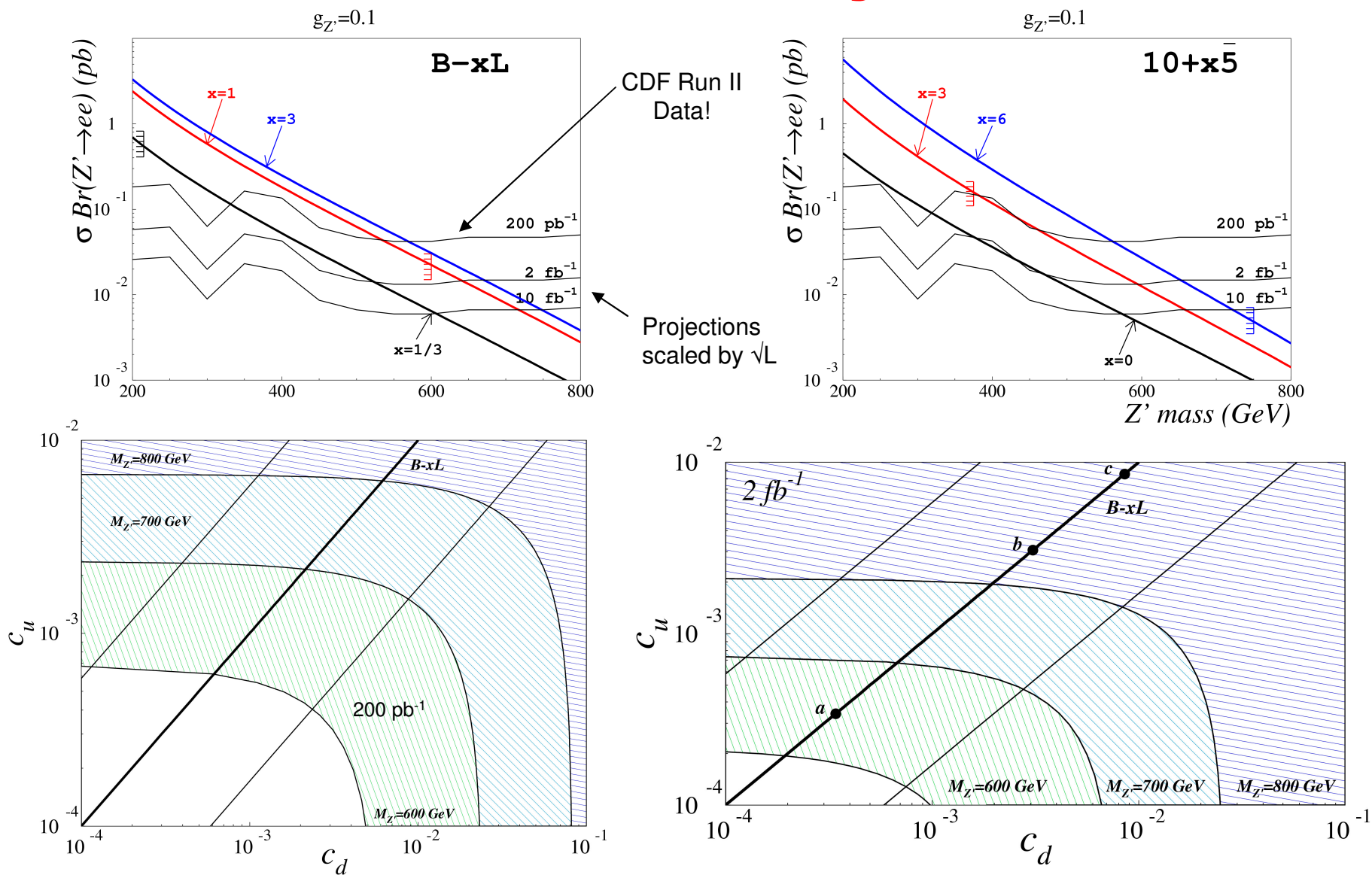
Common Ground

- Run I and existing Run II analyses have presented limits on $\sigma \times BR$. Specific E_6 models are usually overlaid to illustrate the results.
 - This is good because $\sigma \times BR$ is fairly model-independent: results can in principle be applied to other theories with Z' 's.
 - However, it is inconvenient because computing a hadronic cross section is not a completely trivial task : one has to know about QCD, PDFs, etc..
 - These are important for σ , but don't have anything to do with Z' 's.
- An alternate approach tries to factor out the universal QCD dependence, and bound only the Z' quantities themselves.
 - This works perfectly at LO and NLO in QCD; at NNLO there is some (very) small model-dependence.
 - Dominant NLO EW corrections are model-independent (and probably important at the few % level).

$$\sigma(p\bar{p} \rightarrow Z'X \rightarrow \ell^+\ell^-X) = \frac{\pi}{48s} \left[\overset{\text{QCD stuff}}{\color{red}c_u} \overset{\text{EW stuff}}{\color{blue}w_u} + \overset{\text{QCD stuff}}{\color{red}c_d} \overset{\text{EW stuff}}{\color{blue}w_d} \right]$$

$$c_u = g_Z^2 \left(z_q^2 + z_u^2 \right) \text{BR}(Z' \rightarrow \ell^+\ell^-) \quad c_d = g_Z^2 \left(z_q^2 + z_d^2 \right) \text{BR}(Z' \rightarrow \ell^+\ell^-)$$

Bounds & Projections



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TeV4LHC ?

- LHC will eventually reach into the TeV range for Z' masses.
- Equivalently, for lower masses the larger luminosity should eventually allow one to probe **smaller couplings**.
- Of course, our theoretical results for a “common ground” between experiments and model-builders are also of value at the LHC.
- Once the Z' is discovered, we will need more observables to pin down its properties.
- Tevatron inputs for high x PDFs will be important for Z' and all high mass search processes.
- Especially high x gluon PDFs are constrained by jet data, and are particularly important at LHC (LHC k-factors can be $\sim 100\%$).
- Techniques such as τ ID are being developed and can be refined using real data at Tevatron.

Outlook

- Models with massive neutral bosons, Z' 's, are of great interest for a variety of reasons. Many different kinds of models of physics beyond the SM contain Z' 's.
- e^+e^- colliders put strong bounds, and favor a Z' which does not mix with the Z . This implies either $z_H=0$ or some kind of tuning.
- We have introduced several new “model lines” which contain popular E_6 inspired models, but also generalizations of them.
- A new proposal is that we directly bound quantities which contain Z' physics, but none of the QCD inputs, bringing experiment and modelers together easily (for both sides!).

$$c_u = g_Z^2 \left(z_q^2 + z_u^2 \right) \text{BR}(Z' \rightarrow \ell^+ \ell^-) \quad c_d = g_Z^2 \left(z_q^2 + z_d^2 \right) \text{BR}(Z' \rightarrow \ell^+ \ell^-)$$

- Tevatron can provide important inputs to LHC, including high x PDFs and new experimental technique.
- Let's find nature's next force!