

Phenomenology of the chargino and neutralino systems

LCWS 2000

Fermilab

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Question: How to determine fundamental
SUSY parameters in a model
independent way in the chargino
and neutralino systems.

Answer based on:

- Choi, Djouadi, Dreiner, Z.K., Song, Zerwas
EPJC 7, C8, C14, PLB 478
- Moortgat-Pick, Fraas, BarA, Mayerotto
EPJC 7, C9, PRD 59
- Knuu, Moulaka
PRD 59, D61
- Choi, ZK, Moortgat-Pick, Zerwas, in prep.

(earlier work BarA et al '86, '89
Tanaka et al. '95
Feng et al. '95)

Analysis in unconstrained MSSM including CP phases
with R_p -conserving, at e^+e^- collider.

See also Diaz, Blödingen, Plehn

1. The goal

▶ in MSSM many new fundamental parameters

• gauge $\sim g_{\tilde{q}\tilde{q}\tilde{q}}, g_{\tilde{q}\tilde{q}q}, \dots$

• ~~SUSY~~ $\sim M_i, A_i, m_{\tilde{Q}}, m_{\tilde{L}}, \dots, \mu$

▶ after SUSY discovered, measure parameters in a model-independent way

• verify SUSY relations

• check for any relations among them

$\Rightarrow$ learn about SUSY mechanism

▶ two-step procedure



▶ need strategy

-inos

}	charginos neutralinos sleptons $\vdots$
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$M_2, \mu, \tan\beta$
 $\quad \quad \quad + M_1$
 $\quad \quad \quad + \dots$

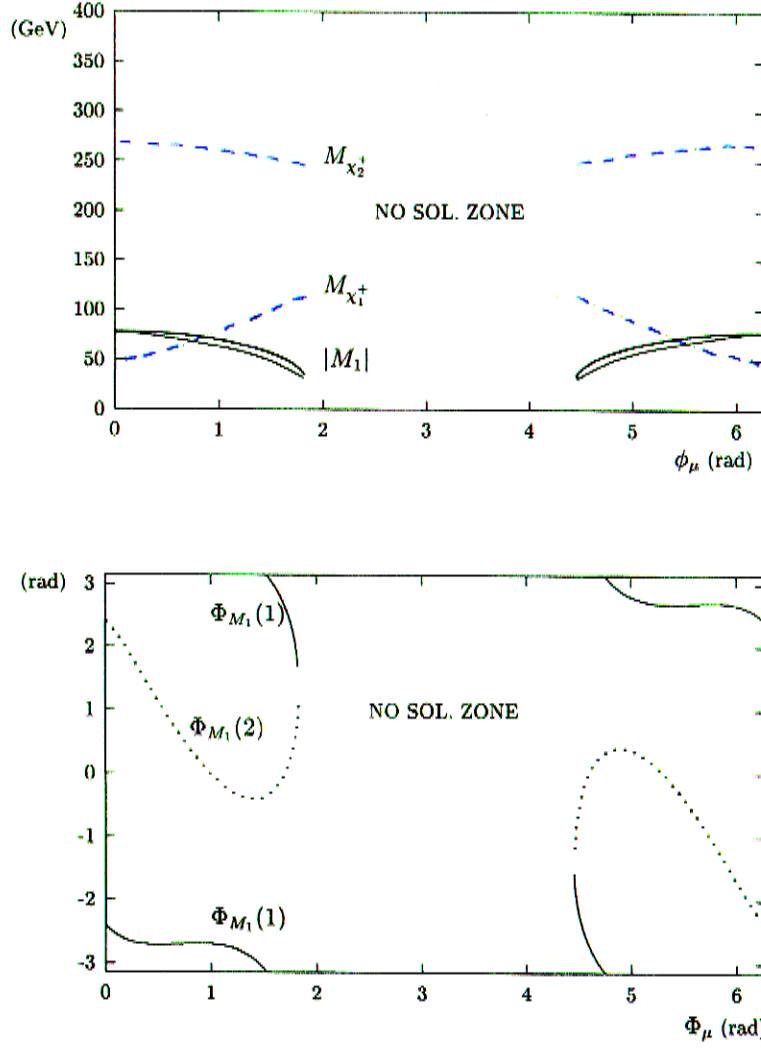
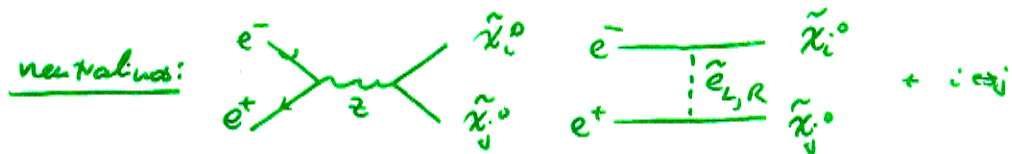
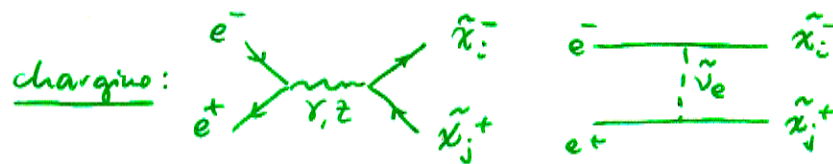


Figure 1: the twofold $|M_1|$ (top figure) and Φ_{M_1} (bottom figure) reconstruction with input choice $|\mu| = 100$ GeV, $M_2 = 120$ GeV, $M_{N_1} = 40$, $M_{N_2} = 80$; $\tan \beta = 2$. Also shown are the corresponding chargino mass values.

3. Two pair-production processes



After Fierzing t - and u -channel exchanges

$$A^{ij} = \frac{e^2}{s} Q_{\alpha\beta}^{ij} [\bar{v}(e^+) \gamma^\alpha P_\alpha u(e^-)] \cdot [\bar{u}(\tilde{\chi}_i^-) \gamma_\alpha P_\beta v(\tilde{\chi}_j^+)]$$

$\alpha, \beta = L, R$

where, for example:

charged $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ $Q_{LL}^{11} = D_L - F_L \cos 2\phi_L$

for neutrals $\tilde{\chi}_i^0 \tilde{\chi}_j^0$ $Q_{LL}^{ij} = \alpha Z_{ij} - \beta g_{ij}^L$

$$Z_{ij} = \frac{1}{2} (N_{i3} N_{j3}^* - N_{i4} N_{j4}^*)$$

$$g_{ij}^L = \frac{1}{4s_W^2 c_W^2} (N_{i2} s_W + N_{i1} c_W) (N_{j2}^* c_W + N_{j1}^* s_W)$$

Note:

(a) for charged: $Q_{\alpha\beta}$ linear in $\cos 2\phi_L, \cos 2\phi_R$

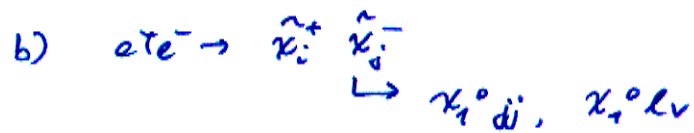
$\tilde{\nu}_e$ only in the LR amplitude

(b) for neutrals: more complicated dep. on mixing

$\tilde{e}_{L,R}$ in all amplitudes.

Measuring chargino masses and mixing angles

masses : a) from energy scan



di-jet energy or mass distribution

error $\sim 0.05 - 0.2$ GeV (Blair, Martin)
 also $m_{\tilde{\chi}_1^0}$ determined. '99

mixing angles

a) if only $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ produced

- exploit $\tilde{\chi}_1^+$ - polarization or
- use transverse beam polarization

b) if all $\tilde{\chi}_i^+ \tilde{\chi}_j^-$

σ_L, σ_R sufficient

$\Rightarrow$ unique solution for $\cos 2\phi_L, \cos 2\phi_R$

note: $\tilde{\nu}_e$ contribution essential

here we assume $m_{\tilde{\nu}_e}$ known from
 other processes Peskin
 or $\tilde{\chi}_i^+ \tilde{\chi}_j^- \rightarrow \tilde{\chi}_1^0 e \tilde{\nu}_e$
 Rantztog-Pich, Franz

m_{SUGRA}	RR1	RR2
$\tan\beta$	3	30
m_0	100	160
$M_{1/2}$	200	200
$m_{\tilde{\chi}_1^\pm}$	128	132
$m_{\tilde{\chi}_2^\pm}$	346	295
$m_{\tilde{\nu}}$	166	206
$m_{\tilde{\chi}_1^0}$	70	72

Choi et al.

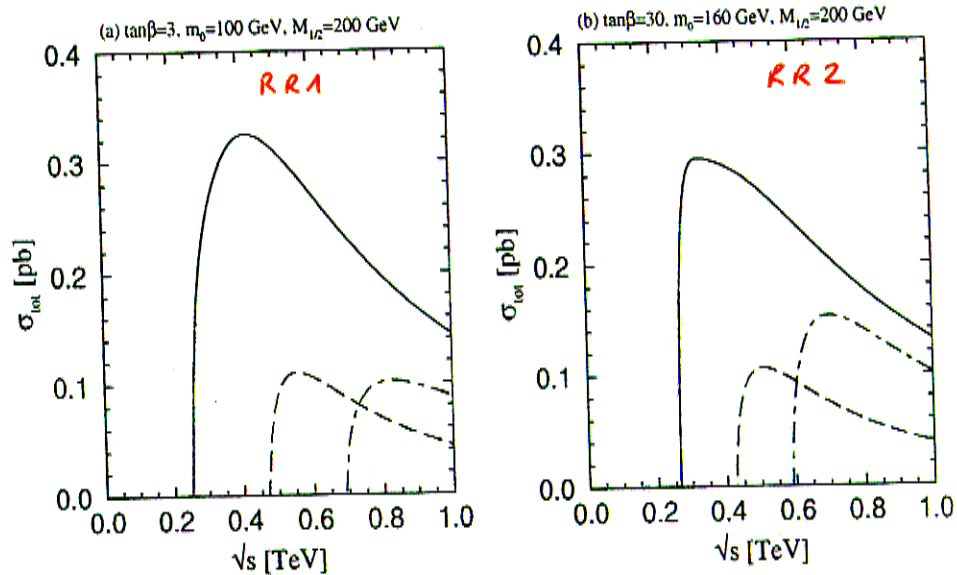


Figure 2: The cross sections for the production of charginos as a function of the c.m. energy (a) with the set $[\tan\beta = 3, m_0 = 100 \text{ GeV}, M_{1/2} = 200 \text{ GeV}]$ and (b) with the set $[\tan\beta = 30, m_0 = 160 \text{ GeV}, M_{1/2} = 200 \text{ GeV}]$: solid line for $\tilde{\chi}_1^- \tilde{\chi}_1^+$ production, dashed line for $\tilde{\chi}_1^- \tilde{\chi}_2^+$ production, and dot-dashed line for $\tilde{\chi}_2^- \tilde{\chi}_2^+$ production.

M_2	152	150
μ	316	263

Chargino Production

$$e^+e^- \rightarrow \chi_1^+ \chi_1^-$$

H.D. Marquardt
G. Blair

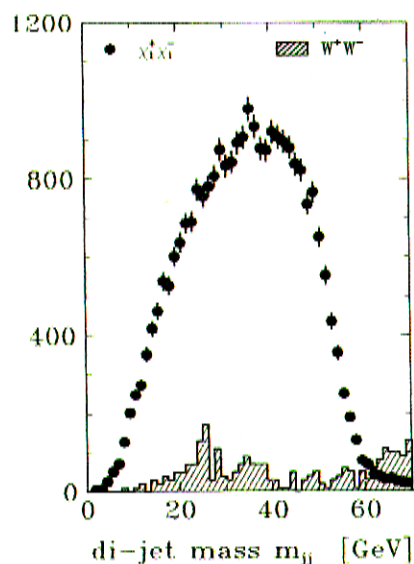
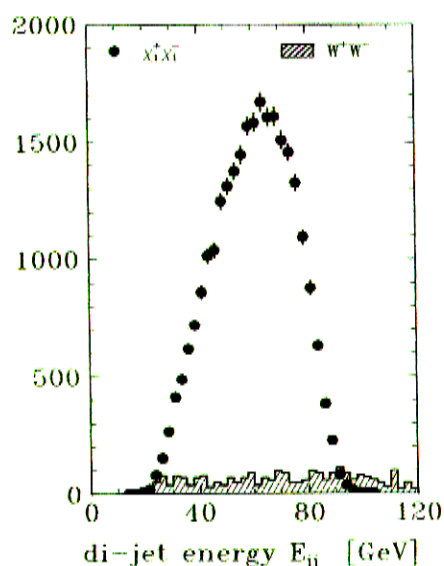
$$\begin{aligned} e^+e^- &\rightarrow \chi_1^+ \chi_1^- \\ &\rightarrow l^\pm \nu \chi_1^0 q \bar{q}' \chi_1^0 \quad Br = 2 \cdot 0.45 \cdot 0.55 \end{aligned}$$

background:

$$W^+W^- < 10\%$$

$\Rightarrow m_{\chi_1^\pm}$ and $m_{\chi_1^0}$ from di-jet energy and mass spectra

$\sqrt{s} = 320 \text{ GeV}$



$$\begin{aligned} m_{\chi_1^\pm} &= 127.7 \pm 0.2 \text{ GeV} & \Delta m(\chi_1^\pm - \chi_1^0) &= 55.8 \pm 0.15 \text{ GeV} \\ \sigma(e_L^- e_R^+) \mathcal{B} &= 330 \pm 1.5 \text{ fb} \end{aligned}$$

Scan at threshold, $10 \text{ fb}^{-1}/\text{point} \Rightarrow \Delta m = 0.04 \text{ GeV}$

- [7] G. Moortgat-Pick, H. Fraas, A. Bartl, W. Majerotto, Eur. Phys. J. C 7 (1999) 113.
- [8] H.E. Haber, G.L. Kane, Phys. Rep. 117 (1985) 75.
- [9] A. Bartl, H. Fraas, W. Majerotto, Z. Phys. C 30, (1986) 441.
- [10] S. Ambrosanio, G.A. Blair, P. Zerwas, ECFA-DESY LC-Workshop, 1998.
- [11] S.P. Martin, P. Ramond, Phys. Rev. D 48 (1993) 5365.
- [12] A. Bartl, H. Fraas, W. Majerotto, B. Mösslacher, Z. Phys. C 55 (1992) 257.

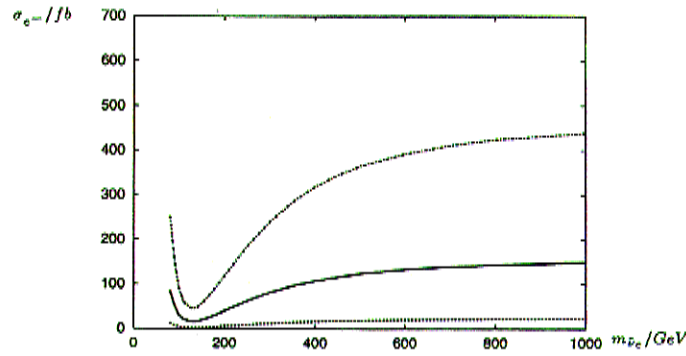


Fig. 1: Cross section σ_{e^-} (see eq. (1)) for $\sqrt{s} = 270$ GeV in scenario A for unpolarized beams (solid line), with only e^- beam polarized $P_- = +85\%$ (dotted line) and with both beams polarized $P_- = -85\%$, $P_+ = +60\%$ (dash-dotted line).

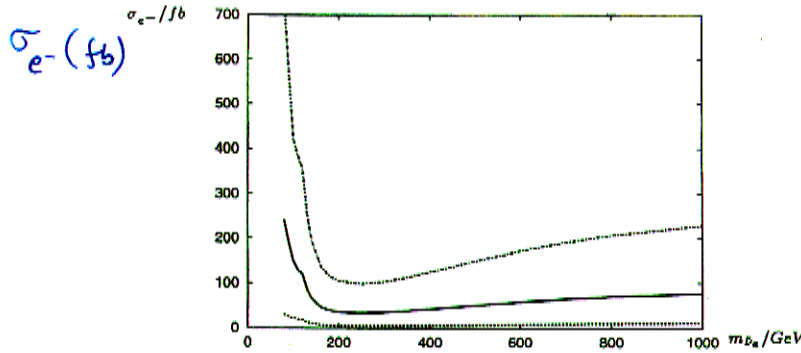


Fig. 2: Cross section σ_{e^-} (see eq. (1)) for $\sqrt{s} = 500$ GeV in scenario A for unpolarized beams (solid line), with only e^- beam polarized $P_- = +85\%$ (dotted line) and with both beams polarized $P_- = -85\%$, $P_+ = +60\%$ (dash-dotted line).

Mixing angles in the neutrino sector:

in general 6 angles + 10 phases

for $|M_1|, |M_2| \gg |\mu|$ or $\ll |\mu|$

approximate analytical expressions can be found
CKM2, in prep

$$N^* M_N N = M_{\text{diag}}$$

$$N^* = \begin{pmatrix} e^{i\alpha_1} & & & \\ & e^{i\alpha_2} & & \\ & & e^{i\alpha_3} & \\ & & & e^{i\alpha_4} \end{pmatrix} \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & c_{34} & s_{34} \\ & & -s_{34}^* & c_{34} \end{pmatrix} \begin{pmatrix} c_{13}c_{14} & s_{12} & s_{13} & s_{14} \\ -\tilde{s}_{12} & c_{23}c_{24} & s_{23} & s_{24} \\ -s_{13}^* & -s_{24}^* & -s_{12}s_{14}^* - s_{23}s_{24}^* & c_{14}c_{24} \end{pmatrix}$$

$$c_{ij} = \cos \theta_{ij}, \quad s_{ij} = \sin \theta_{ij} e^{i\delta_{ij}}$$

$$\text{e.g. } s_{13} = - \frac{m_2 s_w}{|M_1|^2 - |\mu|^2} \left(M_1 e^{i\phi_1} c_\rho + \mu e^{-i\phi_1} s_\rho \right)$$

$$\vdots$$

very complicated

However, use unitarity of diagonalization matrices to check completeness of

$$(\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm) \quad \text{and} \quad (\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0)$$

systems

⇒ first check the completeness of $(\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm)$

* sum rules from 2×2 mixing

for couplings: $\sum_{i,j=1,2} |Q_{\alpha\beta}^{ij}|^2 = 2(|D_\alpha|^2 + |F_\alpha|^2)$

$\alpha\beta = LL, RL, RR$

does not depend on ANY ~~soft~~ param!

for cross sections: (for $\sqrt{s} \gg m_\chi, m_0$)

$$\sum_{i,j=1,2} (\sigma_L^{ij} + \sigma_R^{ij}) = \frac{347 \text{ fb}}{S/(\text{TeV})^2} (1 + O(\frac{m^2}{s}))$$

→ F

at finite energy: $\sigma_L^{ij}, \sigma_R^{ij}$ quadratic

function of $x \equiv \cos 2\phi_L$

$y \equiv \cos 2\phi_R$

⇒ nontrivial consistency conditions for
measured cross sections

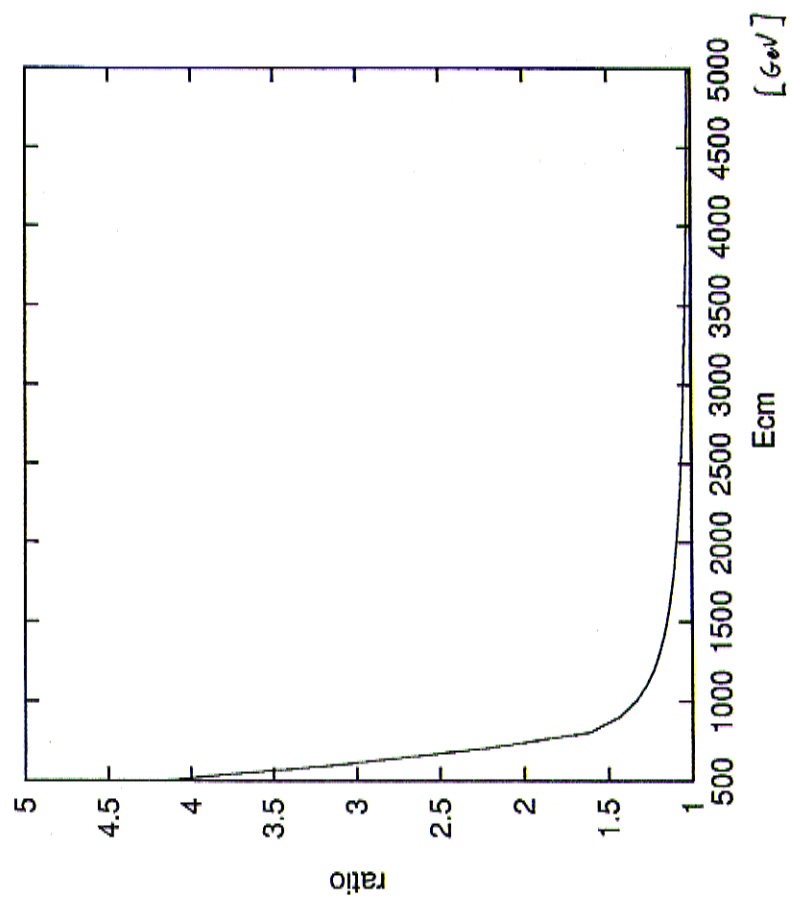
* for neutralinos 4×4 mixing

$$\sum_{i,j=1,\dots,4} (\sigma_L^{ij} + \sigma_R^{ij}) = \frac{325 \text{ fb}}{S/(\text{TeV})^2} (1 + O(\frac{m^4}{s}))$$

→ F

Charginos

$\tan(\beta)=3$



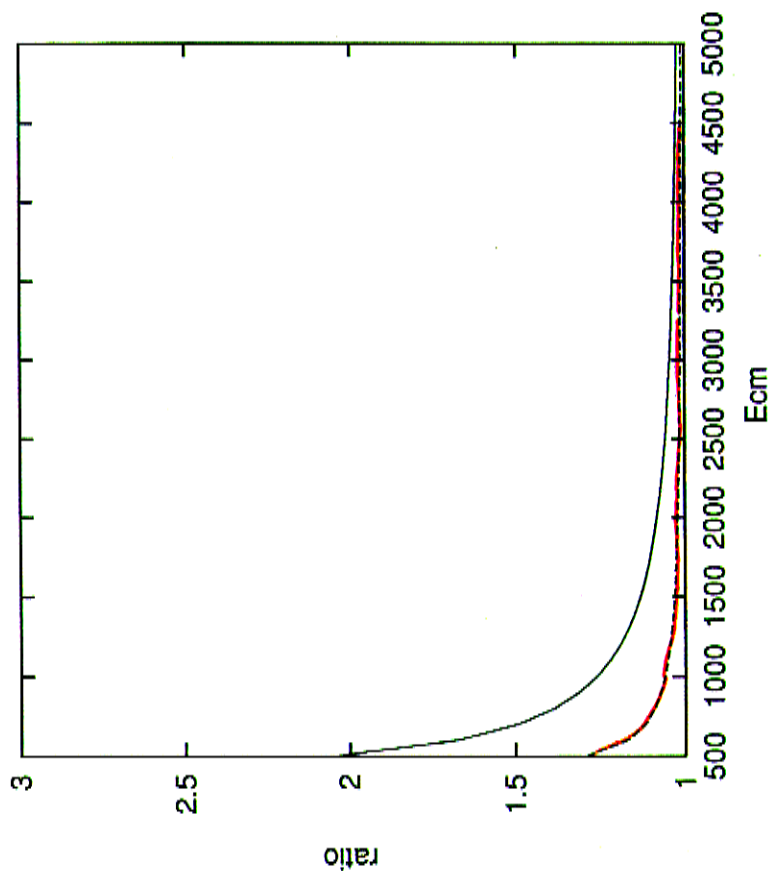
$\frac{\text{asymptotic}}{\text{exact}}$

CKM2

$$\sum_{ij} \sigma_{ij}^{as} = \frac{\pi \alpha^2}{g_s^2} \left[\frac{8}{3} \left(\frac{s_W^2 - \frac{1}{2}}{s_W^2 c_W^2} \right)^2 + \frac{8}{3} \frac{1}{c_W^4} + \frac{16}{c_W^4} + \frac{1}{s_W^4 c_W^4} \right]$$

neutrinos

$\tan(\beta)=3$



Conclusions

- charginos + neutralinos = overconstrained system
- $\sigma_{L,R}$ + masses $\Rightarrow$ couplings $\Rightarrow$ complete determination
- analytic algorithms $\Rightarrow M_2, M_1 e^{i\phi_1}, \mu e^{i\phi_\mu}, \tan\beta$
($\rightarrow$ Plehn)
- expected errors under control
except for large $\tan\beta$
- sum rules for completeness checks
(consistency)
- radiative corrections
 - other parameters enter
 - global analyses needed
($\rightarrow$ Diaz
also T. Blank + W. Hollik)
- other LC modes
 $e\bar{e}, \gamma\gamma \rightarrow$ Blöchlinger