

SUSY and susy breaking Scale

LCWS 00'

Oct. 26th

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** Introduction

** Model predictions for sparticle masses and their relation to susy breaking scale/parameters.

** Capabilities of \int an e^+e^- LC ($\sqrt{s} \gtrsim 500$ GeV) for precision measurements of sparticle properties

*** Lessons we can learn about **the scale** and **mechanism** of susy breaking from these measurements.

** How will LC **crucially complement/extend** the achievements of LHC

** What would be **desired** extensions in **colliding ptcles./ luminosity / energy** from the point of view of susy investigations?

1] A. Accamando et al, Phys. Rep. 289 (1998) 1 ; NEW TESLA TDR

2] Case for 500 GeV e^+e^- linear collider, hep-ph/0009022

3] Web page <http://actahep.keke.jp/>

My charge:

To summarize how

1) LC will aid in establishing SUSY as a viable theory

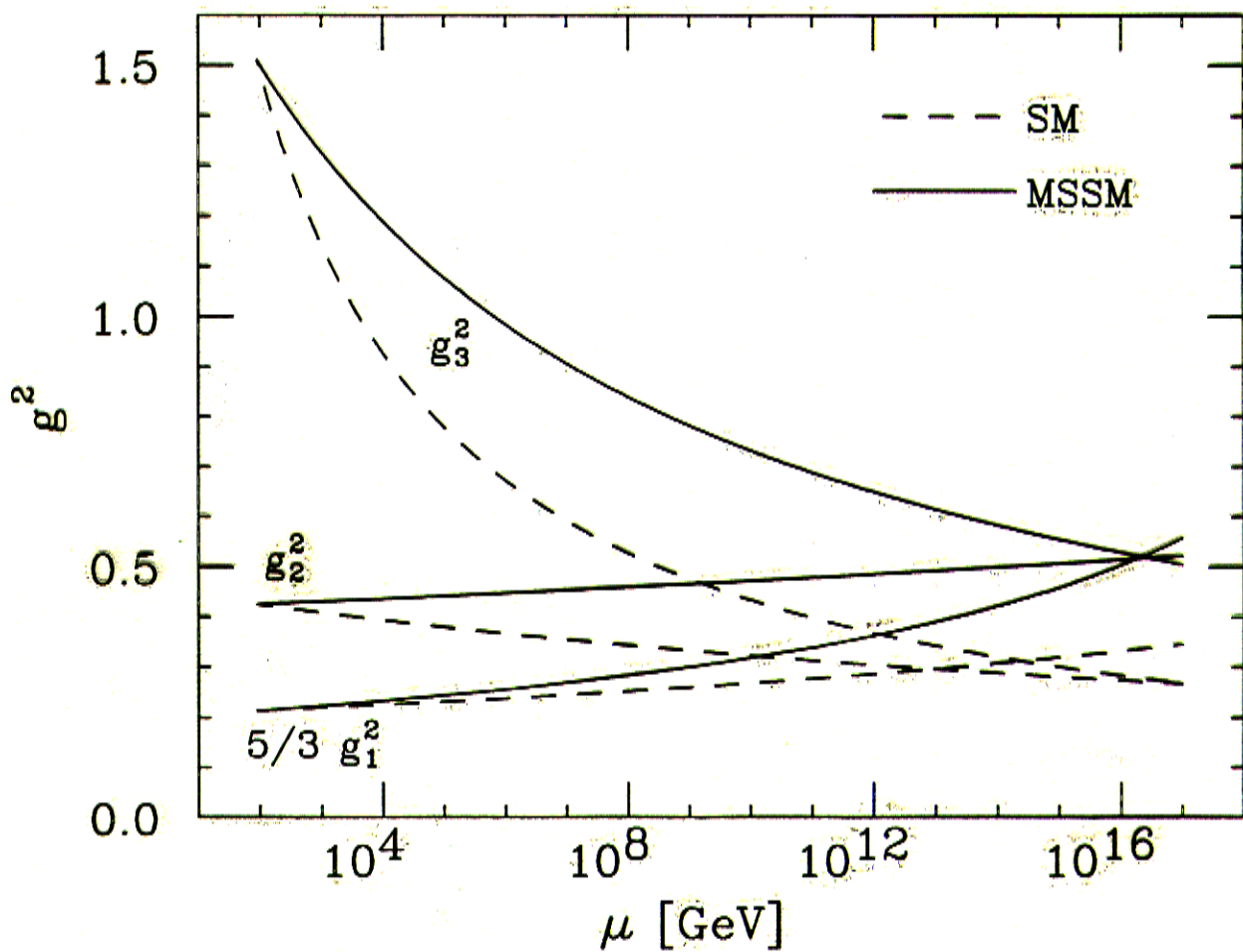
2) LC will help glean information about the scale and the mechanism of SUSY breaking.

Introduction :

Supersymmetry (SUSY) [TeV scale]

- * stabilizes m_h at electroweak scale. The only concrete and worked out solution we have. ***
- * Provides a natural mechanism for the SSB of EW symmetry.
- * Existence of Higgs scalar no longer very special

*** Unification of Couplings.

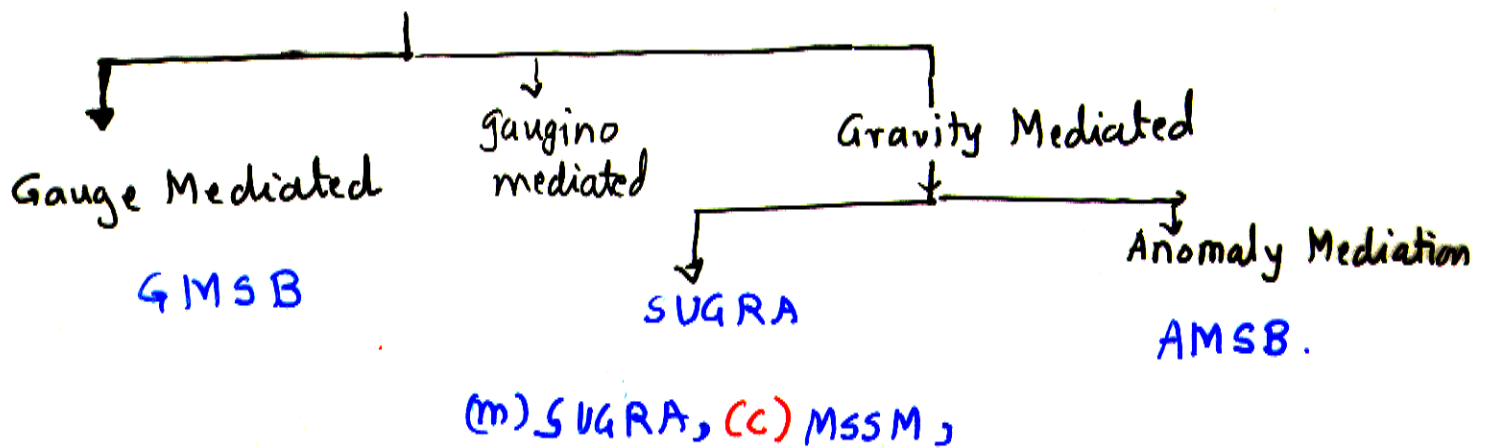


BUT

*** Of course \exists extra dim. then perhaps problem does not exist.

- ** supersymmetry is broken in real life
- ** susy breaking not really understood completely
- ** susy breaking happens in secluded sector.

** Mediation of this breaking to low energy world



** soft susy breaking parameters at high scale decided by the susy breaking mechanism and its mediation. $10^4 \text{ GeV} < M_{\text{SB}} < M_{\text{Pl}}$

** sparticle masses at the EW scale (for all sparticles) and couplings (for some of them) depend on this mechanism.

This is the clue provided to us to point to high scale physics

	$M_{\tilde{g}}$	gaugino m^2	scalar m^2 (5)
I (M) SUGRA (C) MSSM	$\frac{M_{SSB}^2}{\sqrt{3} M_{Pl}} \sim \text{TeV}$ $M_{SSB} \sim 10^{10} - 10^{11}$ GeV	$\left(\frac{\alpha_i}{\alpha_2}\right)^2 M_2^2$	$m_0^2 + \sum_i C_i M_i^2$ \uparrow RG evolution
		Gaugino Unification \uparrow \downarrow	
II GMSB	$\left(\frac{\sqrt{F}}{400 \text{ TeV}}\right)^2 \text{ eV}$	$\left(\frac{\alpha_i}{\alpha_2}\right)^2 M_2^2$	$\sum_i c_i' M_i^2$ \uparrow RG evolution
III AMSB	$\sim 100 \text{ TeV}$	$\left(\frac{b_i}{b_2}\right)^2 M_2^2$ R.G. inv.	$\sum_i a_i b_i \frac{\alpha_i^2}{\alpha_2^2} M_2^2$ R.G. inv.

Weak scale

I, II : $M_1 : M_2 : M_3 = 1 : 2 : 7$

III : $M_1 : M_2 : M_3 = 2.8 : 1 : -8.3$

\downarrow
minimal curve
addition of m_0^2
at some scale

$\beta_i : \beta$ of M . g_{new} which enter RGs = $\frac{-b_i}{(4\pi)^2} g_i^2$

γ_i : Anom. dim. = $-\frac{a_i}{(4\pi)^2} g_i^2$

(m) SUGRA : $m(\tilde{l}) < m(\tilde{q})$ normally ($m_{\tilde{q}}$: mass of 1/2nd gen. squ)
 mostly $\tilde{\chi}_1^0$: bino ; LSP \neq
 $m(\tilde{e}_R) < m(\tilde{l}_L)$

AMSB : $\tilde{\chi}_1^0, \tilde{\chi}_1^{\pm} \approx \tilde{W}_3, \tilde{W}^{\pm}$. Degenerate. Degeneracy lifted only by loop effects.
 $\tilde{\chi}_1^0$ LSP

$\Delta M \equiv m_{\tilde{\chi}_1^{\pm}} - m_{\tilde{\chi}_1^0} < 1 \text{ GeV}$: diff. phenomenology

$\Delta M > 1 \text{ GeV}$: (m) SUGRA situation
 \tilde{l}_L, \tilde{l}_R : degenerate.

GMSB : $\tilde{q} : \text{LSP}$ $2 \times 10^{-5} < m_{\tilde{q}} < 1 \text{ KeV}$
 $\tilde{\chi}_1^0 / \tilde{e}_1 / \tilde{e}_R$: candidates for NLSP

decay length: $c\tau\beta\gamma : 10^{-4} \text{ cm.} \longrightarrow 10^5 \text{ cm.}$

$L = \frac{c\tau\beta\gamma}{\beta\gamma} \propto \frac{1}{(m_{\tilde{\chi}_1^0})^5} (\sqrt{F})^4$



Leptonic colliders which can study sleptons really crucial.

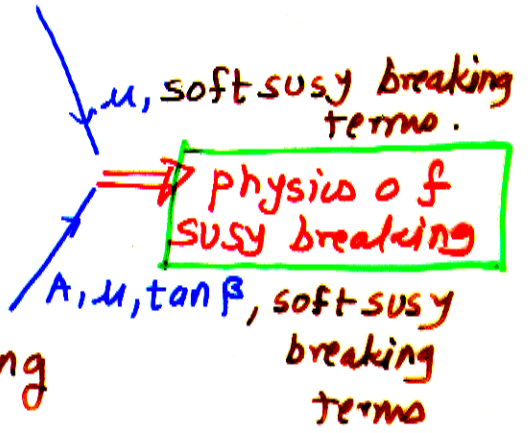
- susy phenomenology in its gory detail looks too complicated. perhaps because we are ^{not} looking at things the right way. Important to try and extract ~~as~~ 'model indep' information as far as possible

We need to

- ① find sparticles, establish quantum nos.
- ② check the interactions and establish coupling equalities implied by supersymmetry.

③ determine sfermion masses

Gaugino-higgsino mixing
Gaugino masses



④ Third generation L-R mixing

LHC will do part of the — things.

(8)

Promises of LHC to SUSY enthusiasts:
(Hopes from LHC of SUSY enthusiasts)

- If SUSY has been realized by nature, LHC will for sure provide proof for it 'Naturalness'

If theories are 'natural' some sparticles must be accessible at LHC. (Gauginos/higgsinos)

Barbieri, ...
Anderson, ...
Feng, Moroi, ...
Allanach, ...

- The strongly interacting heavier sparticles are produced first, lighter sparticles in decays. 'top-down'.
- Rates are very high. 'discovery of SUSY not a problem'. Transp.
- Worst background for SUSY is SUSY itself.
- Very nice analysis methods developed make possible measurement of a variety of sparticle masses but not quantum nos.
- heavier gauginos not accessible.
- Over a large region of parameter space only one higgs (lighter) might be visible. Transp.
- measurements of $m_{\tilde{q}}$, $m_{\tilde{g}}$, $m_{\tilde{\chi}^\pm}$ possible with good degree of accuracy. summarized later
- Possible to get a good idea of eff-SUSY breaking scale. Transp.

2) Accurate information about SUSY breaking scale and mechanism seems generally not easily extractable.

2) Model indep. determination of $m_{\tilde{\chi}_2^0} \approx 10\%$ level.

3) Reach in sleptons limited unless sleptons are produced in decays

4) Analysis and determination of many parameters model dependent. Within a model a high deg.

of accuracy. Model indep. analyses have begun.

Paige, Hinchliffe...

look promising -----

5) If $m_{\tilde{l}_R} < m_{\tilde{\chi}_2^0}$, LHC studies may be able to

determine all the parameters to predict density of thermal LSP relics.

Drees, Nojiri,

[Gaugino/Higgsino content of the neutralino $\tilde{\chi}_2^0$ is what one wants to determine.]

This information is also crucial in LHC analyses
LC \Rightarrow LHC.

A high $\&$ luminosity machine to make precision measurements of sparticle properties.

BUT

We need

$$\sqrt{s} > 2m_s + 30 \sim 40 \text{ GeV.}$$

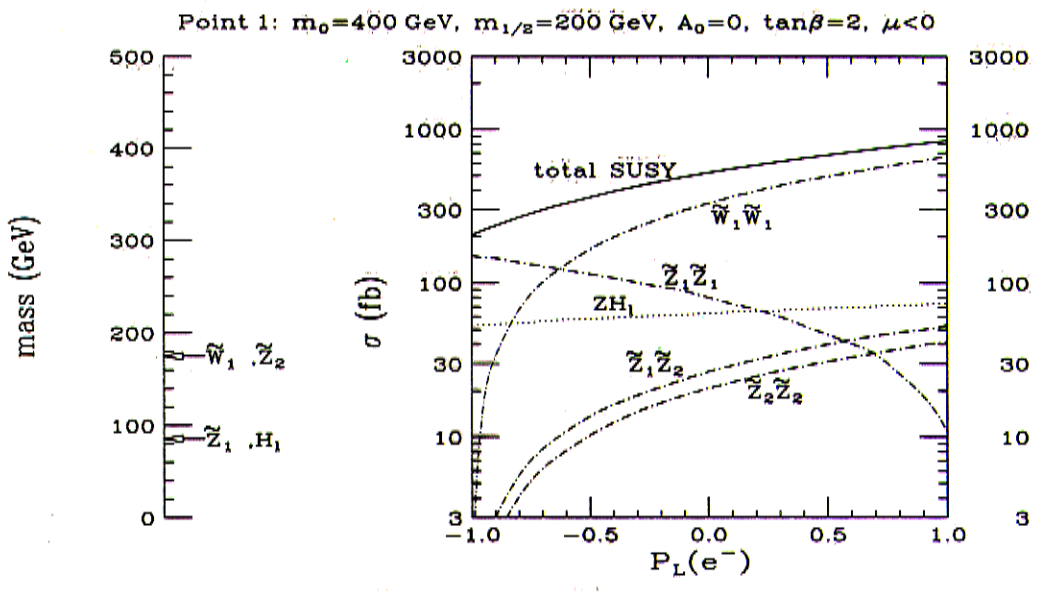
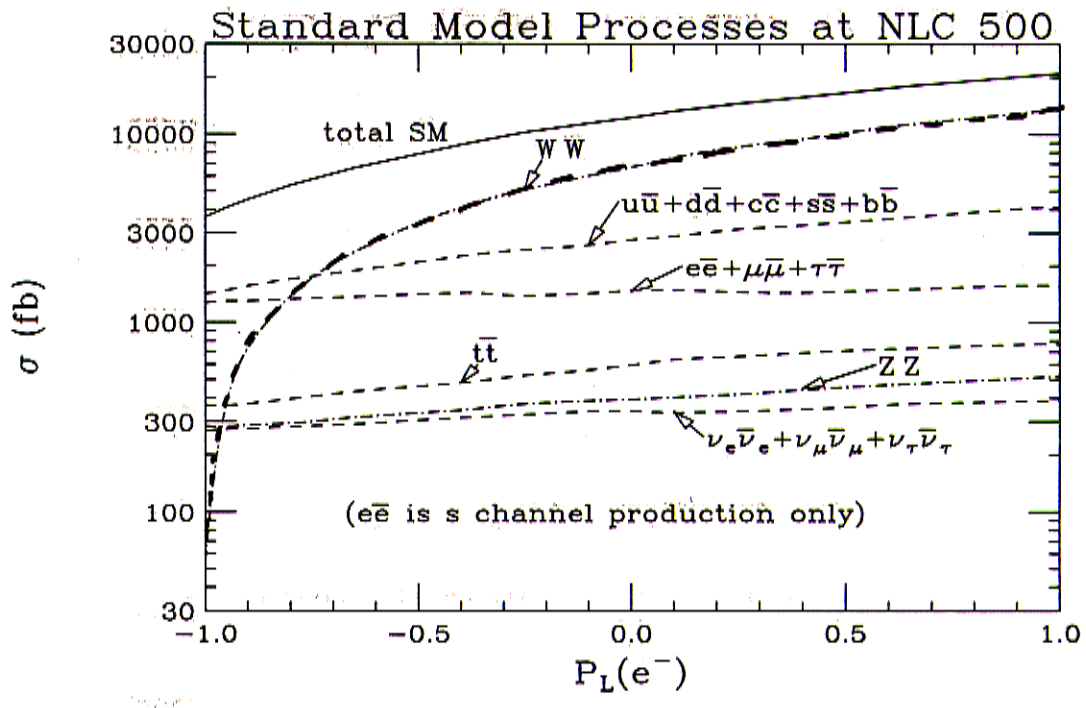
Hope: To build on the results obtained by LHC

mixing

- a) Determine sparticle masses to high precision
- b) Determine quantum nos. such as spin, hypercharge
- c) Measure equality of couplings predicted by SUSY.
- d) with (a) + info. from LHC \Rightarrow ^{info. about} SUSY breaking at high scale.

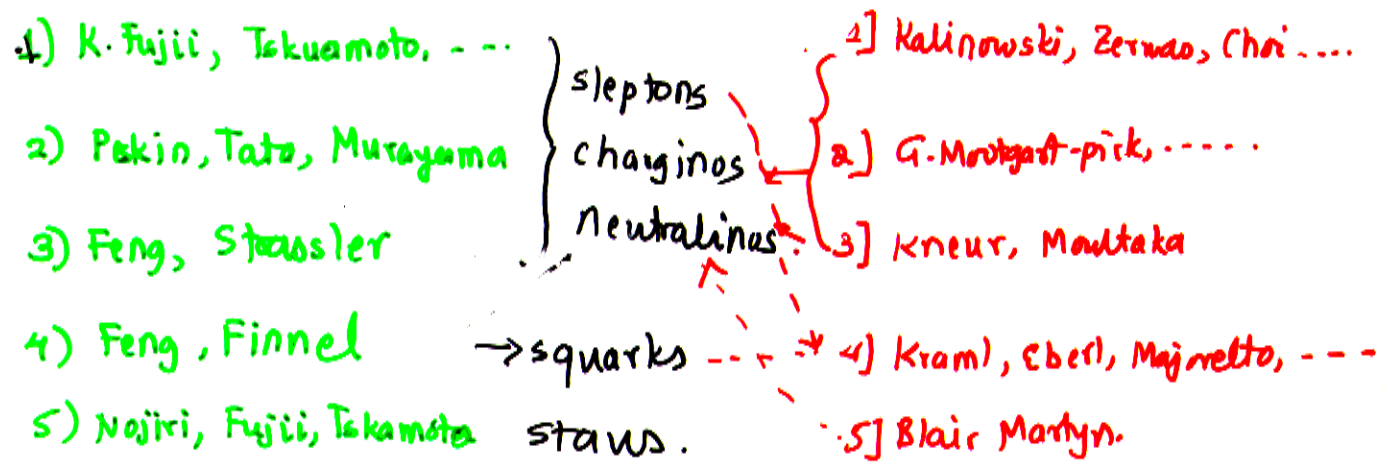
What are the special features which help extra?

- 1) Energy can be tuned.: bottom up approach. Start with lightest. Input from LHC will of course help in choices of energy.
- 2) Use of polarisation: SUSY deals with chiral fermions & their partners. Imp. info. can be projected. WW bkgd. can be killed. useful for (a)
- 3) Running in $e^-e^- / \gamma\gamma / e\gamma$ mode can add to useful information example later



What and how well can LC measure?

* A large body of **old** and **new** work.

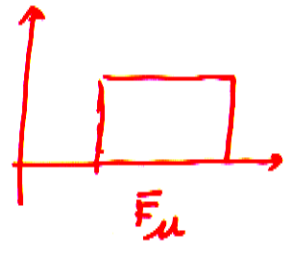
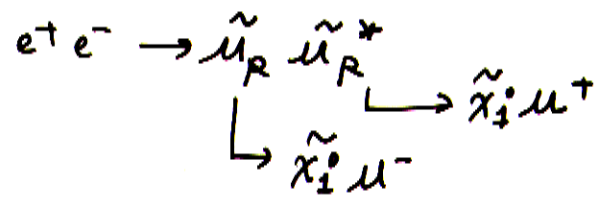


- some are full fledged MC analyses
- some use only statistical errors
- some use statistical errors + conservative estimates of efficiency
- All use polarisation crucially.
- Use

Examples

i) Mass determination from end pts.

Transp.



TESLA, high luminosity : $\delta m_{\tilde{u}_R} \approx 0.3\%$

ii) Threshold Mass scan Fermions $\sim \beta$, scalars $\sim \beta^3$

Transp.

$$\delta m_{\tilde{u}_R} \approx 0.2\%$$

\tilde{u}_E : degradation due to detection eff. of τ .

Advantage of threshold scan remarkable in this case.

iii) Study of higher order effects is now state of the art.

Martyn : finite width of \tilde{u}

Freitas : full calculation of $e^+e^- \rightarrow \mu^+ \mu^- \tilde{\chi}_1^0 \tilde{\chi}_1^0$

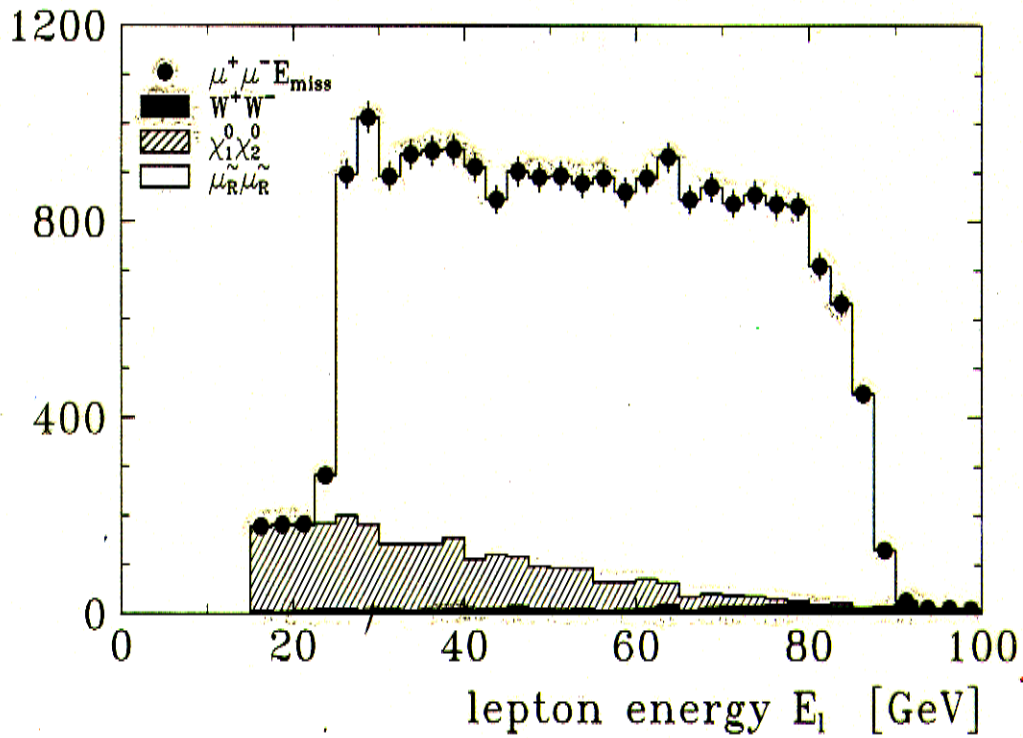
Drees, Ebeli, Kraml, R.G : Effect of higher order

corrections on $m_{\tilde{q}}$ determination

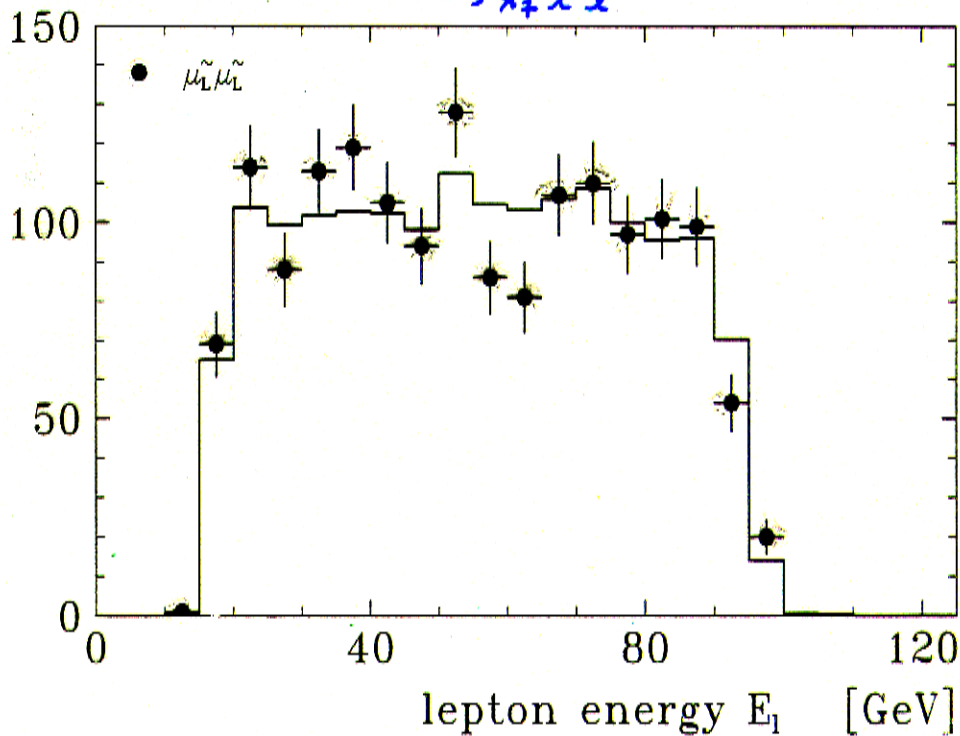
Transp.

For TESLA : $P_{e^-} = 0.8$, $P_{e^+} = 0.6$, Int. Lum. $\sim 500 \text{ fb/yr}$

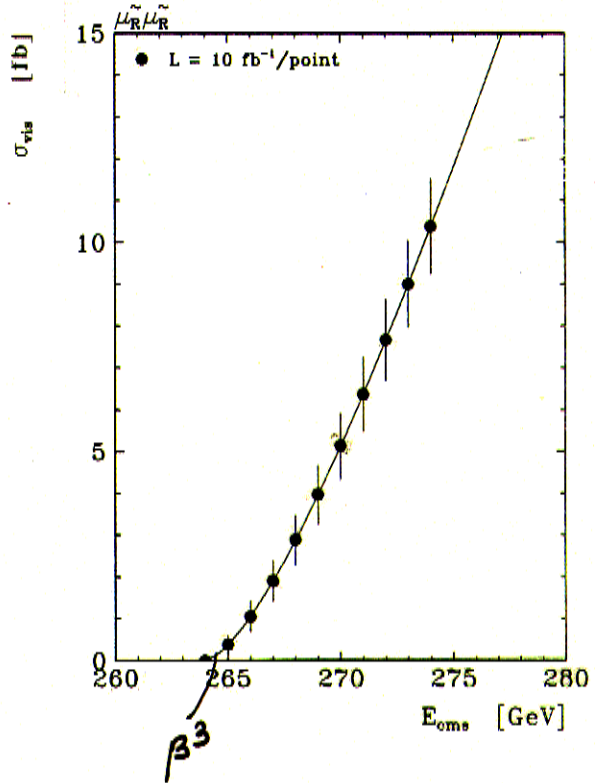
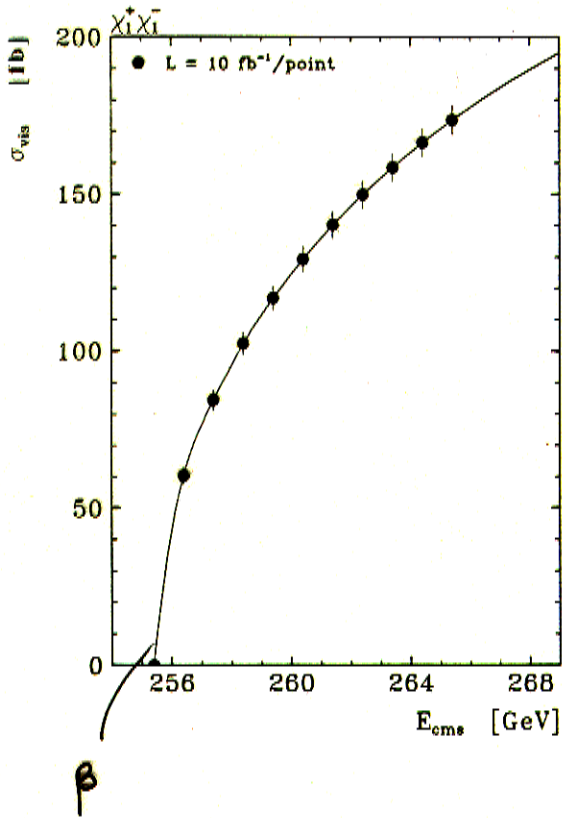
Matyn, Blair



$\tilde{\mu}_L \tilde{\mu}_L^* \rightarrow \mu \tilde{\chi}_2^0 \mu \tilde{\chi}_2^0$
 $\rightarrow \tilde{\chi}_2^0 e^+ e^-$ 6 lepton final states

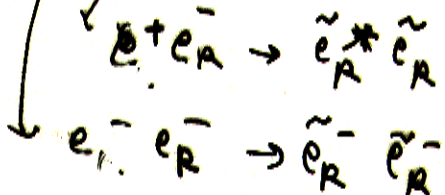
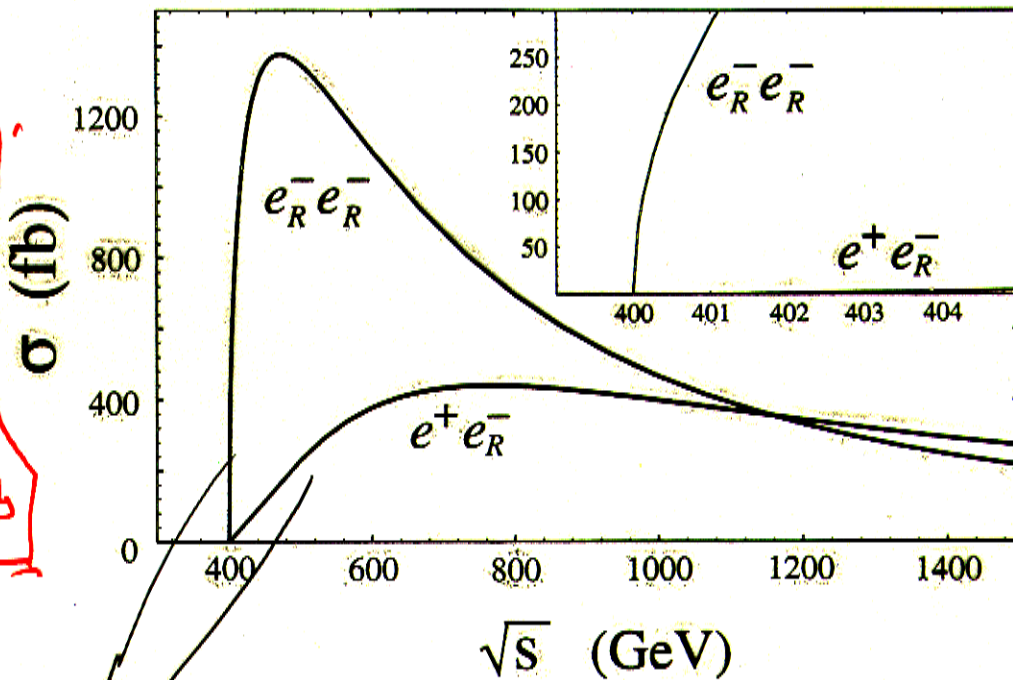


Marty, Blair



Feng

This fig. does not include Beamstr. effects, ISR effects



precision for mass measurement ↑ using $e^- e^-$

Kinematic mass determinations at TESLA
 for threshold scans 10 fb^{-1} per point

Blair, Martyn

	δm	m
$\tilde{\mu}_R$	0.09	132.0
$\tilde{\mu}_L$	0.4	176.0
$\tilde{\nu}_\mu$	0.8	160.6
\tilde{e}_R	0.05	132.0
\tilde{e}_L	0.18	176.0
$\tilde{\nu}_e$	0.07	160.6

Assume universal masses

$\Rightarrow m_0, m_{1/2}, \tan\beta$

$\tilde{\nu}_2, \dots$

$\tilde{\tau}_1$	0.6	131.0
$\tilde{\nu}_\tau$	0.6	160.6

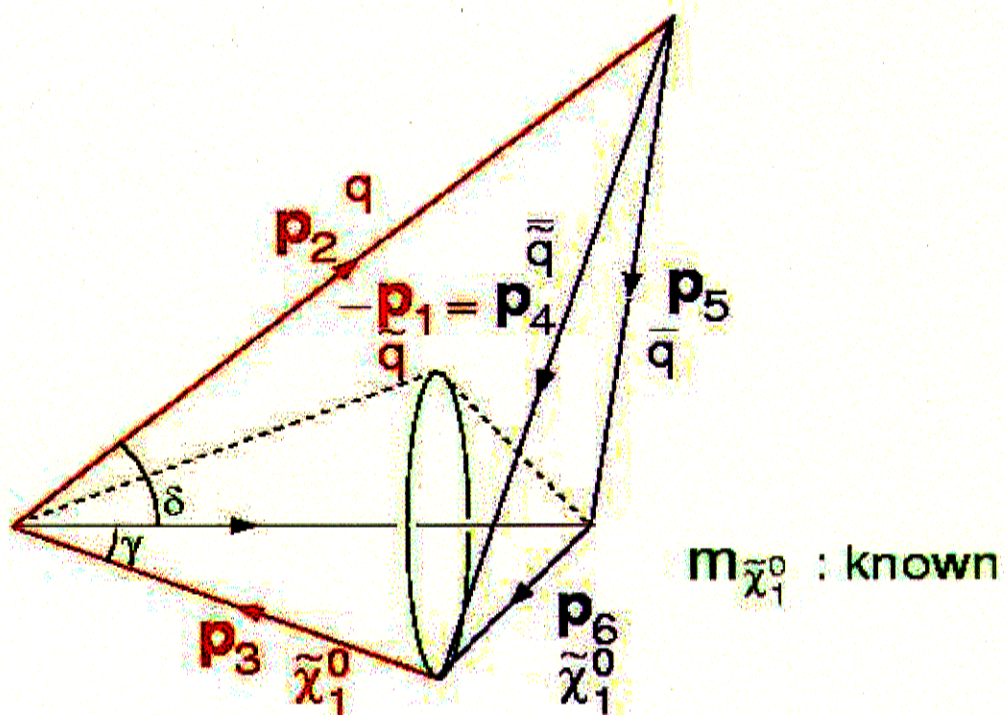
$\Rightarrow m_0, m_{1/2}, \mu, \tan\beta$

$\tilde{\chi}_1^\pm$	0.04	127.7
$\tilde{\chi}_2^\pm$	0.25	345.8

$\Rightarrow M_2, \mu, \tan\beta$

$\tilde{\chi}_1^0$	0.05	71.9
$\tilde{\chi}_2^0$	0.07	130.3
$\tilde{\chi}_3^0$	0.30	319.8
$\tilde{\chi}_4^0$	0.52	348.2

$\Rightarrow M_1, M_2, \mu, \tan\beta$

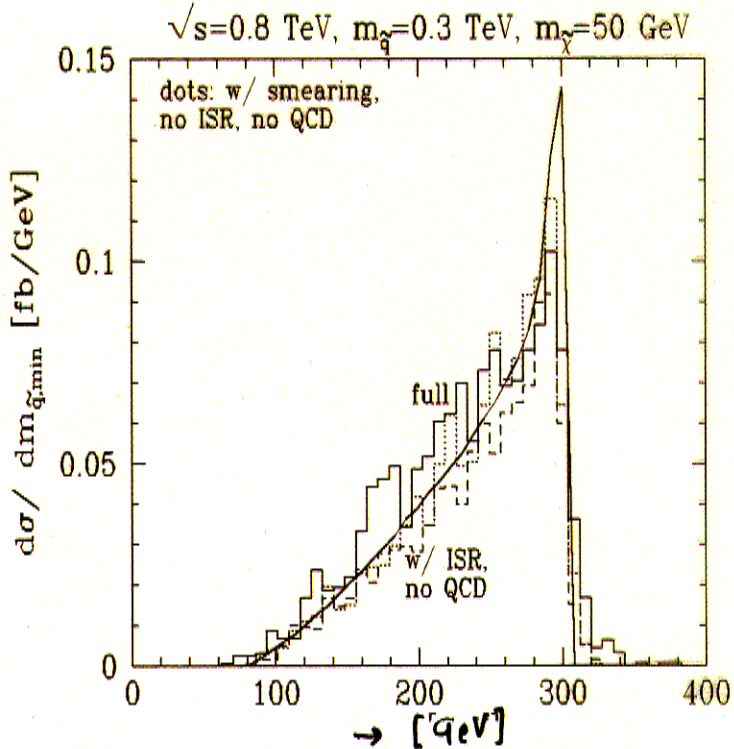
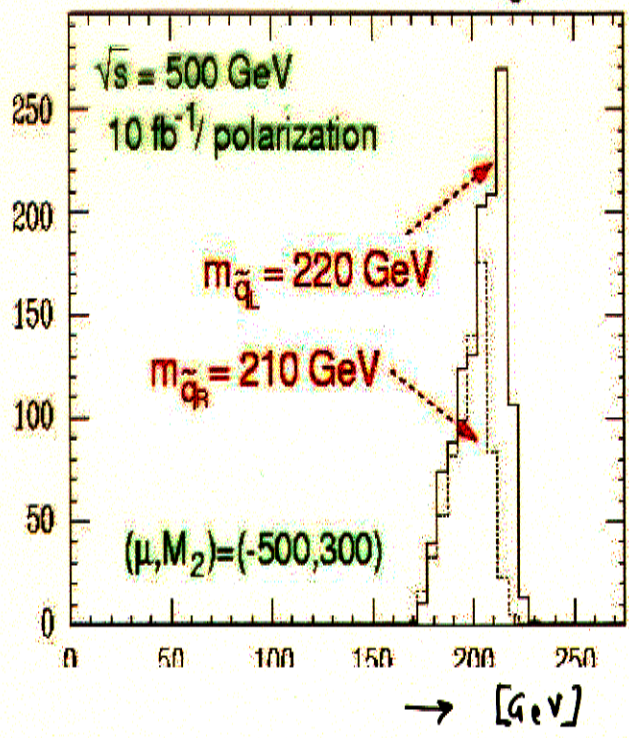


distⁿ. in minimum kinematically allowed mass $m_{\tilde{q}, \min}$ peaks at true $m_{\tilde{q}}$.



Drees, Godbole, Ebeli, Kraml

Feng-Finnell



Effect of rad. corrections.

Third generation squarks, sleptons :

Want to determine

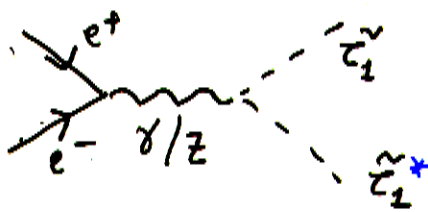
$\theta_{\tilde{e}}, \theta_{\tilde{\tau}}, m_{\tilde{\tau}_1}, m_{\tilde{\tau}_2} \Rightarrow$ clue to high scale phys.

$\tilde{\tau}$: (Nojiri)

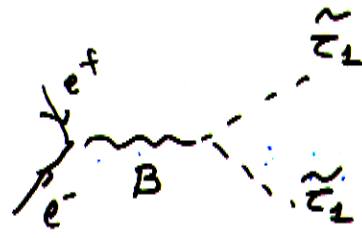
* $\cos \theta_{\tilde{\tau}}$: mixing angle $\tilde{\tau}_2 = \cos \theta_{\tilde{\tau}} \tilde{\tau}_L + \sin \theta_{\tilde{\tau}} \tilde{\tau}_R$

Use polarized beams.

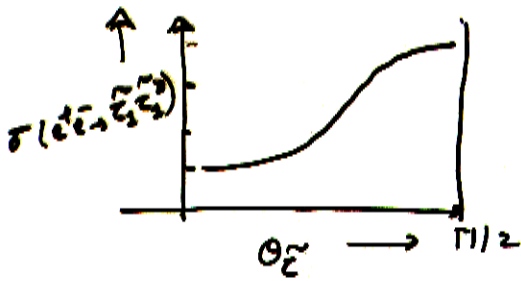
(say) $P_e = 1$



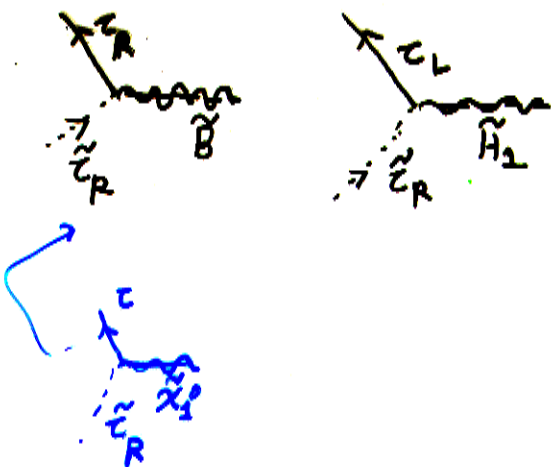
$$\frac{P_e = 1}{\sqrt{s} \gg m_Z}$$



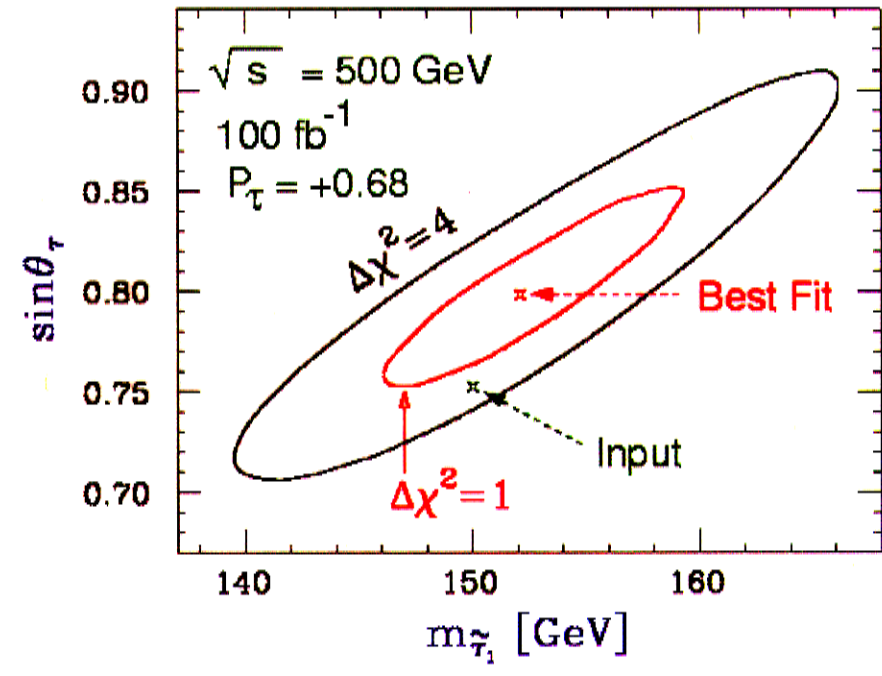
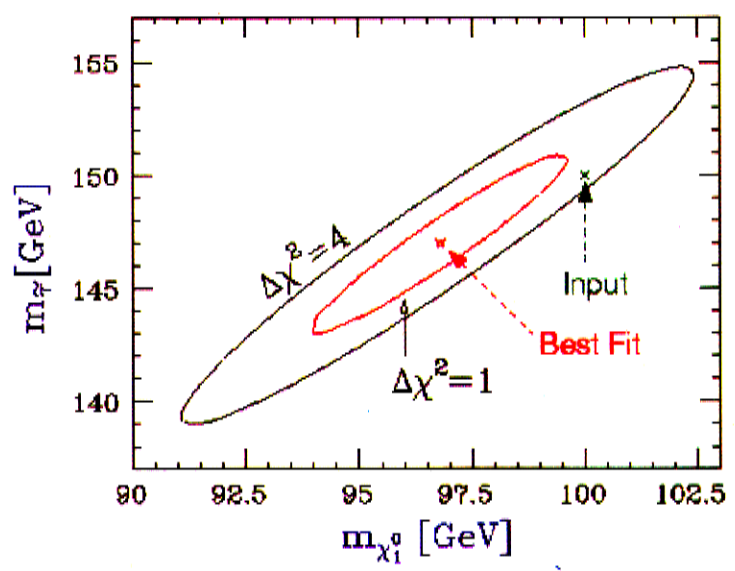
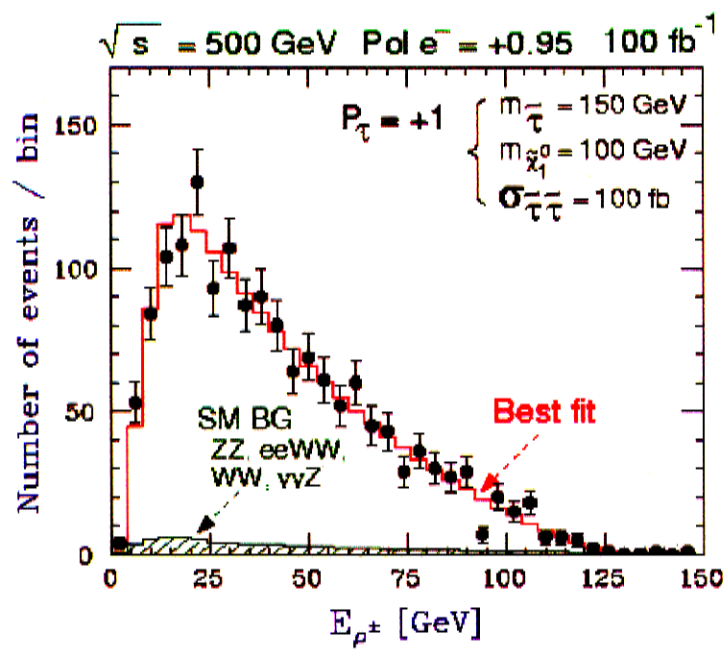
$$4 \sigma(\tilde{\tau}_L) = \sigma(\tilde{\tau}_R)$$



* $\tilde{\tau} \rightarrow \tau \tilde{\chi}_\pm^0$

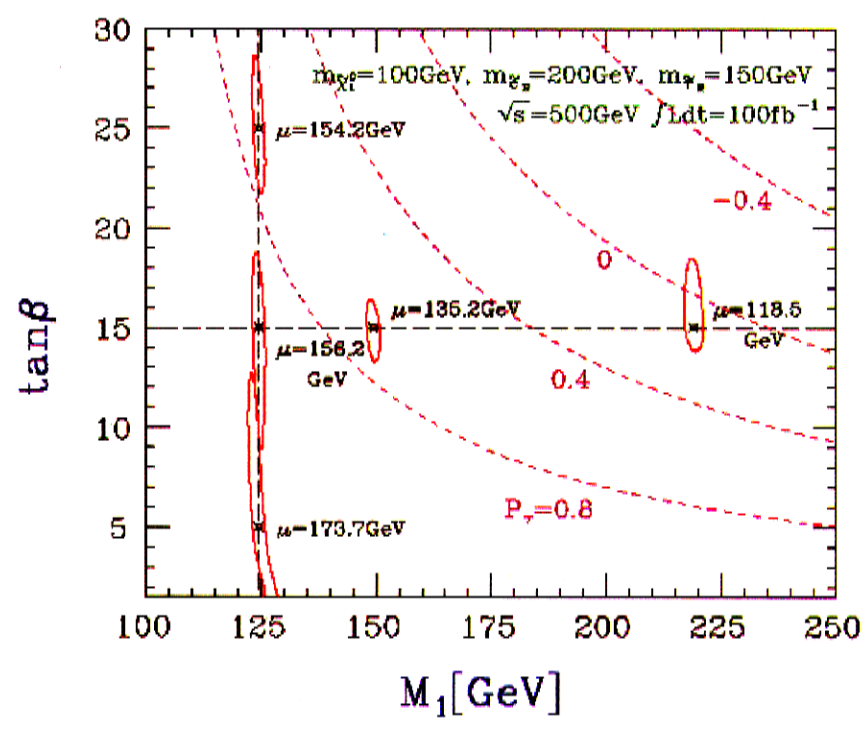
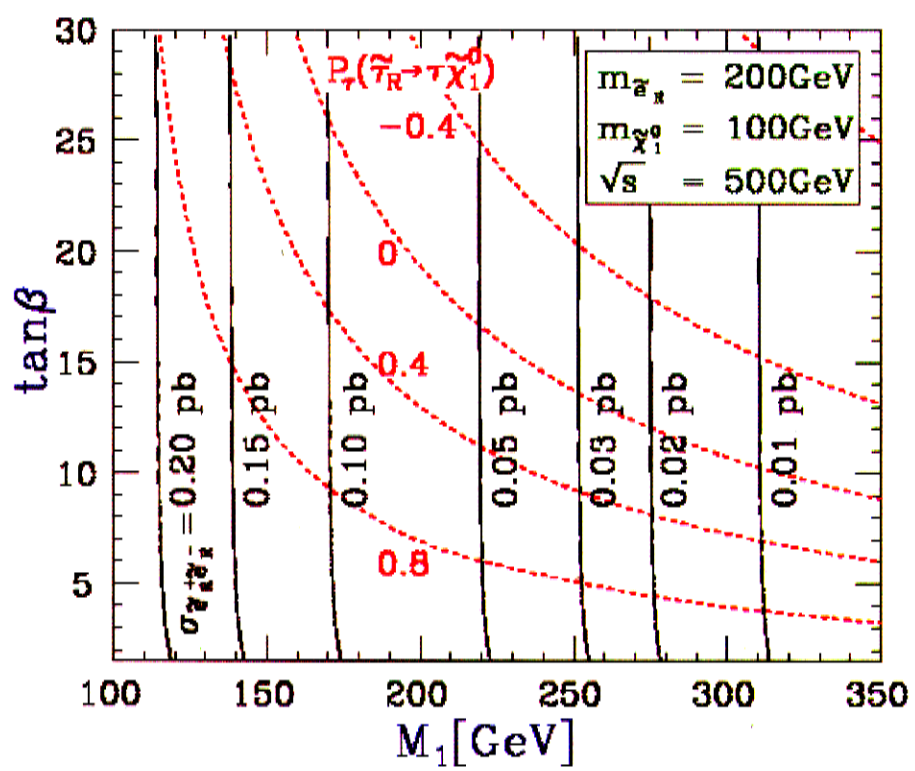


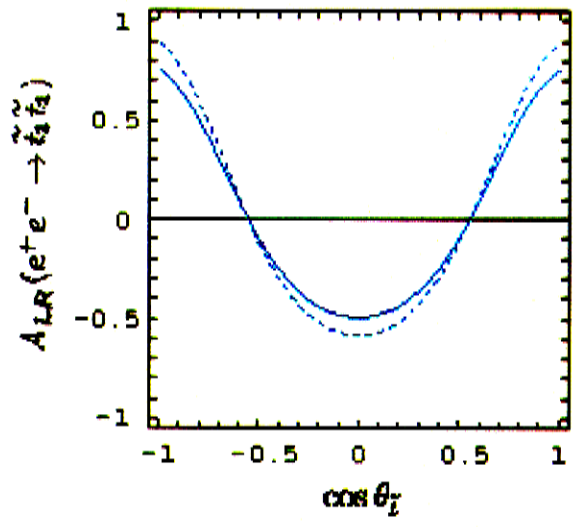
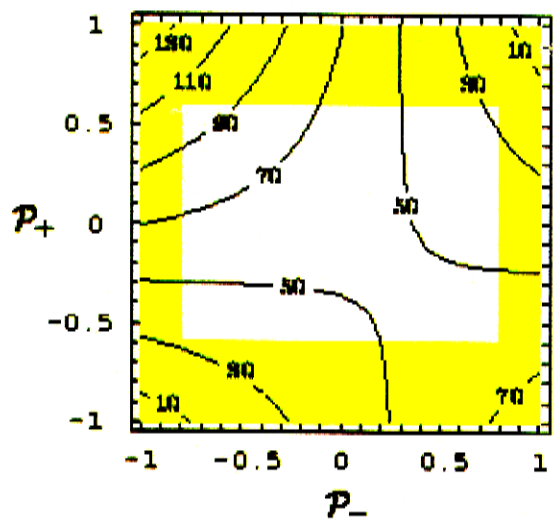
\therefore polarisation of τ sensitive to $\tan \beta$.



$\Delta(\sin\theta_{\tau}) < 0.03$ (if one ~~was~~
assumes $m_{\tilde{\chi}_1^0}$ known $\Rightarrow \Delta m_{\tilde{\tau}_1} \downarrow$)

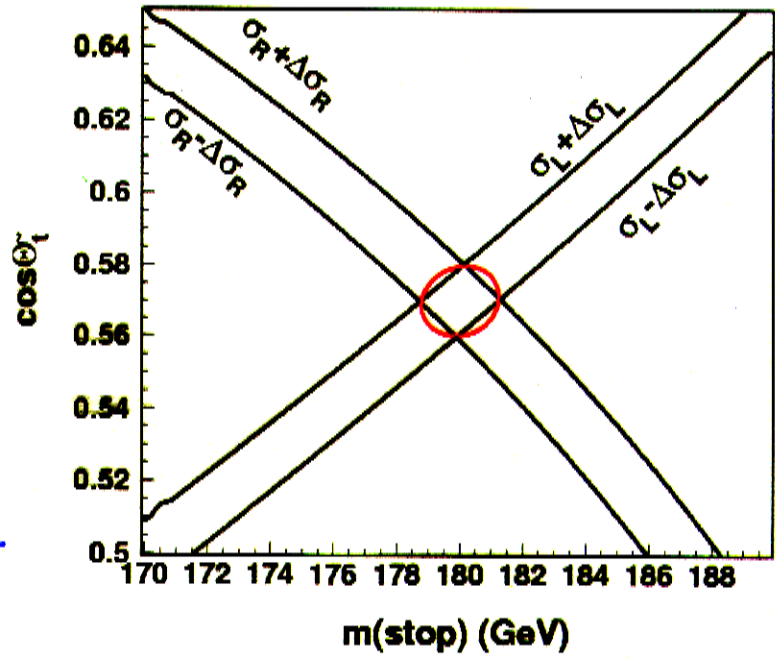
$$\sigma(e^+e^- \rightarrow \tilde{e}_L \tilde{e}_R^*), P_\tau(\tilde{e} \rightarrow \tau \tilde{\chi}_1^0)$$





TESLA

stop into c neutralino 80/60 pol



$\Delta(\cos \theta_t) = 0.003$
 $\Delta(m_E) = 0.8$

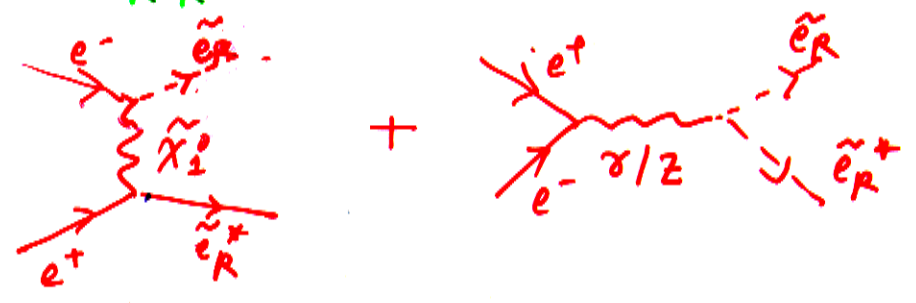
If $m_{\tilde{E}}$ in this range
 LHC analysis?

- 1) Kraml et al
- 2) A-Soparkale ... MC studies

- 1: polarization dep. of σ -sections.
2. spin 0 of \tilde{U} can be inferred by reconstructing angular dist. of \tilde{U} . (Fujii et al)
3. Equality of couplings:

Nojiri

1) $g_{\tilde{e}_R e_R \tilde{B}} = g_{B e e} = \sqrt{2} g \tan \theta_w = g' \sqrt{2} = \sqrt{2} g' Y_B$



$\Delta Y_b \approx 1\% \quad (100 \text{ fb}^{-1})$

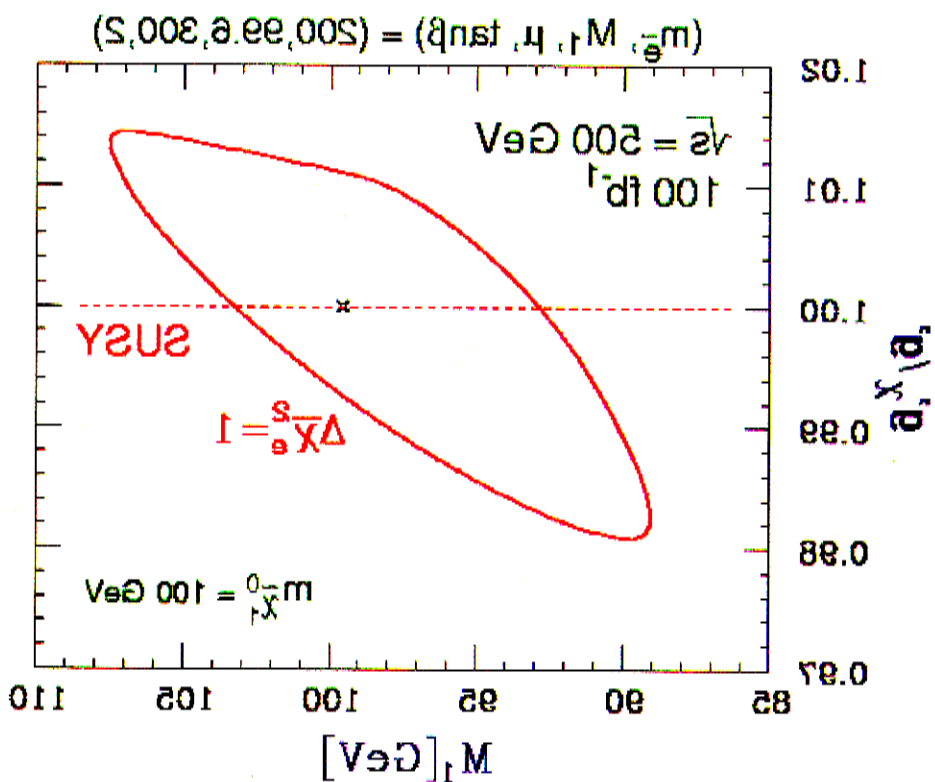
Fig.

size comparable to the SUSY oblique corrections to coupling equalities **

** Nojiri, Pierce... } get back later.
 Feng, Polonsky, ... }

2) $g_{e \nu W} = g_{e \tilde{\nu} \tilde{W}}$ Test (Choi, Zerwas, ...)

for the point they considered $\pm 0.1\%$ test.
 1 ab^{-1}



M. Nojima, R. Kitajima, T. Kuroki

$$e^+ e^- \rightarrow \tilde{\nu}_e \tilde{\nu}_e^*$$

$\frac{d\sigma}{d\Omega}$
 sensitive to M_1

$$\tilde{\nu}_e \tilde{\nu}_e^* \rightarrow b \bar{b}$$

Higgsino / Gaugino Content

Lightest chargino $\tilde{\chi}_1^+$

Measurements of $m_{\tilde{\chi}_1^+}$, $m_{\tilde{\chi}_1^0}$ and

$\sigma_{\tilde{\chi}_1^+ \tilde{\chi}_1^-}^{\text{tot}}$ and (i) Angular correlations among $\tilde{\chi}_1^+$ decay products
(from polarization of $\tilde{\chi}_1^+$)

OR (ii) polarized cross-sections $\sigma_R^{11}, \sigma_L^{11}, \sigma_{R,L}^{22}, \sigma_{R,L}^{22}$

1 ab^{-1}

\Rightarrow Determination of $\cos 2\phi_L, \cos 2\phi_R \lesssim 1-3\%$

[Choi, Zerwas, Kalinowski, ...]

ϕ_L, ϕ_R : chargino-higgsino mixing angles.

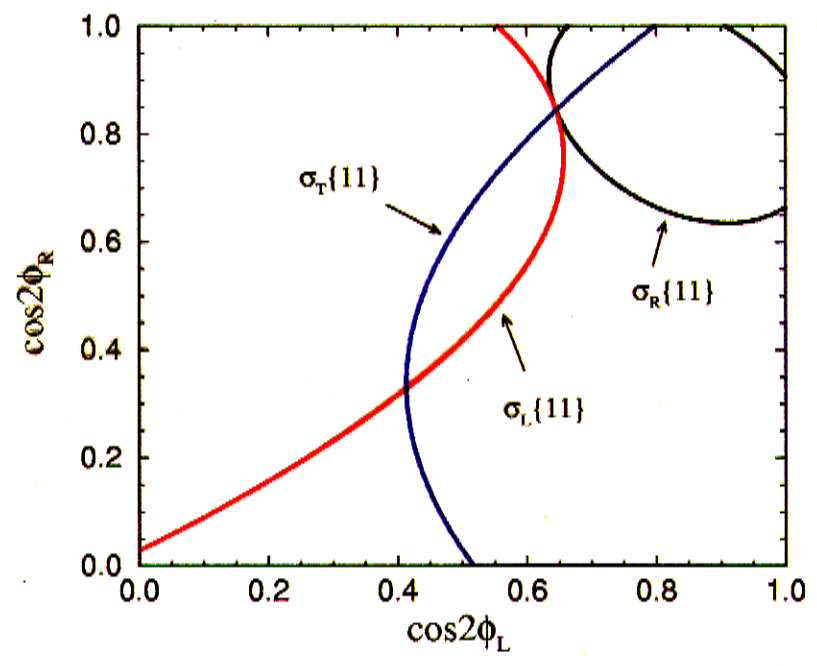
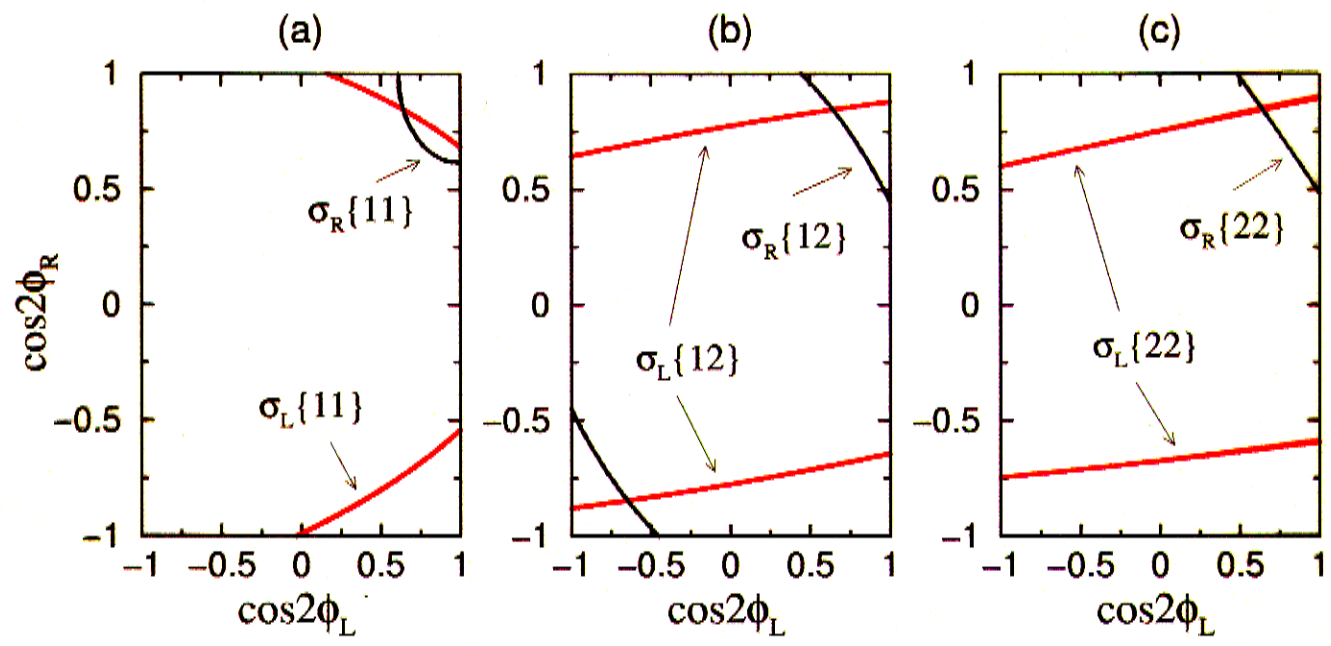
\Rightarrow Reconstruction of $\tan \beta, M_2, \mu$.

Since all the variables are proportional to $\cos 2\beta$ accuracy of $\tan \beta$ determination not so good at high $\tan \beta$

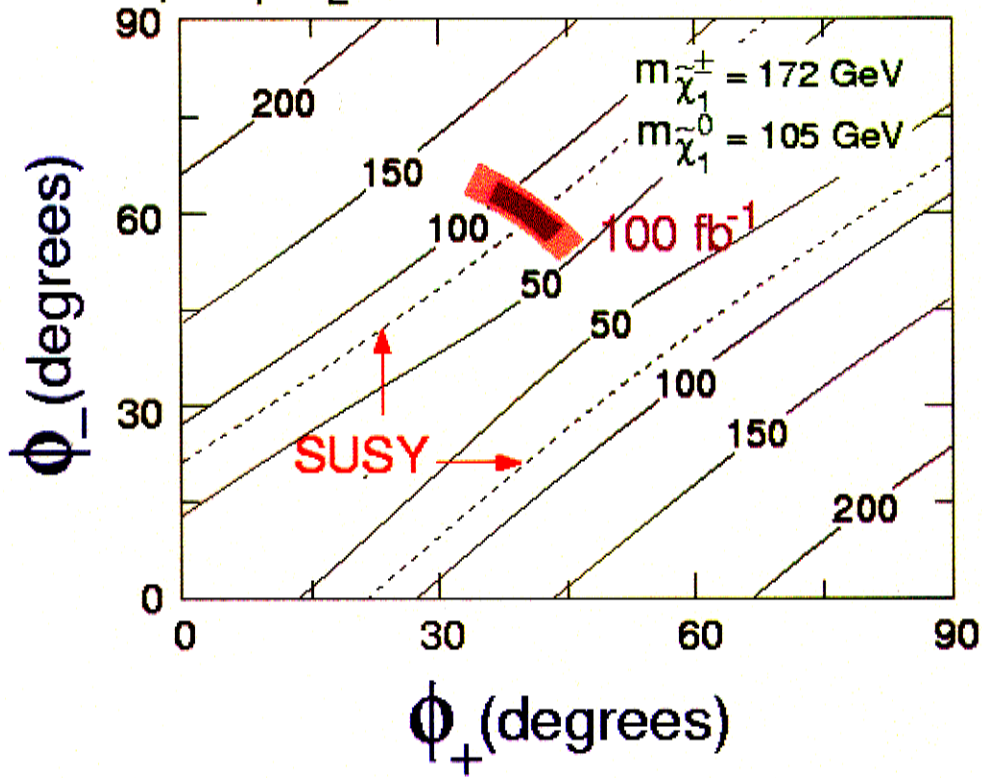
At high $\tan \beta$ $\tilde{\chi}_1^+$ studies \Rightarrow info. about $\tan \beta$
 \hookrightarrow Nojiri

Higgs sector $\Rightarrow \tan \beta$
 \hookrightarrow Feng ...
Han ...

LHC $\rightarrow \tan \beta$



$(m_{\tilde{t}}, m_{\tilde{q}}, M_2, \mu, \tan\beta) = (400, 700, 210, -195, 4)$



Fuji, Feng, Perkin, Tata:

Similar analytical work done for
neutralinos.

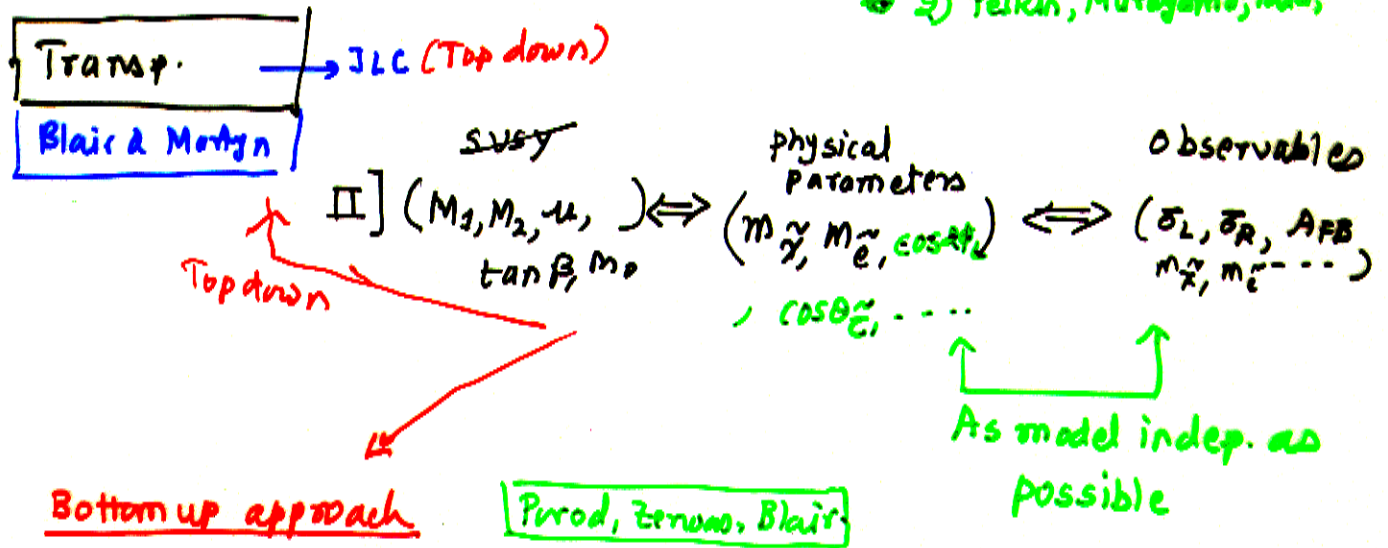
- (i) possible to extract $M_1, M_2, \mu, \tan \beta$ from measurements
Kneur, Mautaka
- (ii) polarization of e^- and e^+ beams can
be used to extract information on M_2 .
Increase accuracy of the info. G. Mootgat-Pick R.....
- (iii) studies where $\tilde{\chi}_2^+$ & $\tilde{\chi}_3^0$ are almost degenerate
wino's ^{pheno.} studies have been performed. Look very
promising
S. Mrenna, Jo-Gunnar.
Prabir Roy, D. Ghosh. --

Determination of susy breaking scale.

+ susy breaking parameters at high scale

Two approaches: I] (susy breaking parameters at high scale $(M_1, M_2, \mu, \tan \beta, m_0)$) \iff (observables $(\sigma_L, \sigma_R, AFB, \dots)$)

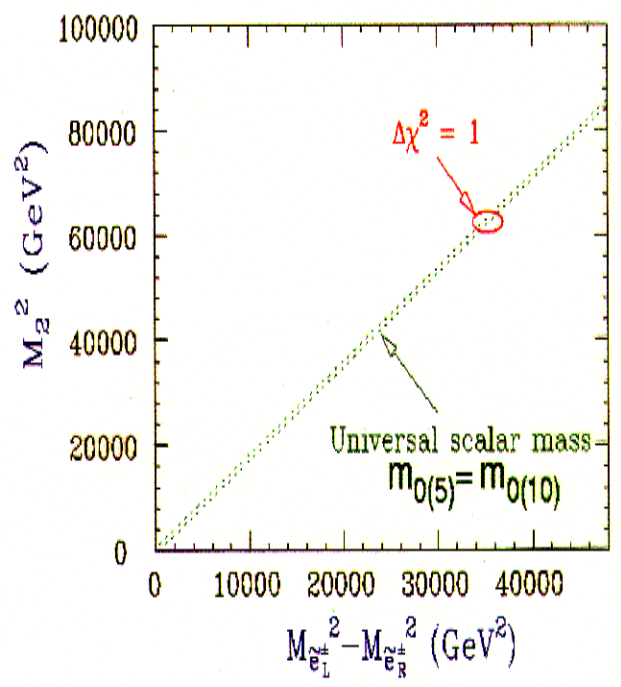
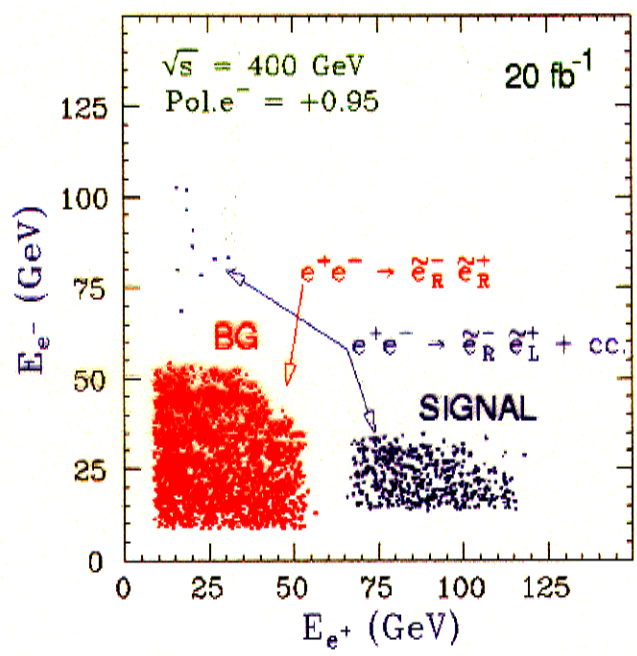
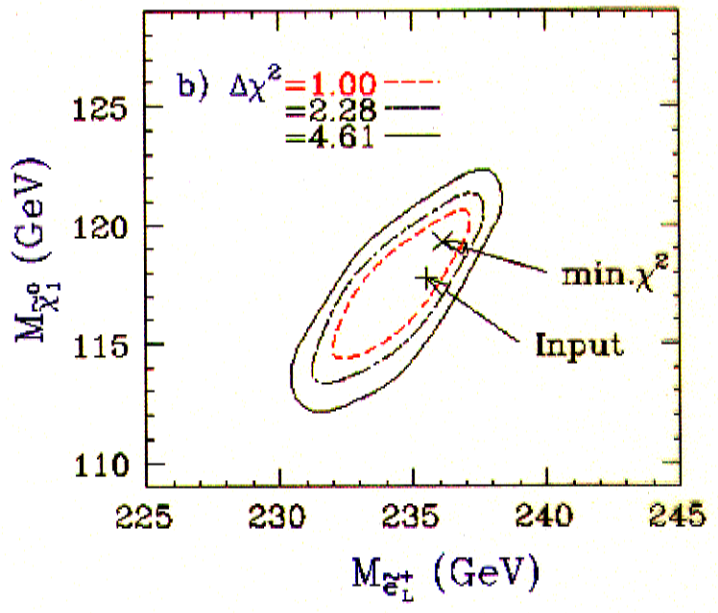
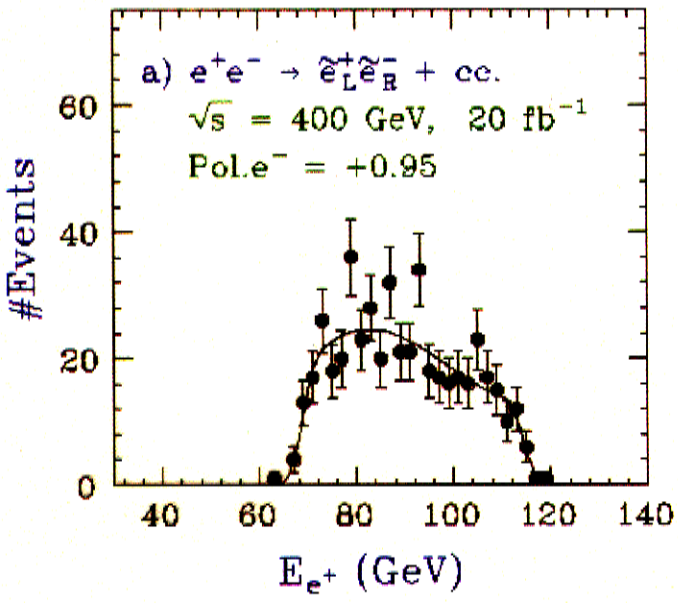
- 1) Feng, Murayama, Fujii, ...
- 2) Perkin, Murayama, Tate, ...



To learn about the scale now, use RG eqns.
 use measured sparticle masses at EW scale to extract susy parameters at EW scale. Try to reconstruct the parameters M_i, M_{scalars} at high scale.

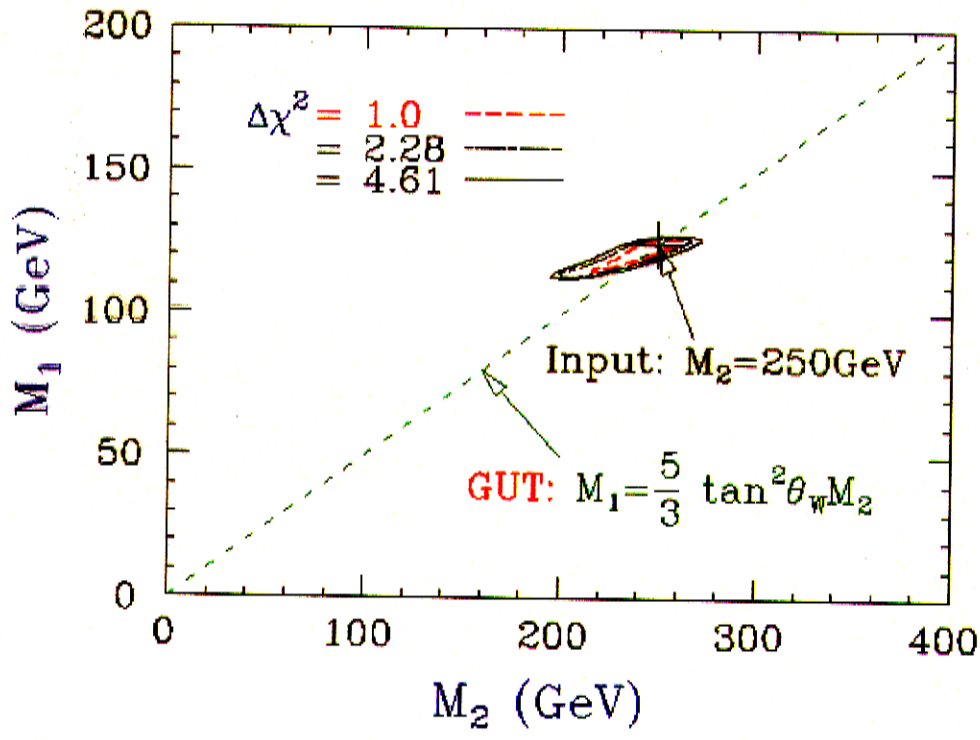
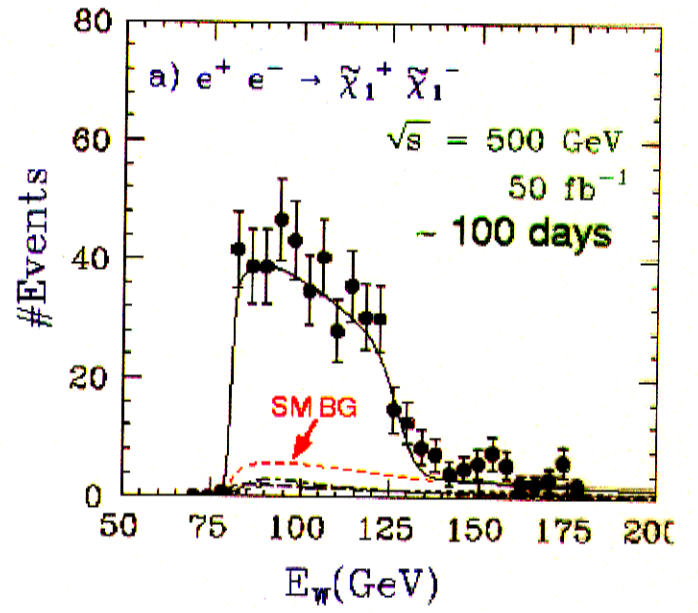
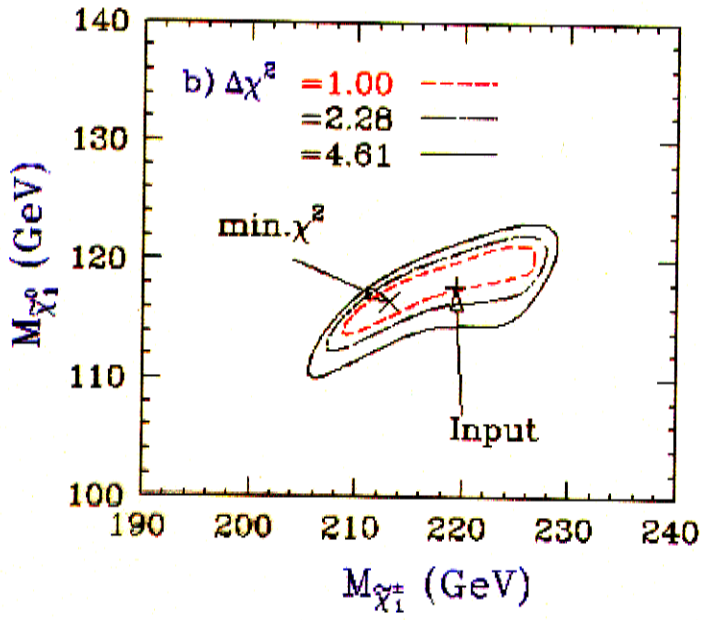
F	Table.	$d_i = z_i d_u$	
		$M_i = z_i M_{1/2}$	— gaugino
		scalars $M_j^2 = M_0^2 + c_j M_{1/2}^2 + \sum_{\beta=1,2} c'_{j\beta} \Delta M_{\beta}^2$	
		$A_k = d_k A_0 + d'_k M_{1/2}$	

lepton
 First two generation scalar masses & gaugino masses unify nicely. squarks, H_2 have larger errors



Assuming a universal mass M_0 .

one expects. $M_{\tilde{e}_L^+}^2 - M_{\tilde{e}_R^+}^2 = 16 M_{1/2}^2 - \frac{1}{2}(1-4\sin^2\theta_W)\cos 2\beta M_Z^2$



$m_{\tilde{\chi}_1^\pm}, \sigma_R(\tilde{\chi}_1^\pm), \text{angular dist}^n, \sigma_R(\tilde{e}_R), m_{\tilde{\chi}_1^0}$

(31)

Blair, Marty n: Top down approach.

Use all the accurately measured masses from threshold scan.

Assume unification of scalar masses.

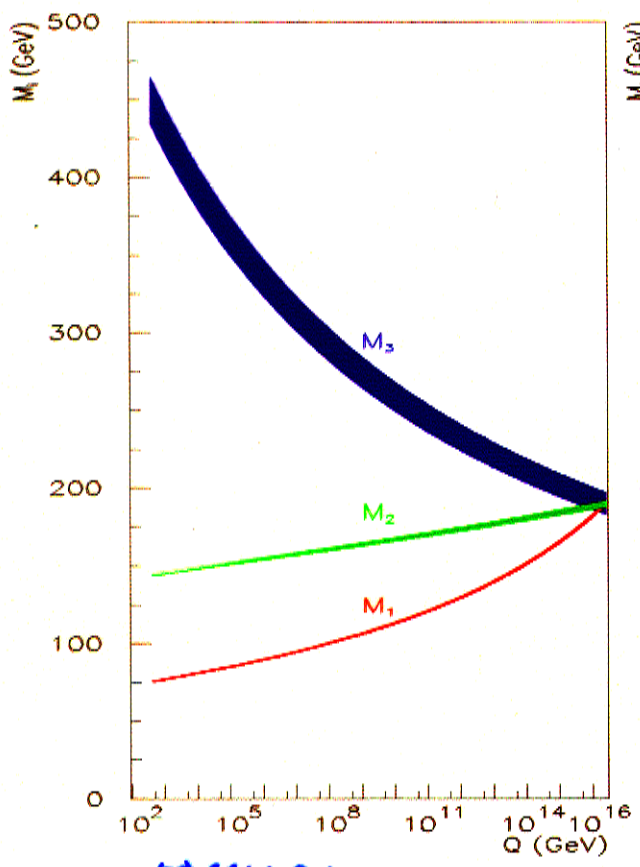
Fit $(M_1, M_2, m_0, A_0, \tan\beta)$,

	True value	Error
m_0	100	0.09
m_{12}	200	0.10
A_0	0	6.3
$\tan\beta$	3	0.02

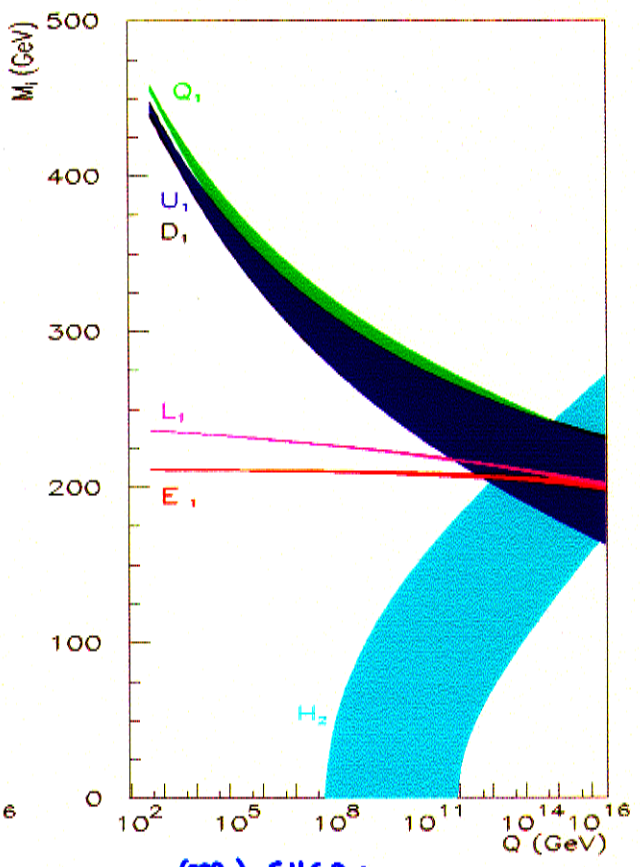
	True value	Error
m_0	100	0.09
M_1	200	0.20
M_2	200	0.20
A_0	0	6.3
$\tan\beta$	3	0.04

High luminosity TESLA and (higher energy)
 $\sqrt{s} = 500 \text{ GeV}$

Errors on initial values LC+LHC

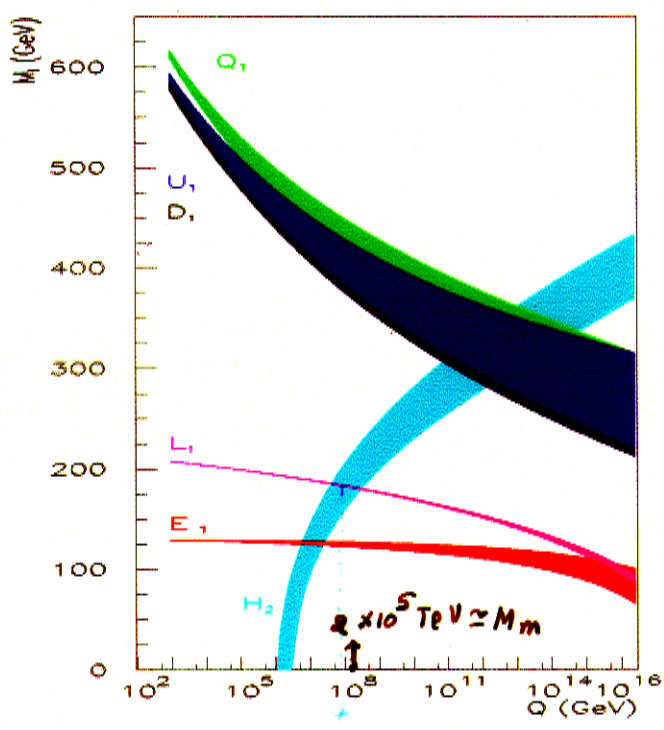


(M) SUGRA.



(M) SUGRA.

1 TeV



→ 1.5 TeV

$2 \times 10^5 \text{ TeV} \approx M_m$

GMSB

Bottom-up approach:

Possible to determine the soft ~~SUSY~~ parameters at high scale 'directly' by ^{R.G.} evolution.

i) Gaugino masses \longrightarrow

ii) 1st, 2nd generation scalar leptons \longrightarrow

Quarks } \longrightarrow 'larger errors'
Higgs }

iii) 3rd generation & A_k : existence of pseudo fixed point increases errors reconstruction of M_0 .

High Accuracy possible at LC (TESLA, high luminosity) absolutely essential input to such studies.

Determination of τ scale:

(34)

GMSB:

Ambrosanio, Blair

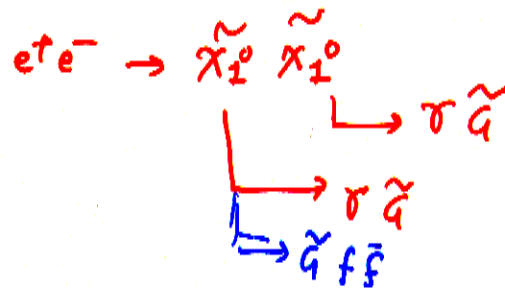


Fig.

Use 'charged decays' of $\tilde{\chi}_1^0$ also

Delayed decay of $\tilde{\chi}_1^0 \Rightarrow$ nonpointing photon.

1) $\gamma\gamma + E$
Central hard τ^s

prompt decay

$\tau \ll 1 \text{ cm}$

2) γ " γ " E
" γ " γ " E

delayed decays

$1 \text{ cm} < \tau < 100 \text{ cm}$

3) E , " γ "

$10 < \tau < 10^5$

Fig.

$\tilde{\chi}_1^0$ mass can be measured
with E_T measurement

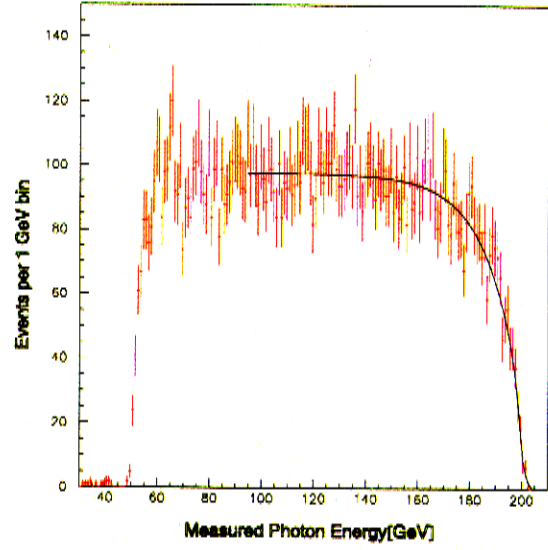
Fig.

Using known value of $m_{\tilde{\chi}_1^0}$, determination of decay length ' λ ' on an event by event basis.

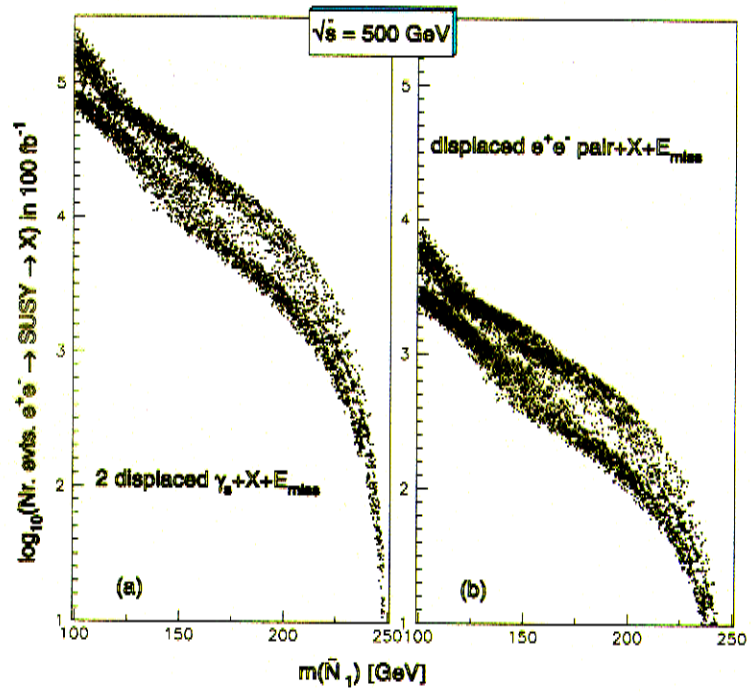
Life time can be determined to $\sim 10\%$ over the entire range expected in GMSB models.

$$\tau \ll \frac{1}{m_{\tilde{\chi}_1^0}^5} (\sqrt{F})^4$$

Model 2 - Photon Energy Spectrum



Inclusive QMSB Signals at the Linear Collider



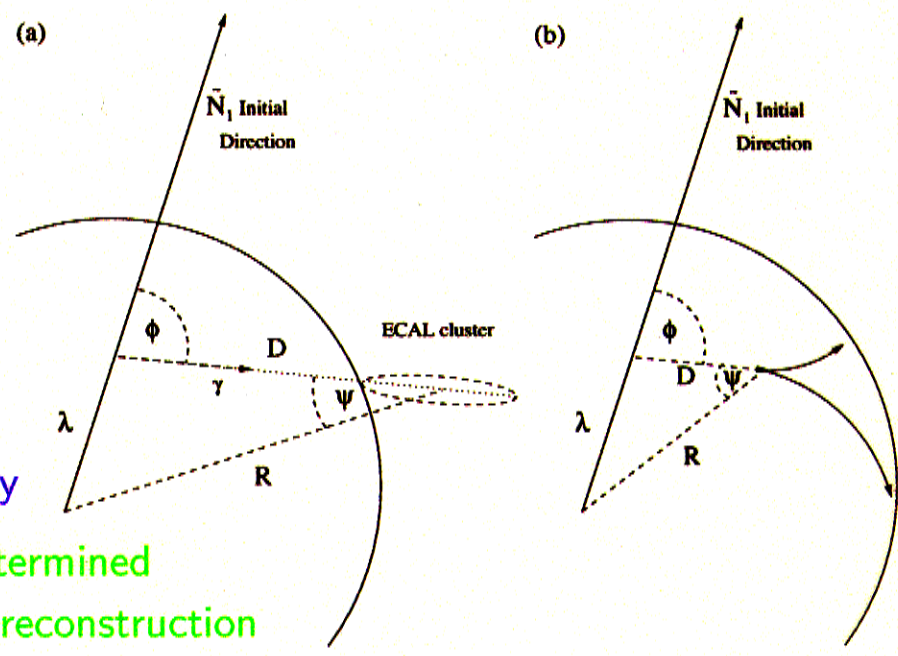
Neutralino NLSP Event Signatures @ LC

◇ Topology of a \tilde{N}_1 decay (one side of each event):

λ = decay length
can be measured
(event-by-event)

ϕ calculated from
measured $m_{\tilde{N}_1}$ and
ECAL shower energy

R and ψ can be determined
from ECAL shower reconstruction
if λ not too small



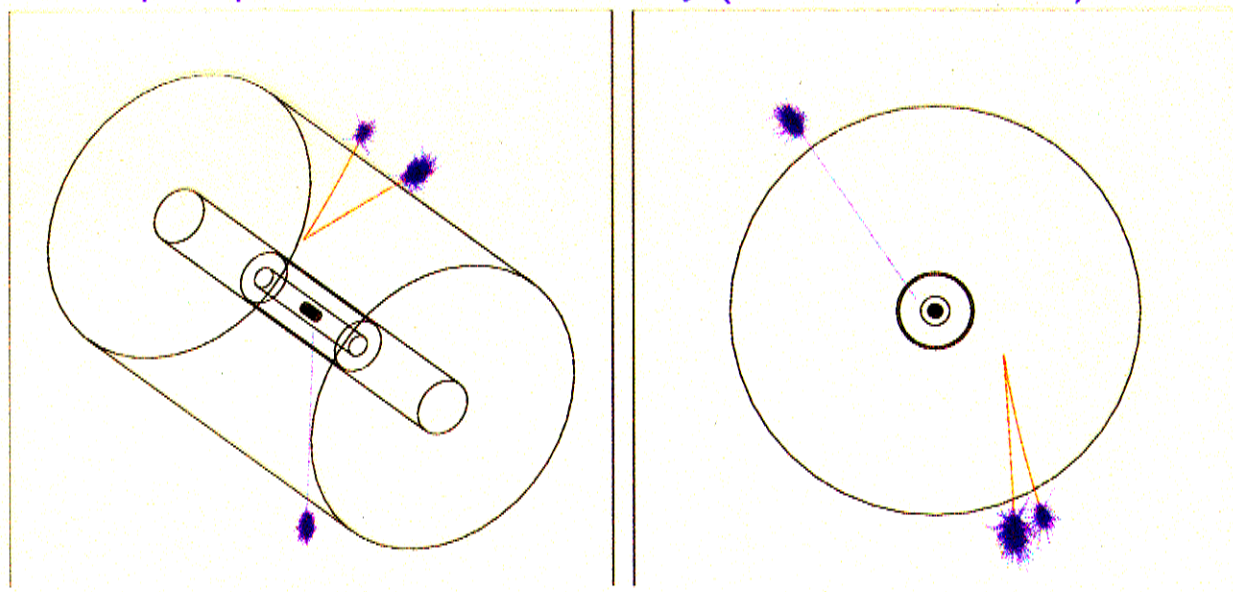
★ (a) Pure photonic \tilde{N}_1 decay

★ (b) γ converts $\rightarrow e^+e^-$ ($D \neq 0$) OR

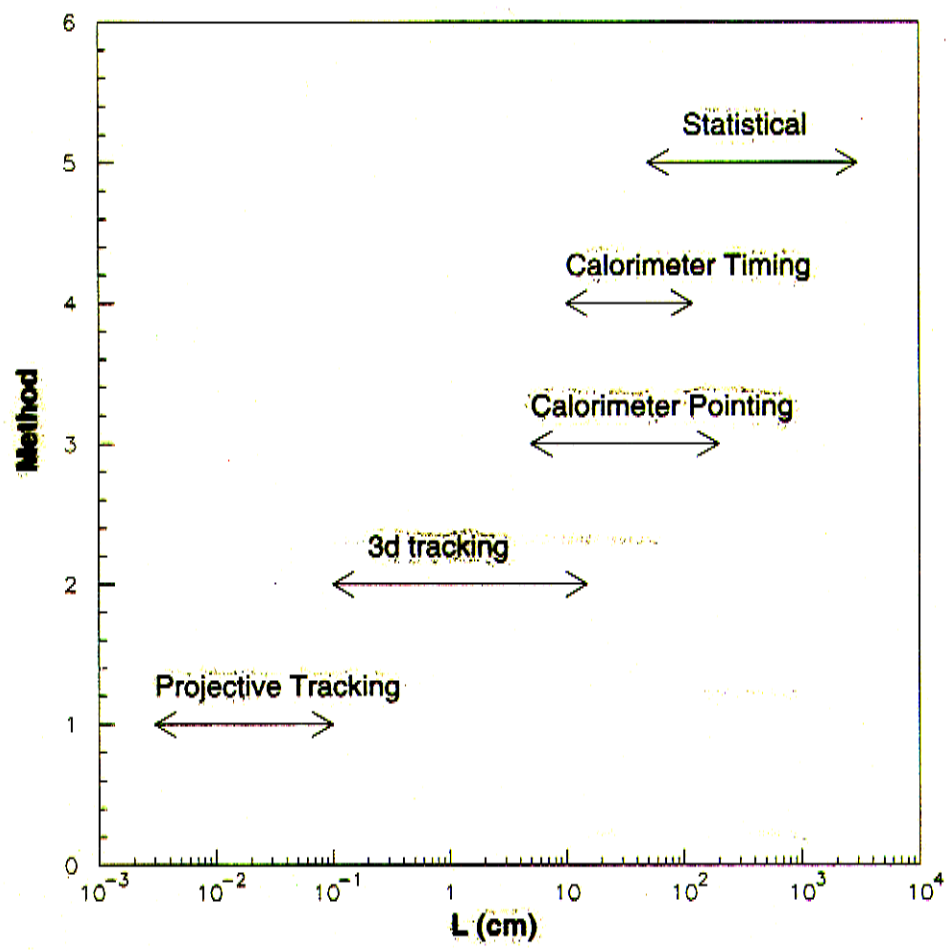
$\tilde{N}_1 \rightarrow 3$ -body "charged decay" ($D = 0$)

\Rightarrow reconstructed pair's vertex/momentum \Rightarrow " γ " flight line

★ A fully simulated $e^+e^- \rightarrow \tilde{N}_1\tilde{N}_1$ event involving both a "charged" and a pure photonic neutralino decay (3D- and end-views):

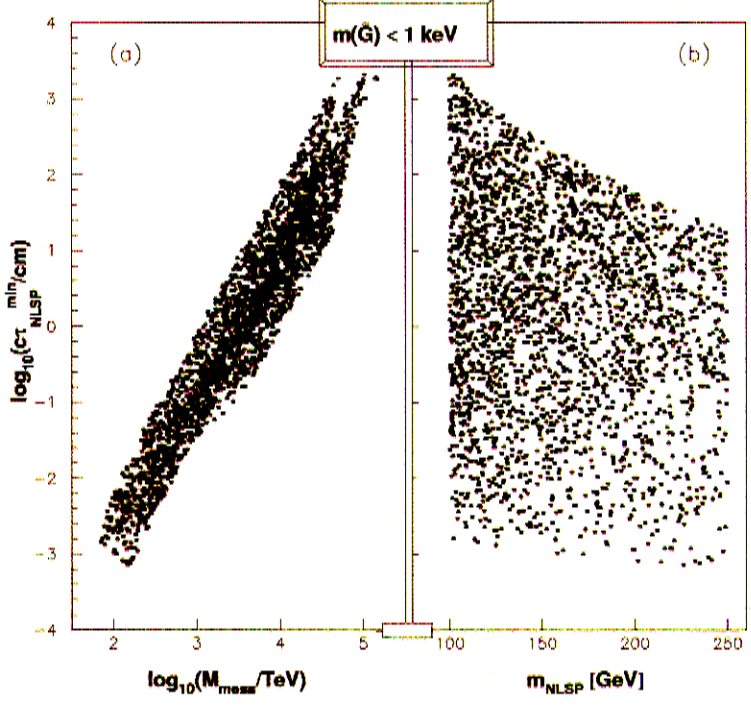


Summary of Techniques

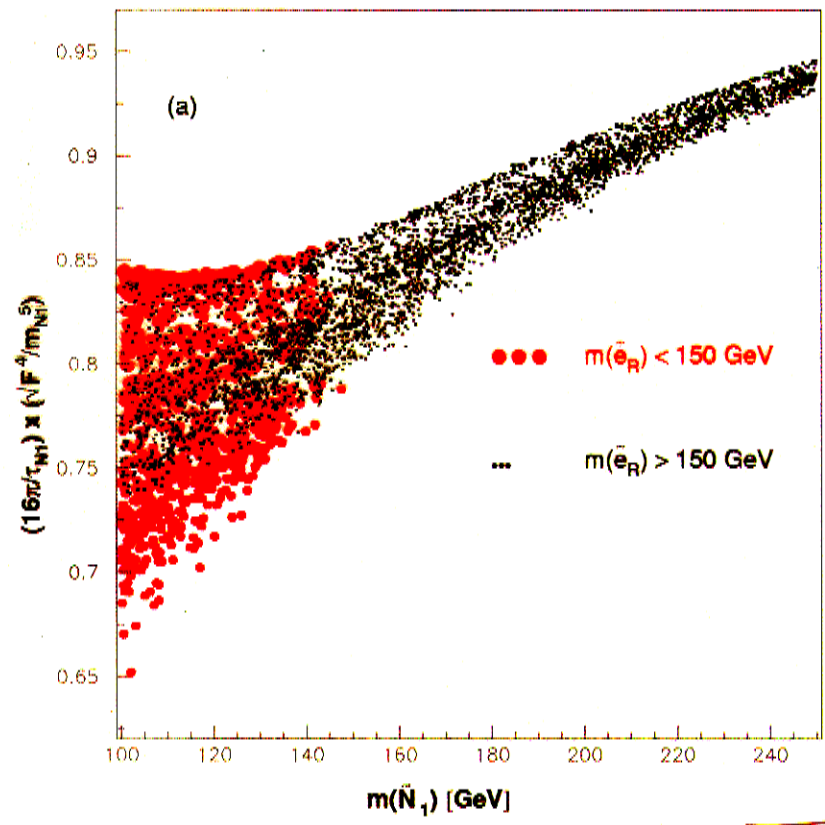


A measurement of L to 10% is possible over the entire range.

GMSB Models with Neutralino NLSP



Neutralino NLSP Lifetime $\rightarrow \sqrt{F}$



Information on $m(\tilde{\chi}_1^0)$, $\tau_{\tilde{\chi}_1^0} \Rightarrow \sqrt{F}$

Super Oblique Corrections

What if square are very heavy?

Can we get a handle on this scale?

'Super' oblique corrections

Nojiri, Pierce, - -
Feng, Polonyi, ...

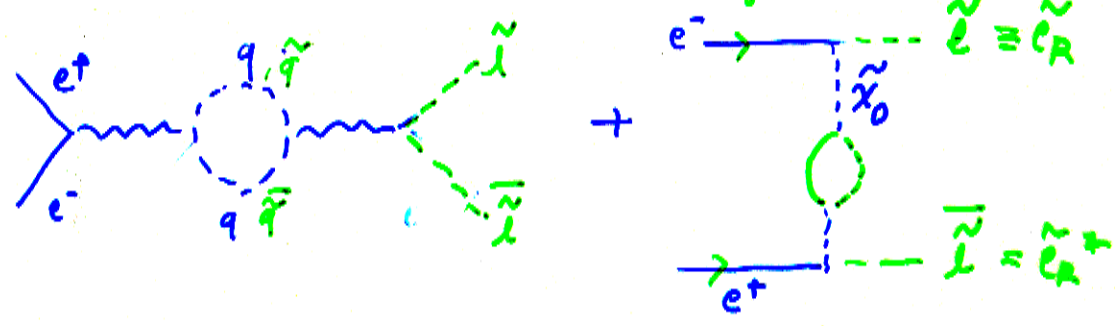
$$g_{e\tilde{e}_R\tilde{B}} = g_e \tilde{e}_R \tilde{B} \equiv g_\gamma \quad \text{SUSY}$$

$$\frac{\delta g_\gamma}{g_\gamma} \approx \frac{11 g_\gamma^2}{48 \pi^2} \ln \left(\frac{m_{\tilde{Q}}}{m_{\tilde{L}}} \right)$$

* splitting bet. $m_{\tilde{Q}}$
and $m_{\tilde{L}} \Rightarrow$ corrections.

SUSY

* corrections even grow with $m_{\tilde{Q}}$!!



$$\frac{m_{\tilde{Q}}}{m_{\tilde{L}}} = 10 \Rightarrow \frac{\delta g_\gamma}{g_\gamma} \sim 0.7\%$$

Feng

$$e^- e^- \rightarrow \tilde{e}_R \tilde{e}_R$$

only the 't' channel diagram exists

Measurement of 'super' oblique correction

to $\sim 0.15\%$ (using only statistics).

Conclusions :

To get precision information about
susy, ~~susy~~ scale, ~~susy~~ mechanism
a high, luminosity e^+e^- collider
 $\sqrt{s} \gtrsim 500$ GeV, upto 1000 GeV
will be absolutely essential.

Such a collider can yield model indep.,
unambiguous information.

'CAN perhaps go a long way towards
putting susy in text books !!'