

Measuring the Higgs CP property  
at a Photon Linear Collider

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# 1. Introduction

Extension of SM in Higgs sector

⇓  
CP-odd Higgs boson(s)

and,

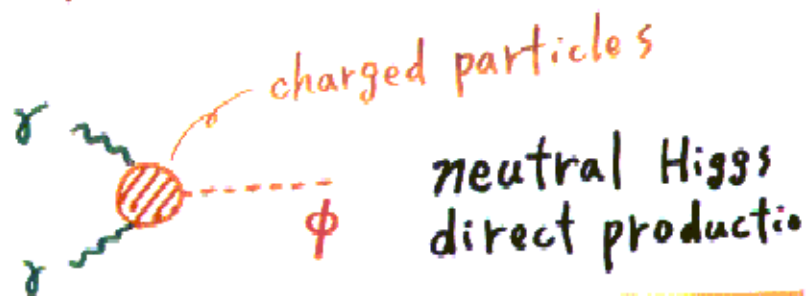
If there exists ~~CP~~ in the Higgs potential,

Mass eigenstates of Higgs bosons  
do NOT carry definite CP-parity.

Measuring CP property of Higgs bosons  
is one of important subjects.

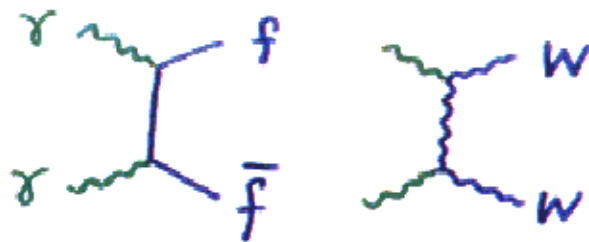
Photon colliders  
give good opportunities.

$\gamma\gamma$  collisions can produce  $J_z = 0$  resonances



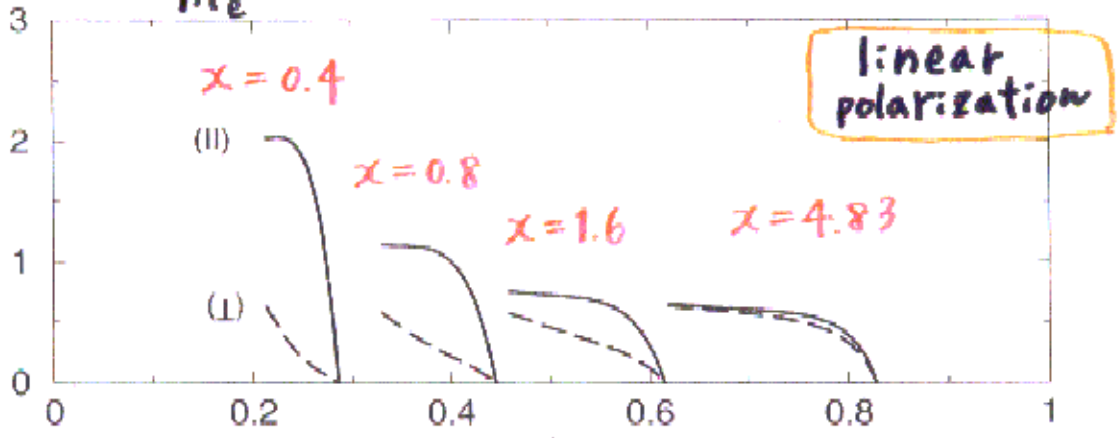
To probe CP property of Higgs bosons,  
we here consider

- initial beam polarization  $\begin{cases} \text{linear} \\ \text{circular} \end{cases}$
- interference effects with background processes
- helicity measurement  
(observation of angular correlation)  
for final particles

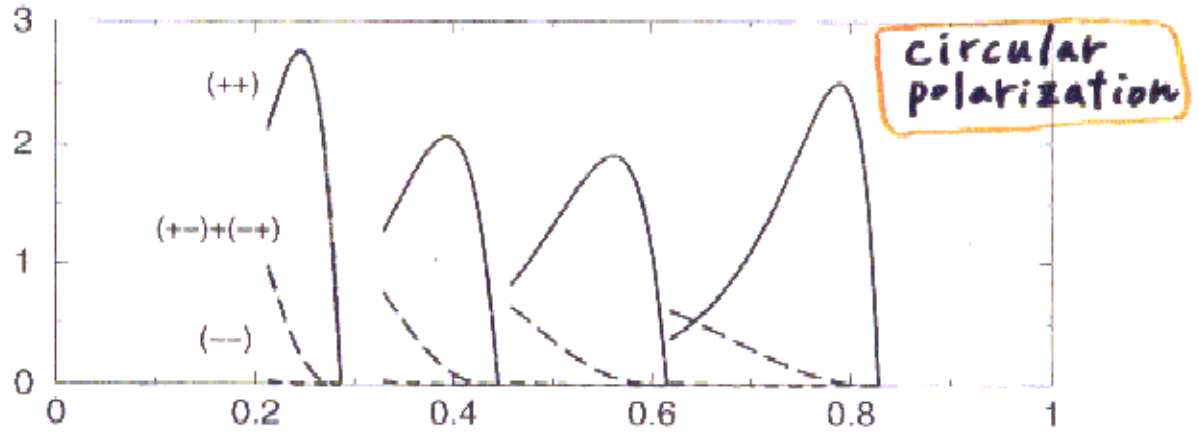


$$\alpha = \frac{4E_e \omega_L}{m_e^2 c^4} \quad \text{initial laser photon energy}$$

$$\frac{1}{L} \frac{dL}{d\sqrt{\tau}}$$



$$\sqrt{\tau} = \sqrt{S_{rr}} / \sqrt{S_{ee}}$$

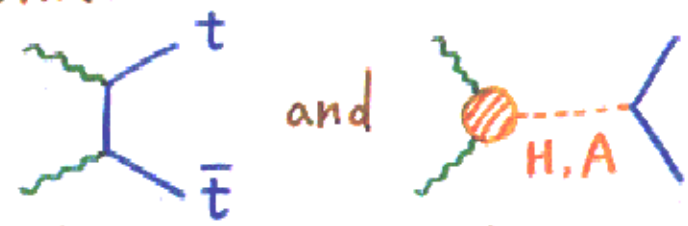


in high  $\sqrt{\tau}$  region  
 linear pol.  $\triangle$   
 circular pol.  $\circ$

## 2. CP invariant case

• E.A., Kamohita, Sugamoto, Watanabe  
EPJC 14 (2000)  
• E.A., Hagiwara

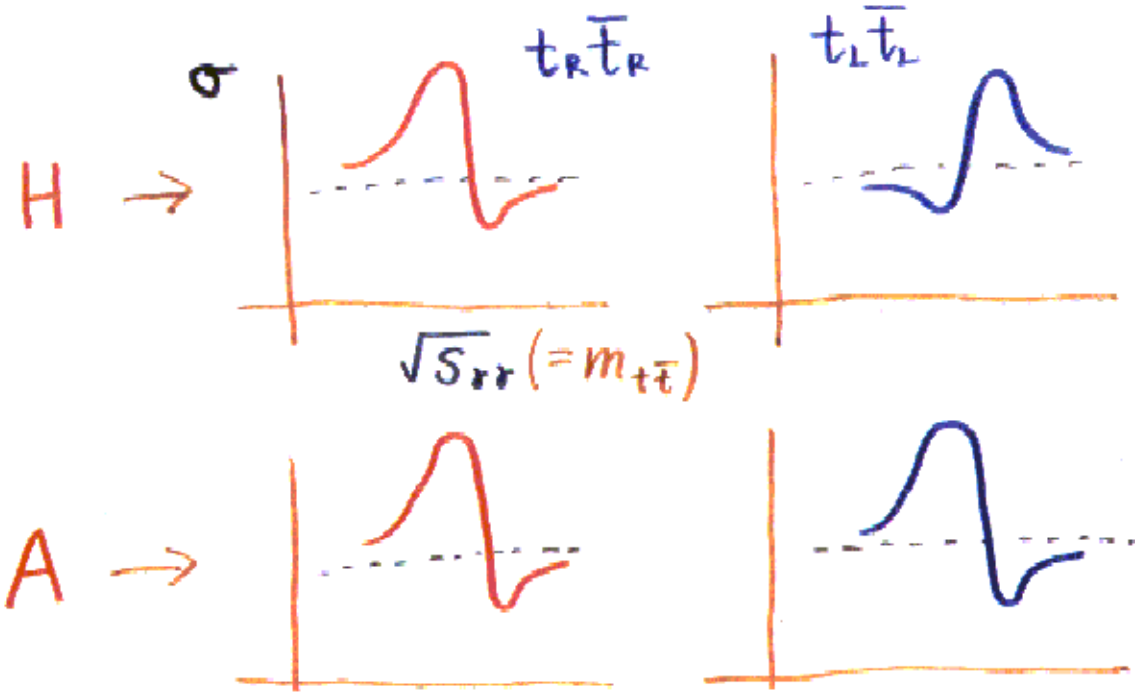
- circular polarization
- interference between



- helicity measurement of top quarks  
(observing angular correlation of decay products)

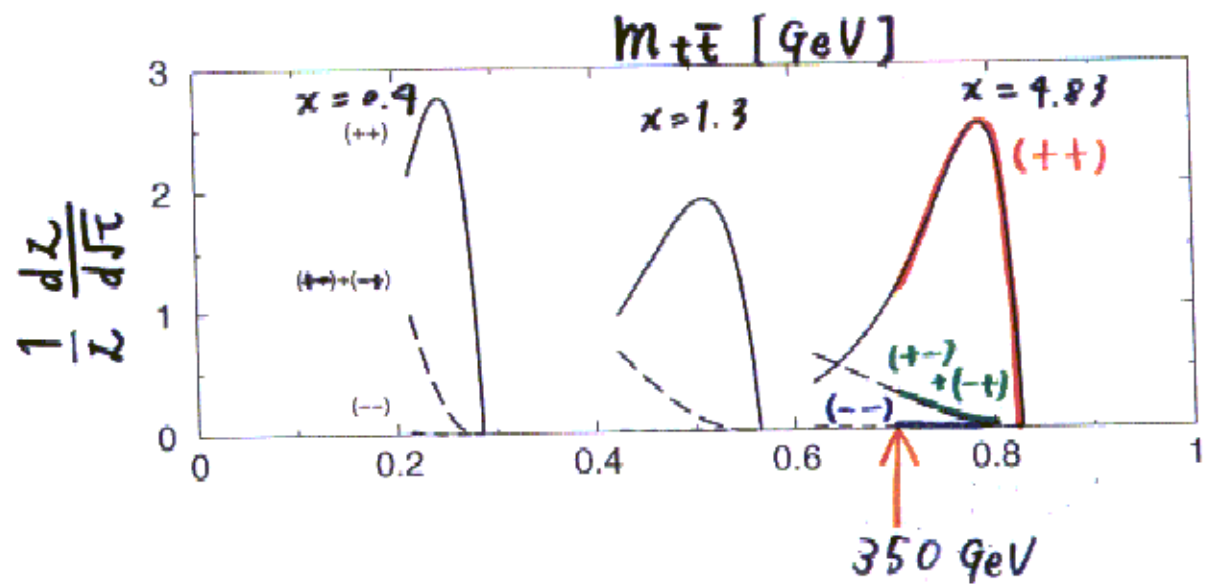
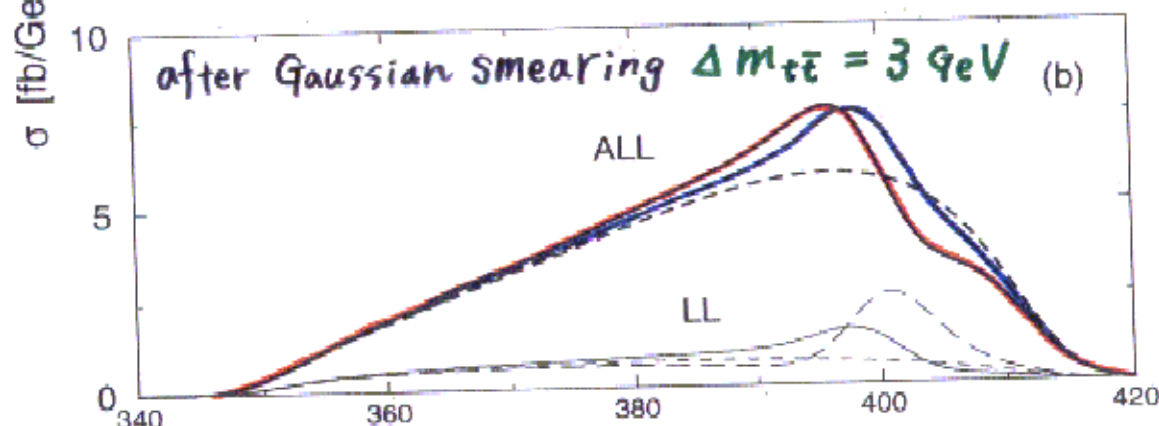
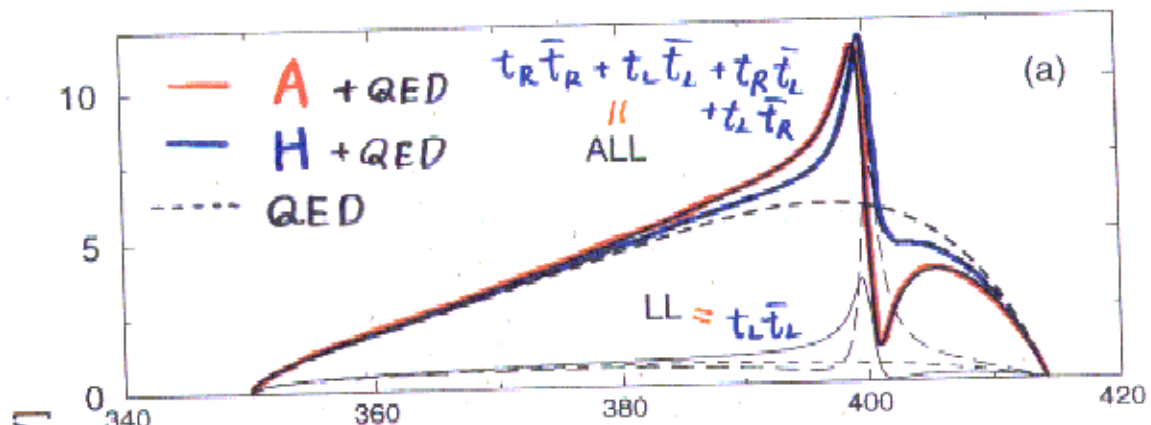
	$t_R \bar{t}_R$	$t_L \bar{t}_L$
$\gamma_+ \gamma_+$	$M_H$	$-M_H$
	$M_A$	$M_A$

$M_H, M_A$ :  
helicity amplitudes  
for  $\gamma_+ \gamma_+ \rightarrow H, A$   
 $\rightarrow t_R \bar{t}_R$



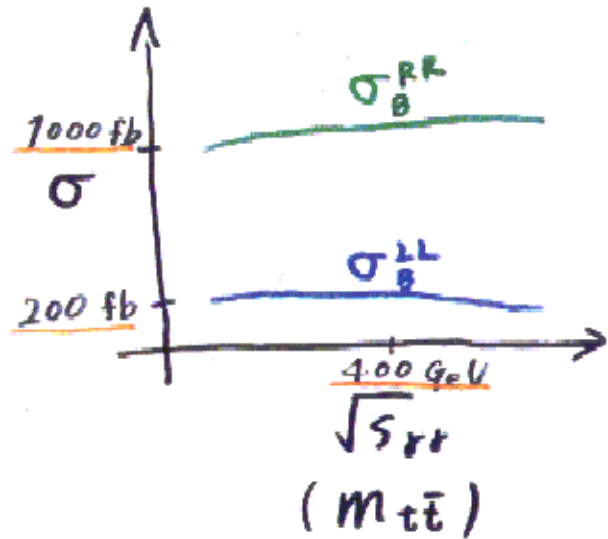
$m_\phi = 400 \text{ GeV}$

$\Gamma_\phi = 1.75 \text{ GeV}$



$$M_{\theta}^{RR} = -8\pi\alpha Q_t^2 \frac{m_t (1 + \beta_t)}{E_t (1 - \beta_t^2 \cos^2 \theta)}$$

$$M_{\theta}^{LL} = -8\pi\alpha Q_t^2 \frac{m_t (1 - \beta_t)}{E_t (1 - \beta_t^2 \cos^2 \theta)}$$



• another observable

$\langle \sin(\bar{\phi} - \phi) \rangle$  is sensitive to CP-parity.

$$\langle \sin(\bar{\phi} - \phi) \rangle \propto \text{Im} \left[ \left( M_B^{RR} + M_\phi^{RR} \right) \left( M_B^{LL} + M_\phi^{LL} \right)^* \right]$$

small

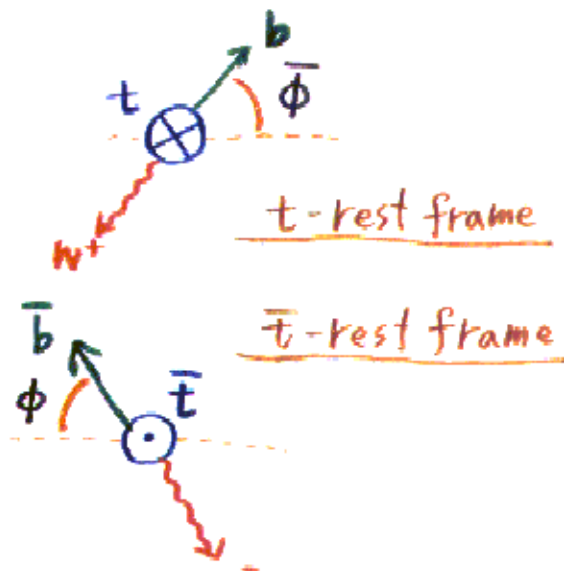
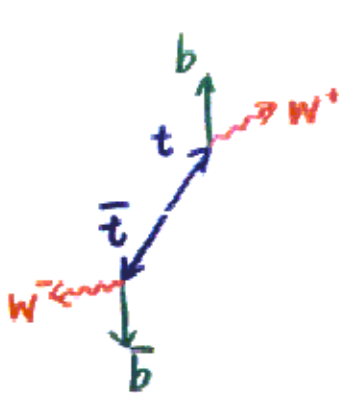
$$= \text{Im} \left[ M_B^{RR} \cdot M_\phi^{LL*} \pm \cancel{\text{real}} \right]$$

for H

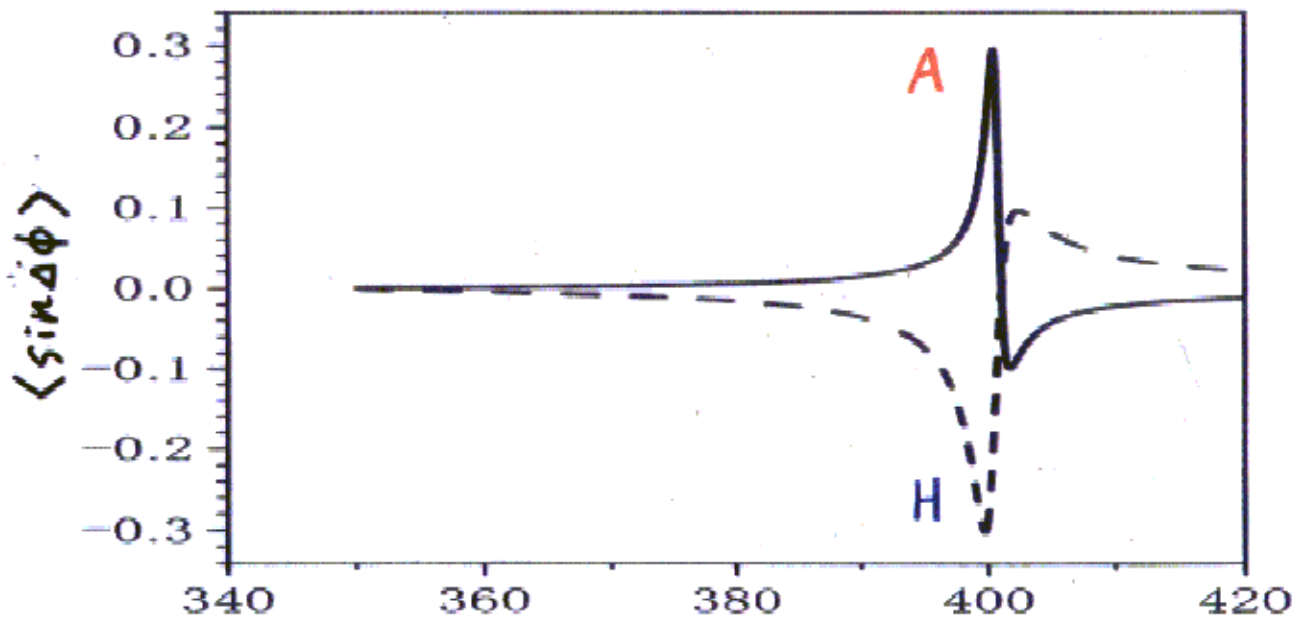
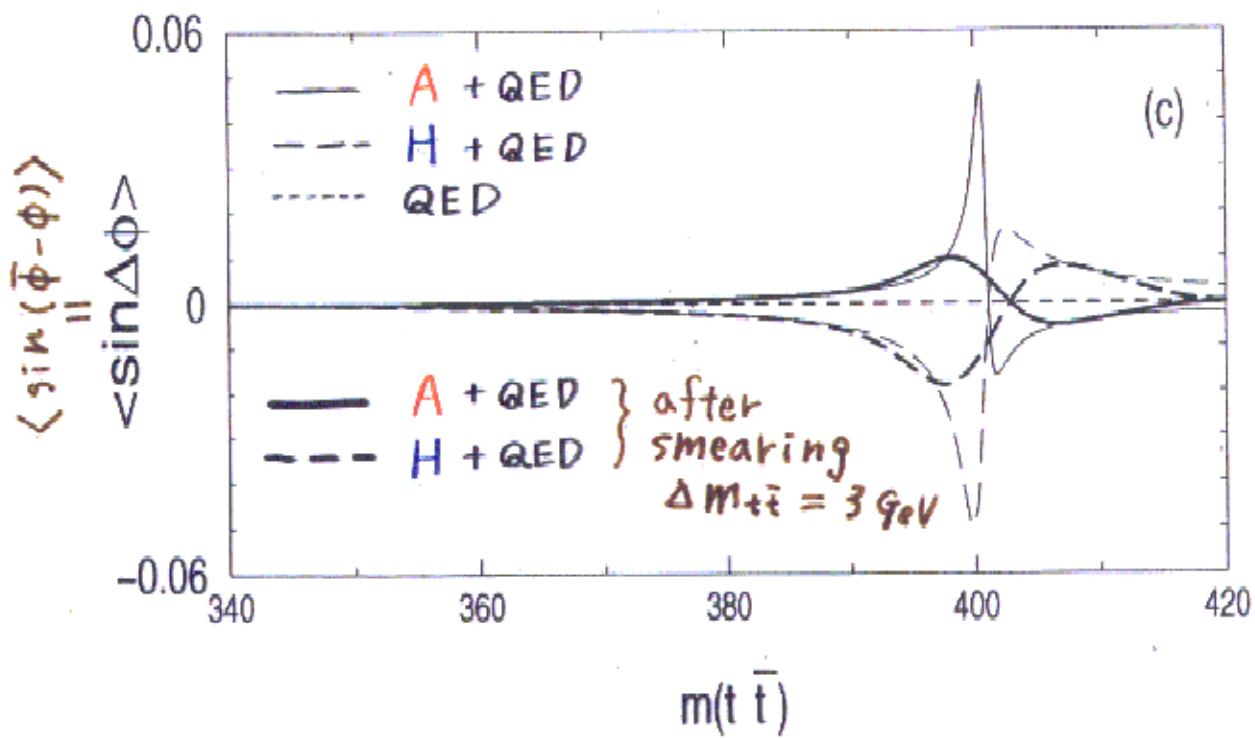
$$\langle \sin(\bar{\phi} - \phi) \rangle \propto -\text{Im} \left[ M_B^{RR} \cdot M_H \right]$$

for A

$$\langle \sin(\bar{\phi} - \phi) \rangle \propto \text{Im} \left[ M_B^{RR} \cdot M_A \right]$$





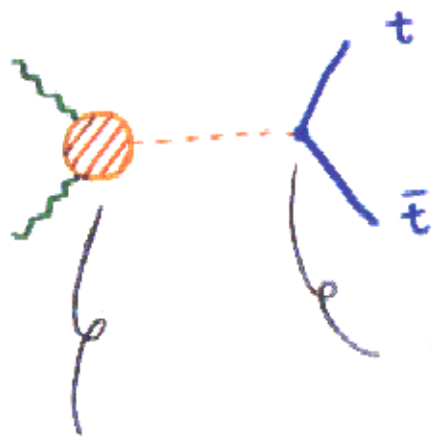


$$m_\phi = 400 \text{ GeV}$$

$$\Gamma_\phi = 1.75 \text{ GeV}$$

### 3. CP non-invariant case

E.A., S.Y. Choi, K. Hagiwara, J.S. Lee  
 hep-ph/0005313  
 (to be published in PRD)



$$\bar{\Psi}_t (S_t + i \gamma_5 P_t) \Psi_t \phi$$

$$S_r \phi F^{\mu\nu} F_{\mu\nu} + P_r \phi F^{\mu\nu} \tilde{F}_{\mu\nu}$$

$S_t, P_t$  : real

$S_r, P_r$  : complex



6 parameters

- circular pol.
- top helicity measurement

photon beam helicities  
t and  $\bar{t}$  helicities

$$\begin{aligned}
 |\mathcal{M}^{(++:++)}|^2 &= \overline{|\mathcal{M}|_0^2} \left[ 1 + \mathcal{A}_0 + \mathcal{A}_1 - (1 + \beta)(\mathcal{A}_2 - \mathcal{A}_3) \right] \\
 |\mathcal{M}^{(--:--)}|^2 &= \overline{|\mathcal{M}|_0^2} \left[ 1 + \mathcal{A}_0 - \mathcal{A}_1 + (1 + \beta)(\mathcal{A}_2 - \mathcal{A}_3) \right] \\
 |\mathcal{M}^{(++:--)}|^2 &= \overline{|\mathcal{M}|_0^2} \left[ 1 - \mathcal{A}_0 + \mathcal{A}_1 + (1 - \beta)(\mathcal{A}_2 + \mathcal{A}_3) \right] \\
 |\mathcal{M}^{(--:++)}|^2 &= \overline{|\mathcal{M}|_0^2} \left[ 1 - \mathcal{A}_0 - \mathcal{A}_1 - (1 - \beta)(\mathcal{A}_2 + \mathcal{A}_3) \right]
 \end{aligned}$$

$$\begin{aligned}
 \overline{|\mathcal{M}|_0^2} &= (1 + \beta^2) A_{\text{cont}}^2 + (\beta^2 S_t^2 + P_t^2) (|S_\gamma|^2 + |P_\gamma|^2) |A_\phi|^2 \\
 &\quad + 2 A_{\text{cont}} \left[ \beta^2 S_t \mathcal{R}(A_\phi S_\gamma) + P_t \mathcal{R}(A_\phi P_\gamma) \right],
 \end{aligned}$$

$$\mathcal{A}_0 = 2\beta A_{\text{cont}} \left\{ A_{\text{cont}} + [S_t \mathcal{R}(A_\phi S_\gamma) + P_t \mathcal{R}(A_\phi P_\gamma)] \right\} / \overline{|\mathcal{M}|_0^2},$$

$$\mathcal{A}_1 = 2|A_\phi|^2 \left\{ (\beta^2 S_t^2 + P_t^2) \mathcal{I}(S_\gamma P_\gamma^*) \right\} / \overline{|\mathcal{M}|_0^2},$$

$$\mathcal{A}_2 = 2\beta A_{\text{cont}} \left\{ S_t \mathcal{I}(A_\phi P_\gamma) \right\} / \overline{|\mathcal{M}|_0^2},$$

$$\mathcal{A}_3 = 2A_{\text{cont}} \left\{ P_t \mathcal{I}(A_\phi S_\gamma) \right\} / \overline{|\mathcal{M}|_0^2},$$

$$A_{\text{cont}} = \frac{16\pi\alpha Q_t^2 m_t}{\sqrt{s} (1 - \beta^2 \cos^2 \theta)}$$

$$A_\phi = \frac{e\alpha}{4\pi} \frac{m_t}{m_w} \frac{s}{s - m_\phi^2 + i m_\phi \Gamma_\phi}$$

Circular + linear pol.

$$|M|^2 = |M|_0^2 \left\{ (1 + \zeta_2 \bar{\zeta}_2) + B_1 (\zeta_2 + \bar{\zeta}_2) + B_2 (\zeta_1 \bar{\zeta}_3 + \zeta_3 \bar{\zeta}_1) - B_3 (\zeta_1 \bar{\zeta}_1 - \zeta_3 \bar{\zeta}_3) \right. \\ \left. + \sin^2 \Theta [-C_0 (\zeta_2 \bar{\zeta}_2 - \zeta_3 \bar{\zeta}_3) + C_1 (\zeta_1 + \bar{\zeta}_1) + C_2 (\zeta_3 + \bar{\zeta}_3) \right. \\ \left. + C_3 (\zeta_1 \bar{\zeta}_2 + \zeta_2 \bar{\zeta}_1) + C_4 (\zeta_2 \bar{\zeta}_3 + \zeta_3 \bar{\zeta}_2) \right\},$$

$P_c$ : degree of circular pol.  
 $P_t$ : degree of linear pol.

Stokes parameters  
 $\zeta_1 \propto P_t \sin 2\kappa$      $\zeta_2 \propto P_c$      $\zeta_3 \propto P_t \cos 2\kappa$

$$\equiv 2 \left\{ |A_\phi|^2 (\beta^2 S_t^2 + P_t^2) I(S_t P_t^*) + A_{cont} [-\beta^2 S_t I(A_\phi P_t) + P_t I(A_\phi S_t)] \right\} / |M|_0^2$$

$$B_1 = (A_1 - \beta A_2 + A_3) |M|_0^2 / |M|_0^2$$

$$B_2 = 2 \left\{ |A_\phi|^2 (\beta^2 S_t^2 + P_t^2) \mathcal{R}(S_t P_t^*) + A_{cont} [\beta^2 S_t \mathcal{R}(A_\phi P_t) + P_t \mathcal{R}(A_\phi S_t)] \right\} / |M|_0^2$$

$$B_3 = \left\{ (-1 + \beta^2 - \beta^2 \sin^4 \Theta) A_{cont}^2 + (\beta^2 S_t^2 + P_t^2) (|S_t|^2 - |P_t|^2) |A_\phi|^2 \right. \\ \left. + 2 A_{cont} [\beta^2 S_t \mathcal{R}(A_\phi S_t) - P_t \mathcal{R}(A_\phi P_t)] \right\} / |M|_0^2$$

CP-odd observables

$$C_0 = 2\beta^2 \sin^2 \Theta A_{cont}^2 / |M|_0^2$$

$$C_1 = 2\beta^2 A_{cont} \left\{ S_t \mathcal{R}(A_\phi P_t) \right\} / |M|_0^2$$

$$C_2 = 2\beta^2 A_{cont} \left\{ A_{cont} + S_t \mathcal{R}(A_\phi S_t) \right\} / |M|_0^2$$

$$C_3 = 2\beta^2 A_{cont} \left\{ S_t I(A_\phi S_t) \right\} / |M|_0^2$$

$$C_4 = -2\beta^2 A_{cont} \left\{ S_t I(A_\phi P_t) \right\} / |M|_0^2$$

- circular + linear pol.
- top helicity measurement

$$\Delta = \frac{|\mathcal{M}|^2(\lambda = \bar{\lambda} = +) - |\mathcal{M}|^2(\lambda = \bar{\lambda} = -)}{|\mathcal{M}|_0^2}$$

$$\Delta = \mathcal{D}_1 (1 + \zeta_2 \bar{\zeta}_2) + \mathcal{D}_2 (\zeta_2 + \bar{\zeta}_2) + \mathcal{D}_3 (\zeta_1 \bar{\zeta}_3 + \zeta_3 \bar{\zeta}_1) - \mathcal{D}_4 (\zeta_1 \bar{\zeta}_1 - \zeta_3 \bar{\zeta}_3) \\ + \sin^2 \Theta \left[ \mathcal{E}_1 (\zeta_1 + \bar{\zeta}_1) + \mathcal{E}_2 (\zeta_3 + \bar{\zeta}_3) + \mathcal{E}_3 (\zeta_1 \bar{\zeta}_2 + \zeta_2 \bar{\zeta}_1) + \mathcal{E}_4 (\zeta_2 \bar{\zeta}_3 + \zeta_3 \bar{\zeta}_2) \right]$$

$$\mathcal{D}_1 = 2\beta A_{\text{cont}} \left\{ -S_t \mathcal{I}(A_\phi P_\gamma) + P_t \mathcal{I}(A_\phi S_\gamma) \right\} / |\mathcal{M}|_0^2,$$

$$\mathcal{D}_2 = 2\beta A_{\text{cont}} \left\{ A_{\text{cont}} + [S_t \mathcal{R}(A_\phi S_\gamma) + P_t \mathcal{R}(A_\phi P_\gamma)] \right\} / |\mathcal{M}|_0^2,$$

$$\mathcal{D}_3 = 2\beta A_{\text{cont}} \left\{ -S_t \mathcal{I}(A_\phi S_\gamma) + P_t \mathcal{I}(A_\phi P_\gamma) \right\} / |\mathcal{M}|_0^2,$$

$$\mathcal{D}_4 = 2\beta A_{\text{cont}} \left\{ S_t \mathcal{I}(A_\phi P_\gamma) + P_t \mathcal{I}(A_\phi S_\gamma) \right\} / |\mathcal{M}|_0^2.$$

$$\mathcal{E}_1 = 2\beta A_{\text{cont}} \left\{ P_t \mathcal{I}(A_\phi P_\gamma) \right\} / |\mathcal{M}|_0^2,$$

$$\mathcal{E}_2 = 2\beta A_{\text{cont}} \left\{ P_t \mathcal{I}(A_\phi S_\gamma) \right\} / |\mathcal{M}|_0^2,$$

$$\mathcal{E}_3 = -2\beta A_{\text{cont}} \left\{ P_t \mathcal{R}(A_\phi S_\gamma) \right\} / |\mathcal{M}|_0^2,$$

$$\mathcal{E}_4 = 2\beta A_{\text{cont}} \left\{ A_{\text{cont}} + P_t \mathcal{R}(A_\phi P_\gamma) \right\} / |\mathcal{M}|_0^2.$$

- circular pol.
- top helicity measurement

$\tan \beta$	$\hat{\sigma}_0[+]$	$\bar{A}_0[+]$	$\bar{A}_1[-]$	$\bar{A}_2[-]$	$\bar{A}_3[-]$
3	0.88 pb	0.45	0.13	-0.17	0.26
10	0.62 pb	0.91	0.00	-0.02	0.03

measurable  
CP-odd obs.  
 $\Rightarrow (-A_2 + \beta_t A_3)$   
 $(A_1 - \beta_t A_2 + A_3)$

- circular + linear pol.

$\tan \beta$	$\bar{B}_1[-]$	$\bar{B}_2[-]$	$\bar{B}_3[+]$	$\bar{C}_0[+]$	$\bar{C}_1[-]$	$\bar{C}_2[+]$	$\bar{C}_3[+]$	$\bar{C}_4[-]$
3	0.46	-0.27	-0.60	0.17	0.13	0.09	0.17	0.06
10	0.03	0.00	-0.47	0.24	0.01	0.30	0.04	0.00

- circular + linear pol.
- top helicity measurement

$\tan \beta$	$\bar{D}_1[-]$	$\bar{D}_2[+]$	$\bar{D}_3[+]$	$\bar{D}_4[-]$	$\bar{E}_1[+]$	$\bar{E}_2[-]$	$\bar{E}_3[-]$	$\bar{E}_4[+]$
3	0.32	0.41	-0.45	0.03	-0.04	0.10	0.09	0.40
10	0.03	0.80	-0.09	0.00	-0.00	0.01	0.01	0.46

$\tan \beta = 3$

$m_\phi = 500 \text{ GeV}$     $\Gamma_\phi = 1.9 \text{ GeV}$

$S_r = -1.3 - 1.2i$     $P_r = -0.57 + 1.1i$

$S_t = 0.33$     $P_t = 0.15$

$|A| = 1 \text{ TeV}$     $\phi = \frac{\pi}{2}$   
 $|\mu| = 2 \text{ TeV}$   
 $m_{\text{susy}} = 0.5 \text{ TeV} \parallel \arg(A_t \mu e^{i\phi})$

$\tan \beta = 10$

$m_\phi = 500 \text{ GeV}$     $\Gamma_\phi = 1.1 \text{ GeV}$

$S_r = -0.39 - 0.35i$     $P_r = -0.06 + 0.14i$

$S_t = 0.11$     $P_t = 0.02$


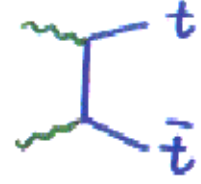
## 4. Summary

We have studied

how to measure Higgs CP property  
in high  $\sqrt{s}$  region.

↳ degree of  
linear polarization  
is not so good.

To observe interference effect

between  and 

can be a powerful method

especially in CP invariant case.

In CP violating case,

linear polarization is also needed  
to determine  $S_r, P_r, S_t, P_t$  completely.  
even without linear polarization  
possible to judge whether  
there exists  $\wedge$  CP or not.  
sizable