

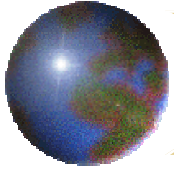
ICFA Seminar

ICFA Seminar on “Future Perspectives in High Energy Physics”

LHC Detector Upgrade - Oct. 9

Dan Green

Fermilab



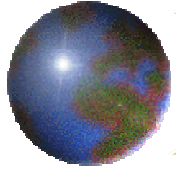
Outline

- **Physics Basics**
 - Z' vs \sqrt{s}
 - Rapidity Range
 - Minbias
 - Pileup and Jets
- **Occupancy and Radiation Dose**
- **Tracker Upgrade**
- **Calorimetry**
- **Muons**
- **Trigger and DAQ**

CERN-TH/2002-078

**“Physics Potential and
Experimental
Challenges of the LHC
Luminosity Upgrade”**

10x will be challenging!



Mass “Reach” and L

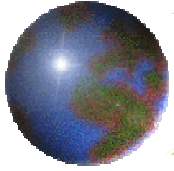
- The number of Z’ detected in leptonic decays is:

$$N_{Z'} = \frac{L \sigma_{Z'}(M)}{4\pi} \times M$$

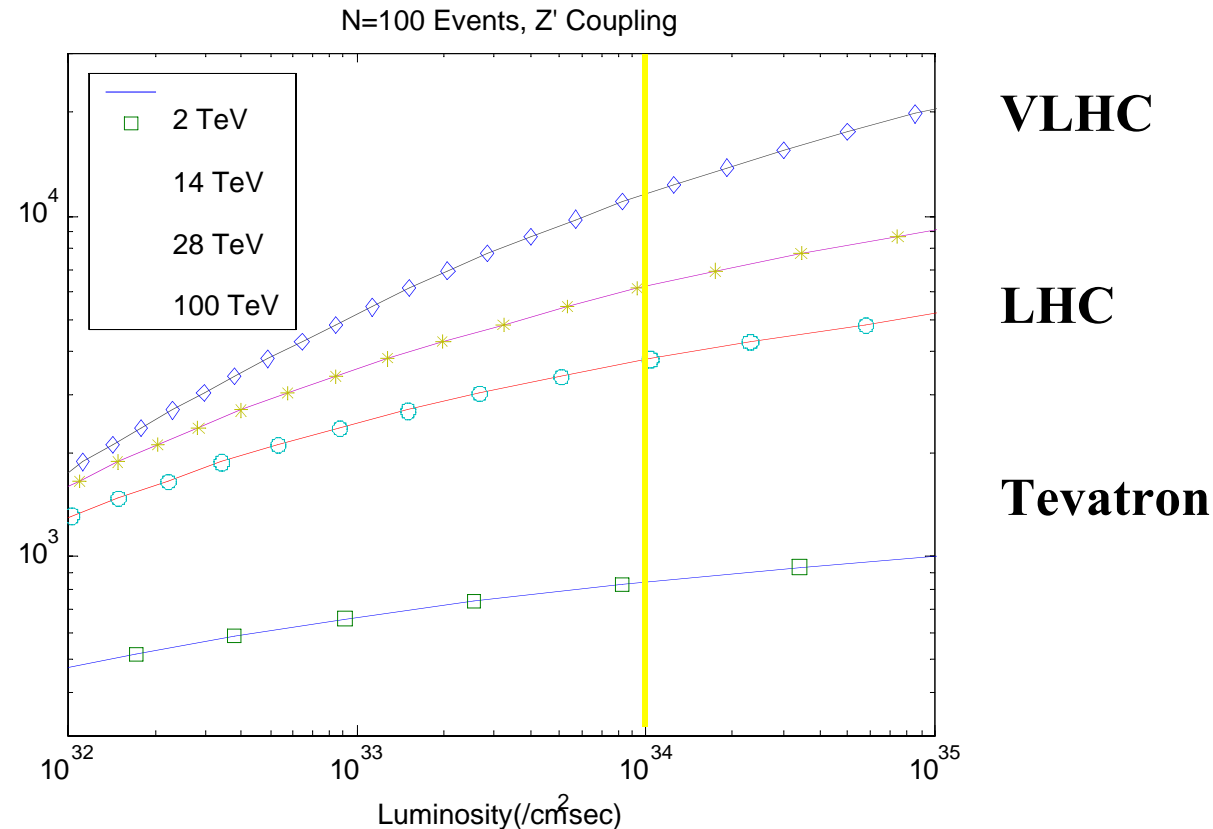
$$= \frac{L \sigma_{Z'}(M)}{4\pi} \times M$$

$$= \frac{L \sigma_{Z'}(M)}{4\pi} \times M$$

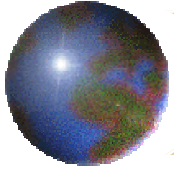
- For $\frac{L \sigma_{Z'}(M)}{4\pi} \times M \sim 10.36(1)$, if $N = 100$ is discovery level then $M \sim 5.3$ TeV is \sim the mass “reach” in 1 year ($M=4 \rightarrow 5.3$ TeV).
- The leptons will be sharply limited to low y or large angles (“barrel”) which helps.



Mass Reach vs L

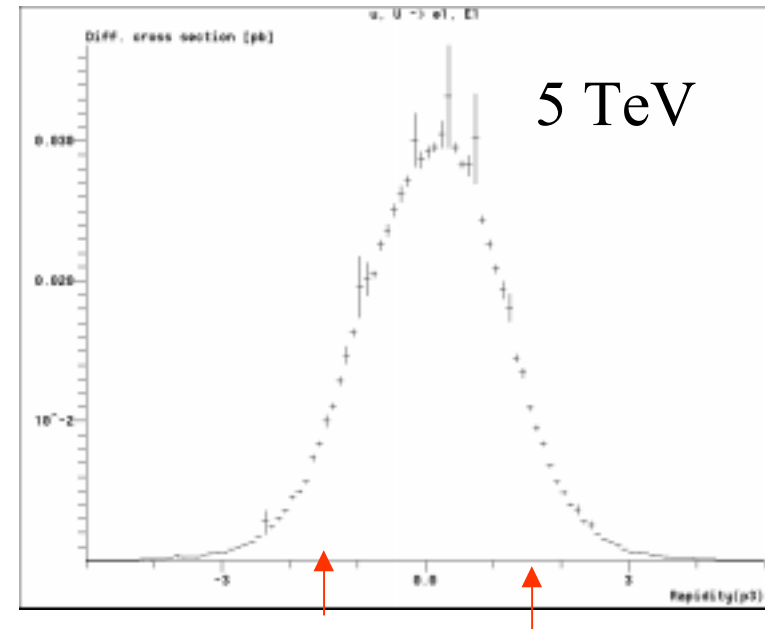
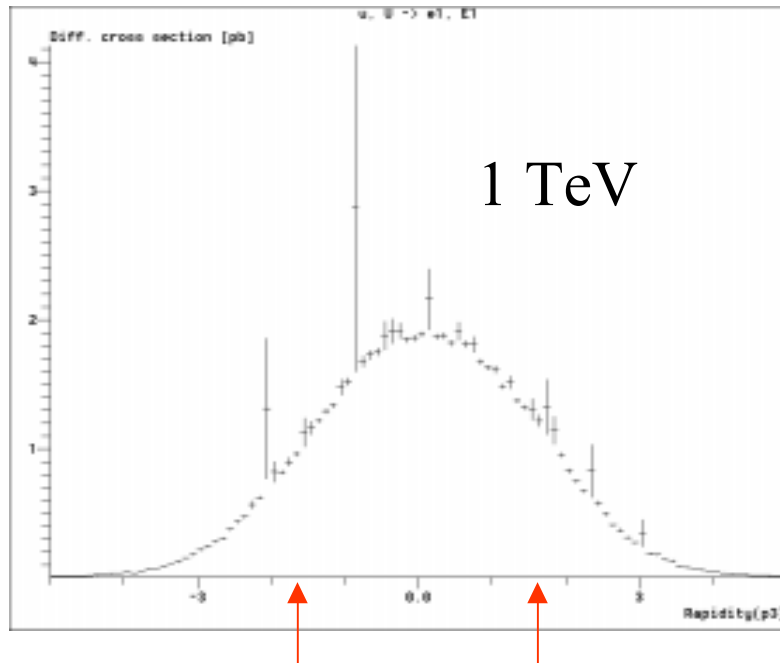


In general mass reach is increased by ~ 1.5 TeV for Z' , heavy SUSY squarks or gluinos or extra dimension mass scales. A $\sim 20\%$ measurement of the HHH coupling is possible for Higgs masses < 200 GeV. However, to realize these improvements we need to maintain the capabilities of the LHC detectors.

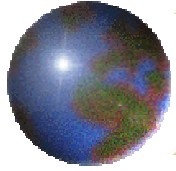


Kinematics

$d\sigma/dy$

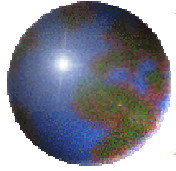


Heavy States decay at wide angles. For example Z' of 1 and 5 TeV decaying into light pairs. Therefore, we need to concentrate on wide angle detectors.



Inclusive Interactions

- The inclusive p - p interaction has an inelastic cross section ~ 50 mb.
- It produces \sim equal numbers of $\pi^+\pi^+\pi^+$ which are distributed \sim uniform in rapidity, y , with a “density” ~ 9 pions per unit of y .
- The pions have a distribution in transverse momentum with a mean, ~ 0.6 GeV.

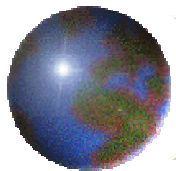


Pileup and Luminosity

- For $\sigma \sim 50$ mb, and $\rho \sim 40$ charged pions/unit of y with a luminosity of 10^{31} cm²/sec and a crossing time of 25 nsec (12.5 nsec ?):

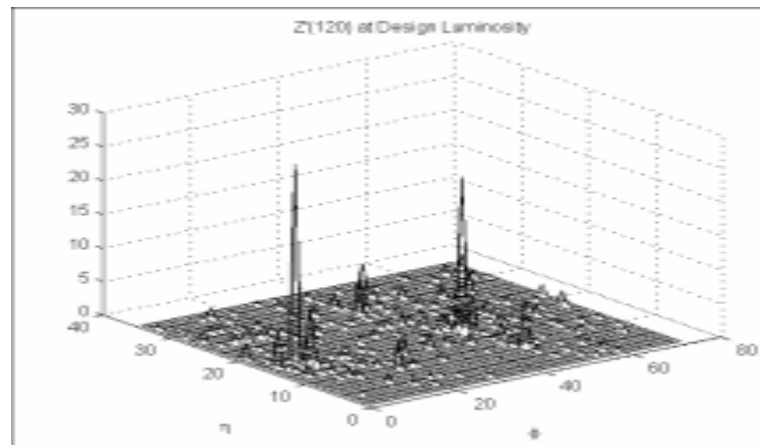
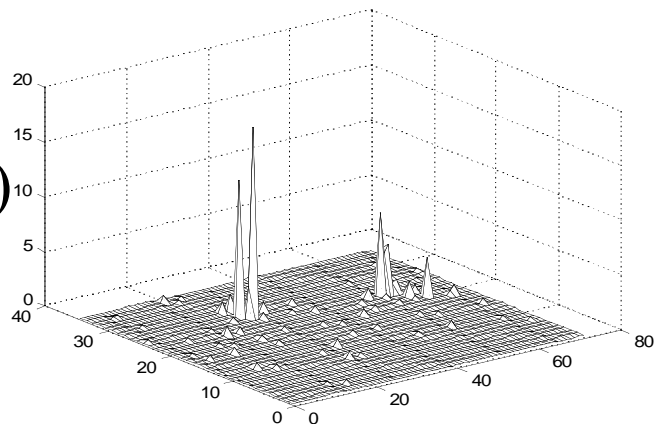
$$\sim 141 \text{ pions/unit of } y$$

- In a cone of radius = 0.5 there are ~ 141 pions, or ~ 84 GeV of transverse momentum. This makes low E_t jet triggering and reconstruction difficult.

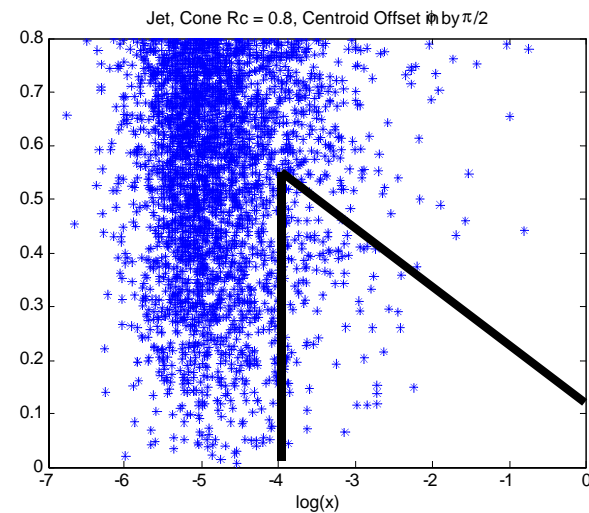
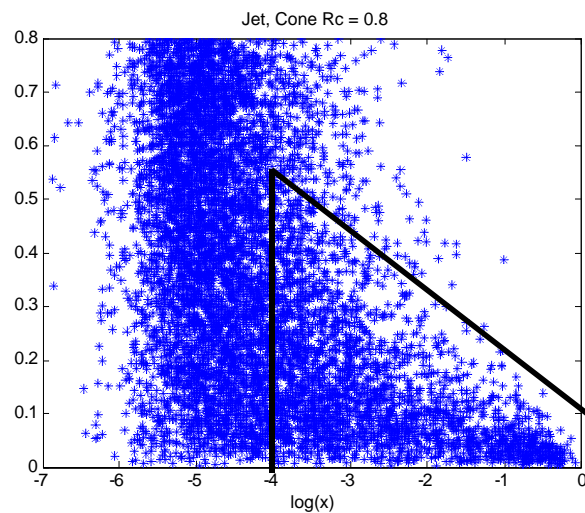


$Z'(120)$ at $L/5$ and L

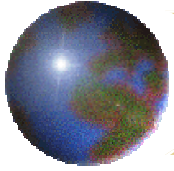
$E_t(\text{GeV})$



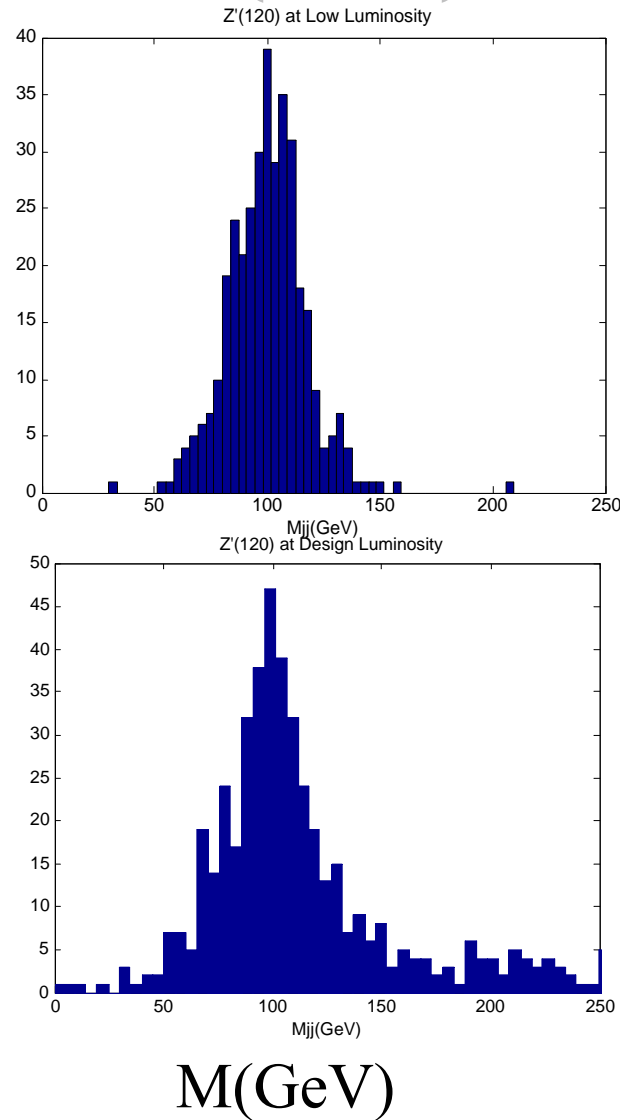
dR



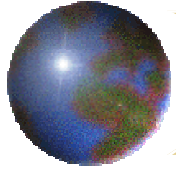
$\text{Log}(x)$



Z'(120) Mass Resolution

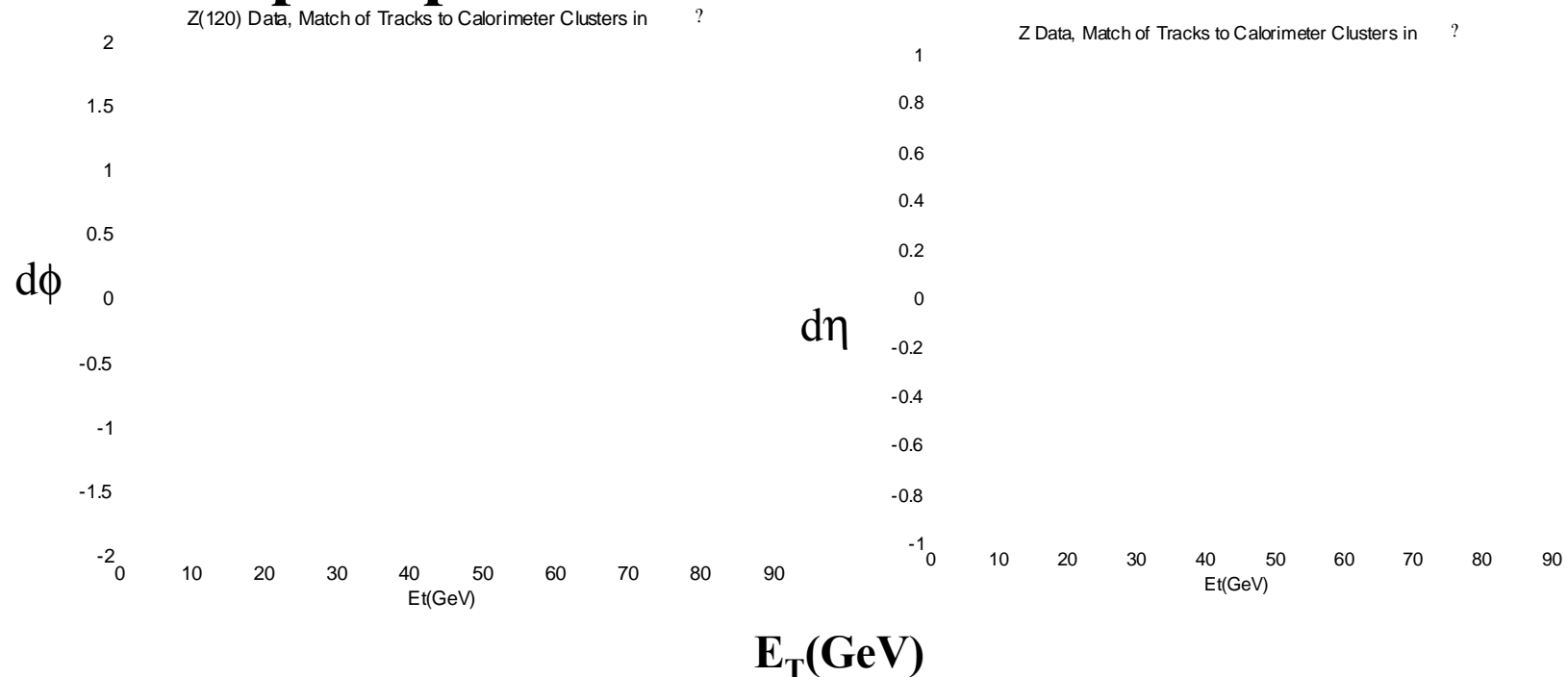


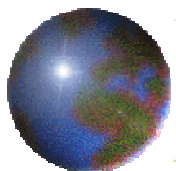
Note that the calorimeter cells are still fairly sparsely populated (granularity (10×8)) at 10^{34} . With the cuts shown, the dM/M with Gaussian fits is the same at $L/5$ and at L . Use the fact that QCD implies that there is a core of the jet at small dR and large x . Extend to $10 \times L$ using tracker and energy flow inside the jet?



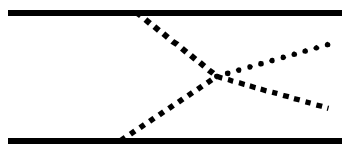
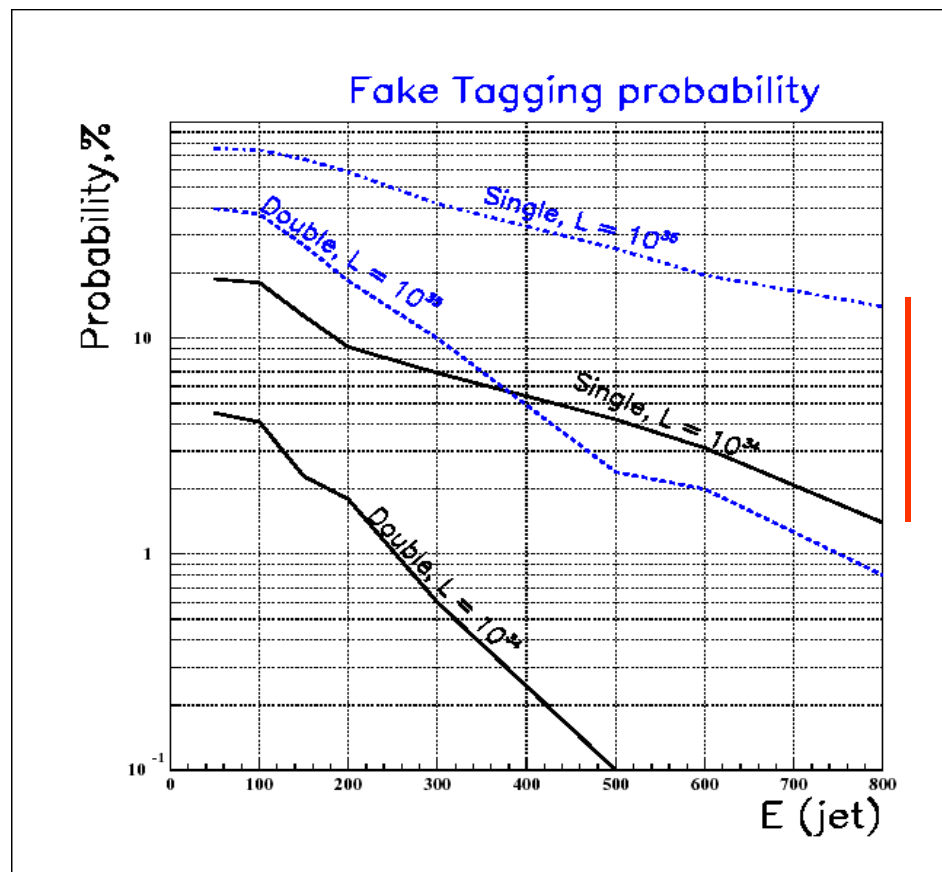
Tracker and Energy Flow

- For 120 GeV Z' match tracks in η and ϕ to “hadronic” clusters within the jet. Improves dijet mass resolution. Units are HCAL tower sizes. Also use track match to remove deposits from pileup vertices ?





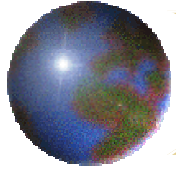
WW Fusion and “Tag Jets”



WW
fusion

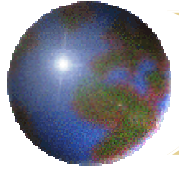
~~MM~~
These jets have

and $\langle y \rangle \sim 3$. Lose
10x in fakes. We
must use the energy
flow inside a jet cone
to further reduce the
fake jets due to
pileup. We can also
use the tracker to
remove charged
energy from pileup
vertices (in principle)



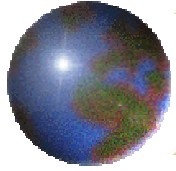
Tracking Detectors

- **Clearly, the tracker is crucial for much of the LHC physics [e.g. e , μ , jets (pileup, E flow)].**
- **The existing trackers will not be capable of utilizing the increased luminosity as they will be near the end of their useful life.**
- **It is necessary to completely rebuild the LHC tracking detectors.**



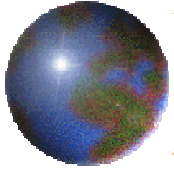
Tracker - Occupancy

- The occupancy, O , for a detector of area dA exposed for time dt at (r,z) is $\frac{dN}{dA dt}$
- e.g. Si strip 10 cm x 100 μm in a 25 nsec crossing at $r = 30$ cm is 1.3 %
- For higher luminosity, decrease dA , or decrease dt or increase r – smaller, faster or further away.



Tracker Occupancy

- **Preserve the performance using $\frac{1}{z}$:**
 - Push Si strips out to ~ 60 cm. – development
 - Push pixels out to 20 cm. – development
 - For $r < 20$ cm. Need new technologies – basic research
- **Shrink dAdt 10x at fixed r to preserve b tagging?**
- **Possibilities**
 - 3-d detectors – electrodes in bulk columns
 - Diamond (RD42) - radhard
 - Cryogenic (RD39) – fast, radhard
 - Monolithic – reduced source capacity.



Monolithic Pixel - DEPFET

Combine
the
detector
and the
readout
for
pixels?

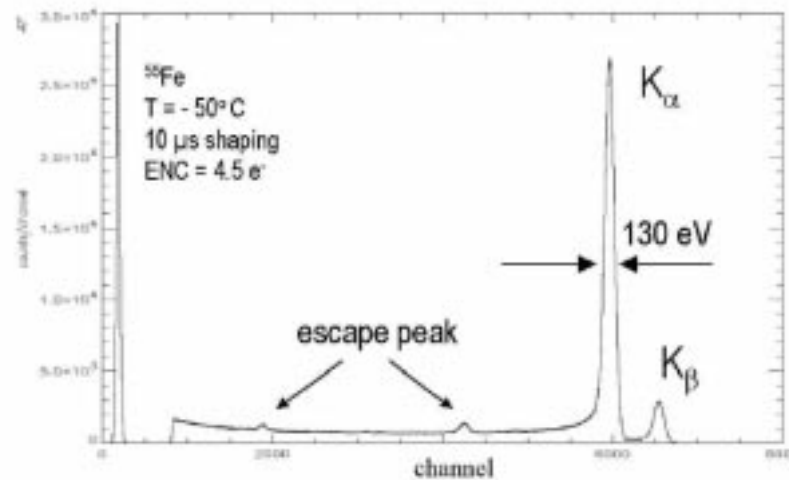
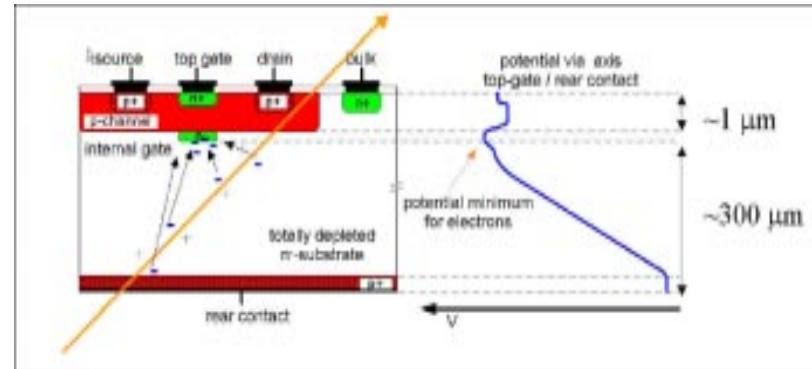
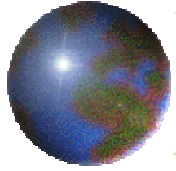
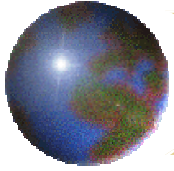


Fig. 10. Response of a DEPFET pixel detector to an ^{55}Fe 6keV X-ray source obtained using a single pixel device operated at -50°C with a shaping time of $10 \mu\text{s}$.

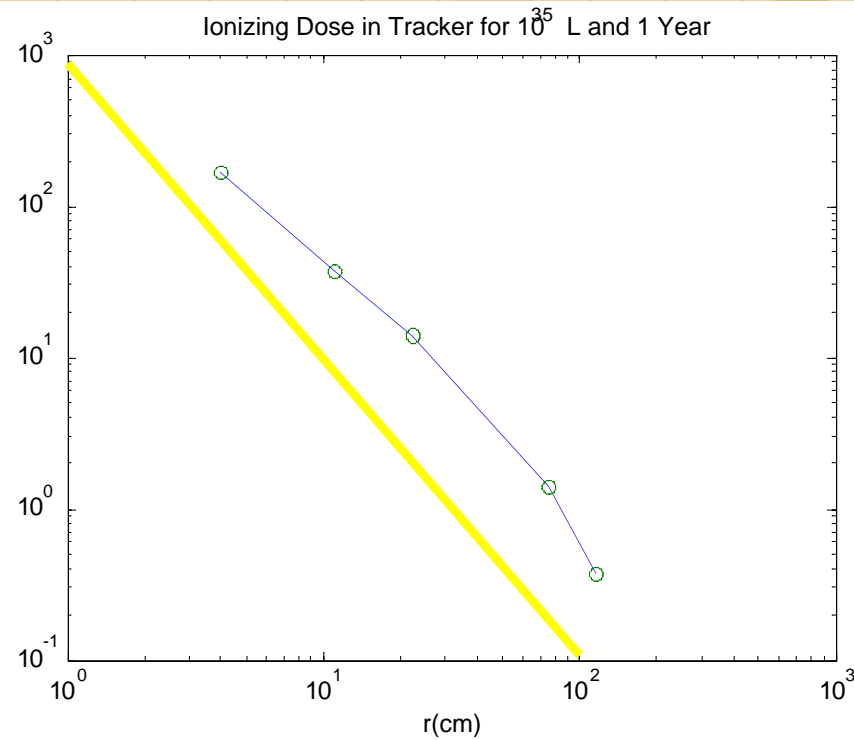


Tracker – Ionizing Dose

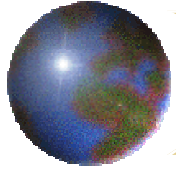
- The ionizing dose due to charged particles is:
$$\int \frac{L}{4\pi r^2} dt$$
- The dose depends only on luminosity, r , and exposure time dt .
- For example, at $r = 30$ cm, the dose is 1.3 Mrad/yr – ignoring loopers, interactions,



Tracker ID vs. Radius



- **DSM is radiation hard in 0.25 um**
- **0.13 um is available – use Moore’s “law”**
- **Development needed – track and adapt commercial processes**
- **Optical links – development needed**



ECAL – Shower Dose

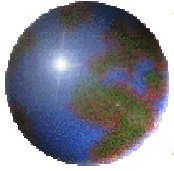
- The dose in ECAL is due to photon showers

and is:

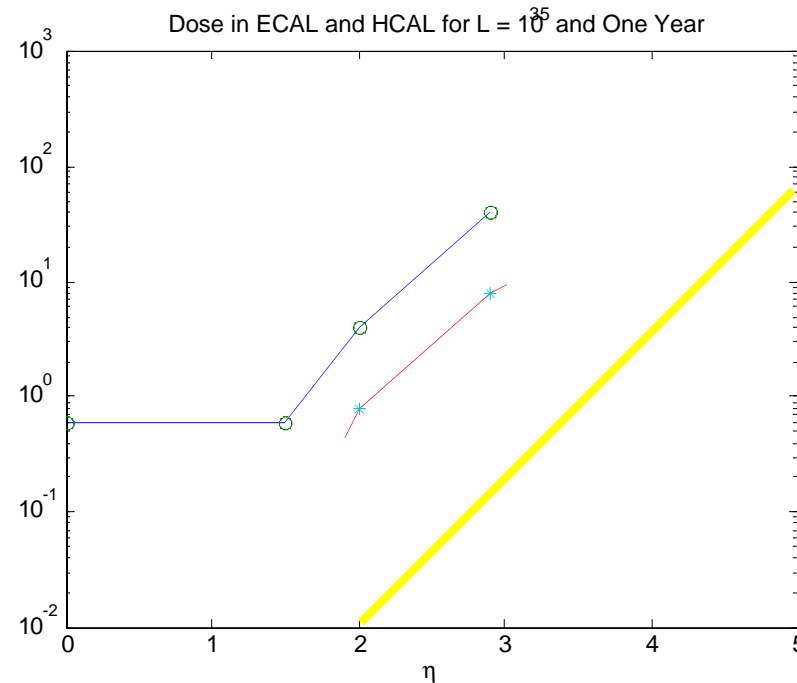
$$\int \frac{dN}{dE dx pE}$$

- In the barrel, SD is $\sim \frac{1}{4} \sin^2 \theta$. In the endcap, SD $\sim \frac{1}{4} \sin^2 \theta$

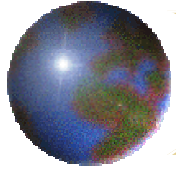
- At $r = 1.2$ m, for Pb with $E_c = 7.4$ MeV, the dose at $y=0$ is 3.3 Mrad/yr, at $|y|=1.5$ it is 7.8 Mrad/yr.



HCAL and ECAL Dose

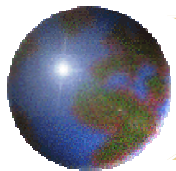


The dose ratio is \sim ~~1~~ Barrel doses are not a problem. For the endcaps a technology change may be needed for $2 < |y| < 3$. For CMS endcap switch to quartz fiber?



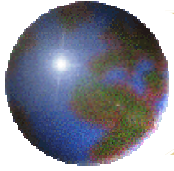
ECAL

- **For both ATLAS and CMS the barrel will probably tolerate the increased dose. There are issues of $\sim 3x$ increased pileup noise and poorer isolation for electrons.**
- **ATLAS LA has space charge and current draw issues. CMS has APD leakage current noise issues in the barrel. The CMS endcap needs development.**

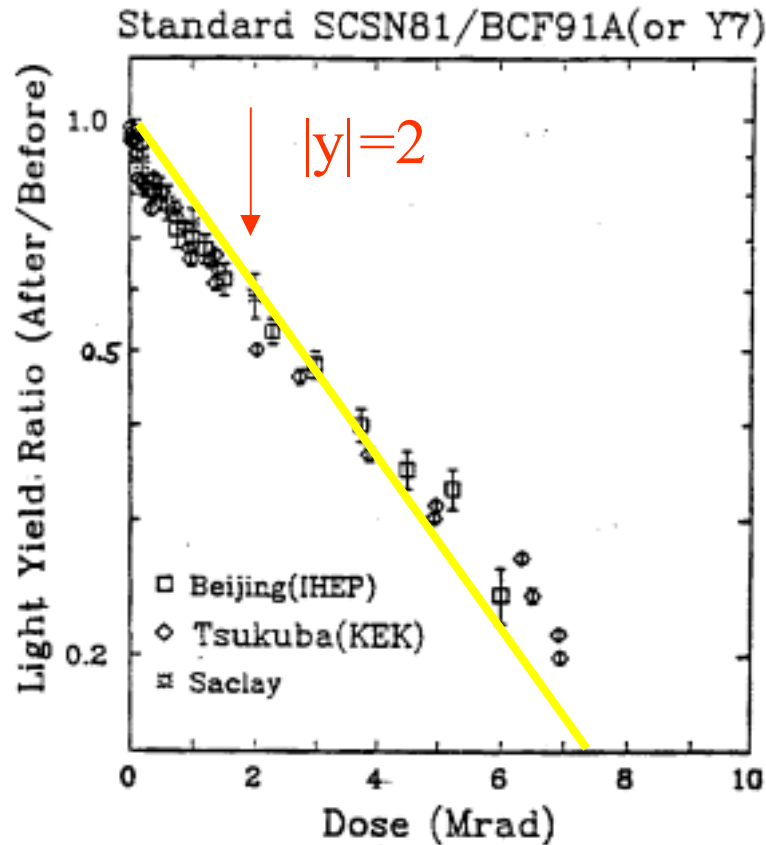


HCAL

- **Both ATLAS and CMS will function in the barrel region.**
- **In the $3 < |y| < 5$ region, a reduction to $y < 4.2$ keeps the dose constant. The loss of efficiency is not terrible (peak “tag” rate at $|y|=3$). Or replace quartz fibers with high pressure gas? Better tower granularity might be needed due to pileup and “fake” jets.**
- **At $|y| \sim 3$ the CMS scintillator needs development – better longitudinal segmentation and/or improved scintillator.**



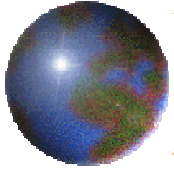
Scintillator - Dose/Damage



Scintillator under irradiation forms Color centers which reduce the Collected light output (transmission loss).

$LY \sim \exp[-D/D_0]$, $D_0 \sim 4 \text{ Mrad}$

This technology will not survive gracefully at $|y| \sim 3$ without development of improved scintillator.



Muons and Shielding

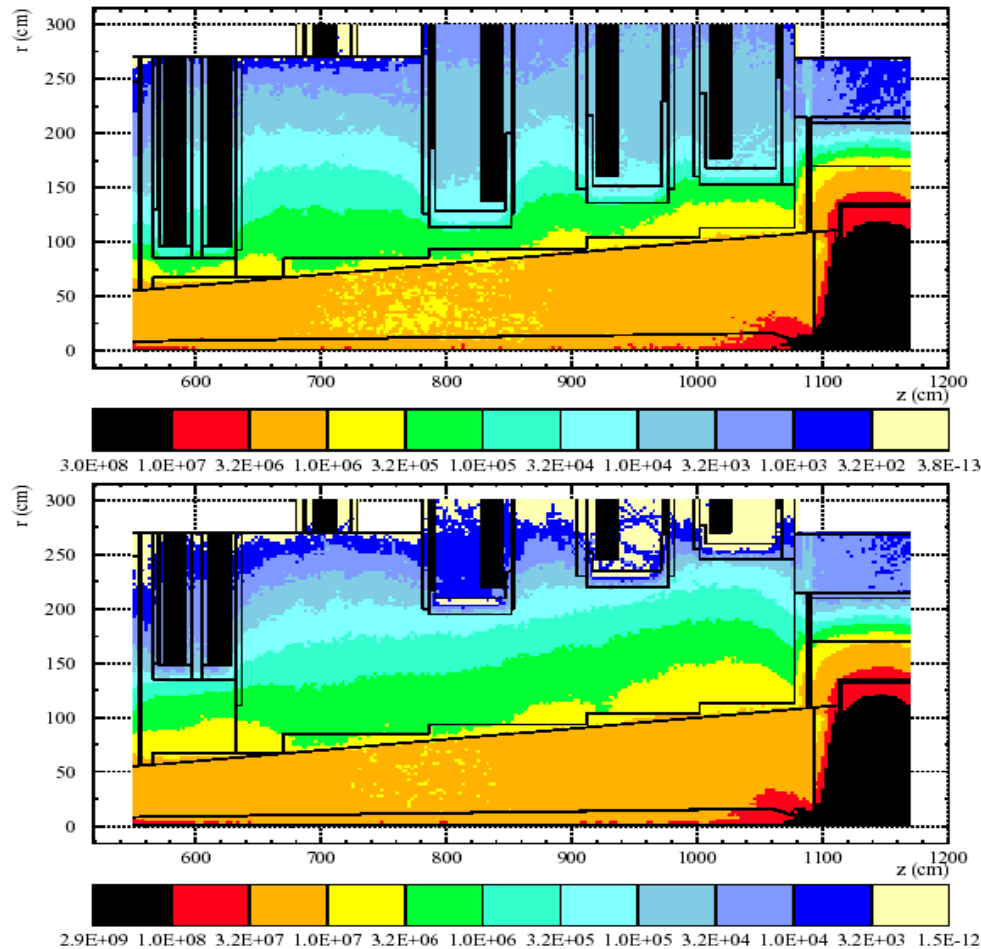
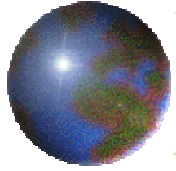


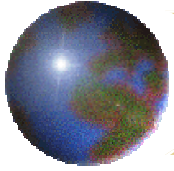
Fig. 24: Upper: the neutron fluxes in the low radius, high η -region of the CMS endcap muon detector for $\eta < 2.4$ and present LHC shielding $10^{34} \text{ cm}^{-2}\text{s}^{-1}$. Lower: same, for $\eta < 2$ and possible shielding for SLHC $10^{35} \text{ cm}^{-2}\text{s}^{-1}$.

With added shielding, dose rates can be kept constant if angular coverage goes from $|y| < 2.4$ to $|y| < 2$. CSC migrate into barrel?



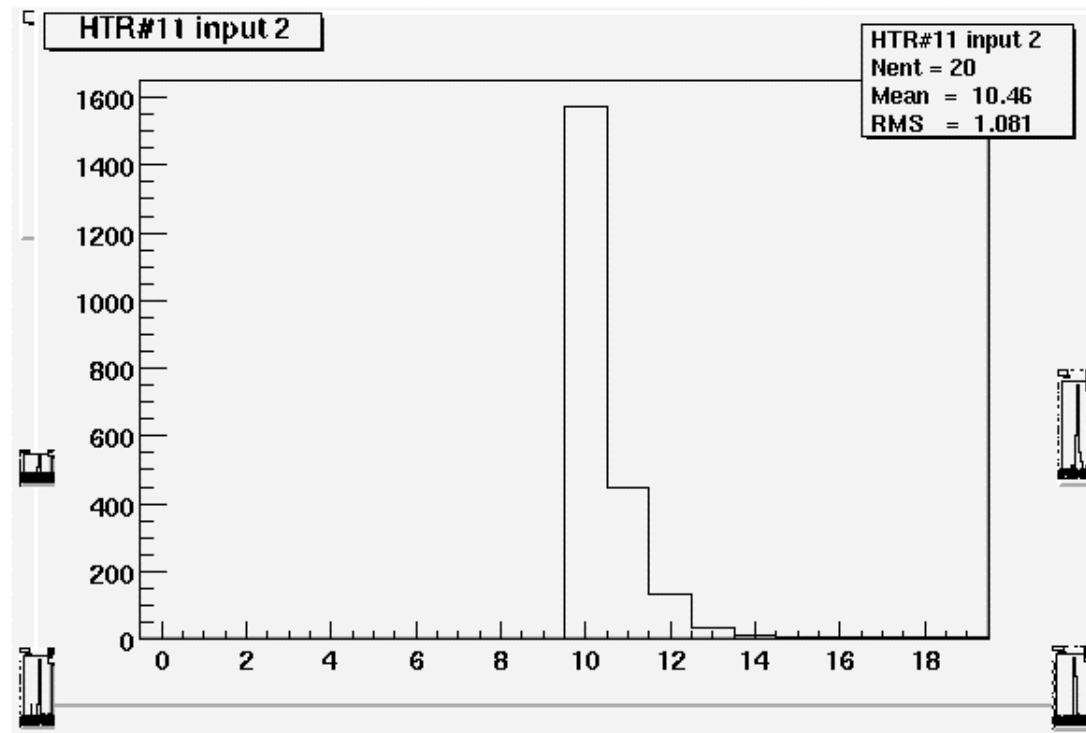
Trigger and DAQ

- **Assuming LHC initial program is successful, raise the trigger thresholds.**
- **Rebuild trigger system to run at 80 MHz. Utilize those detectors which are fast enough to give a BCID within 12.5 nsec.**
- **Examine algorithms to alleviate e.g. degraded e isolation.**
- **Design for the increased event size (pileup) with reduced L1 rate and/or data compression.**
- **For DAQ track the evolution of communication technologies, e.g. 10 Gb/sec Ethernet.**



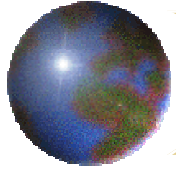
300 GeV Pion – H2 test Beam

E



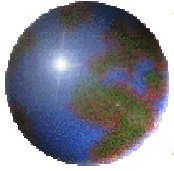
HTR - Bunch crossing
number

The shape of the pulse in time is ~ as expected – due to scint flours and HPD drift. Bunch crossing ID can be extended to 12.5 nsec (80 MHz) ?



Summary

- **The LHC Physics reach will be substantially increased by higher luminosity.**
- **To realize that improvement, the LHC detectors must preserve performance.**
- **The trackers must be rebuilt – with new technology at $r < 20$ cm.**
- **The calorimeters, muon systems, triggers and DAQ will need development.**
- **The upgrades are likely to take $\sim (6-10)$ years. The time to start is now, and the people to do the job are those who did it for the present detectors.**



HH Study

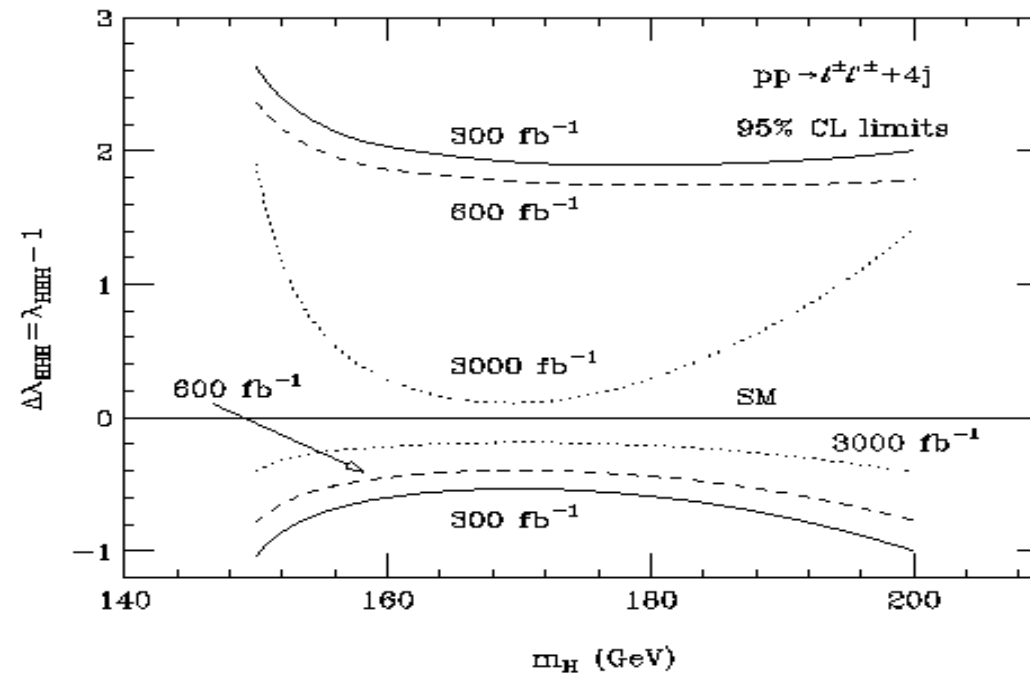
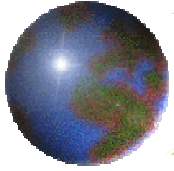
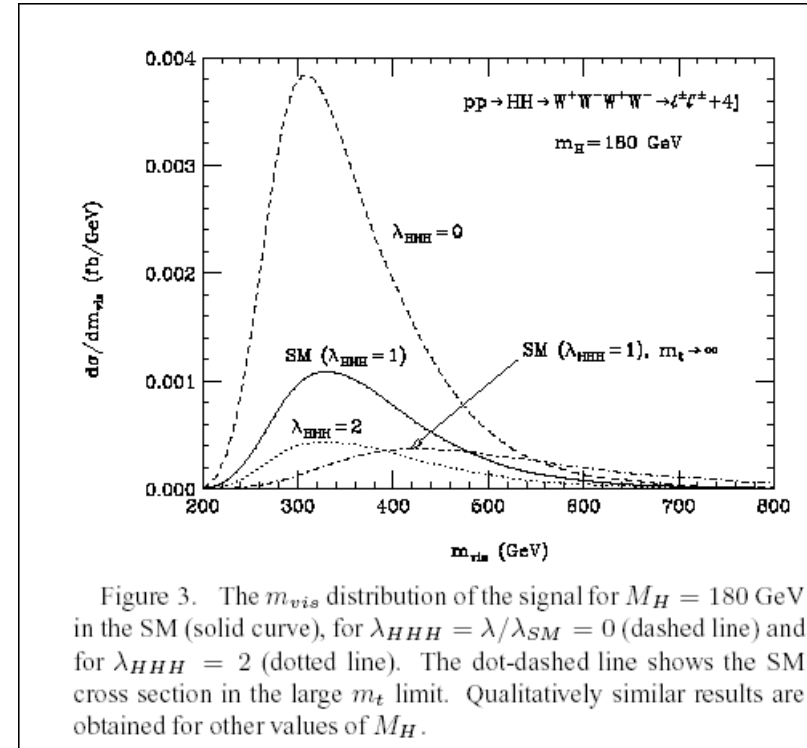
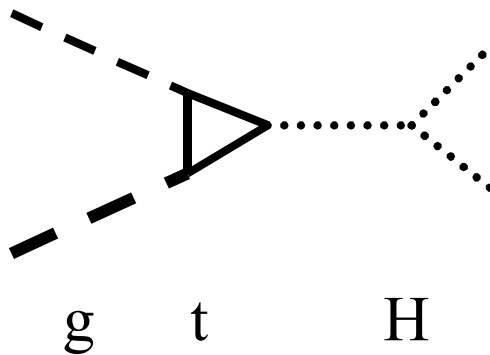
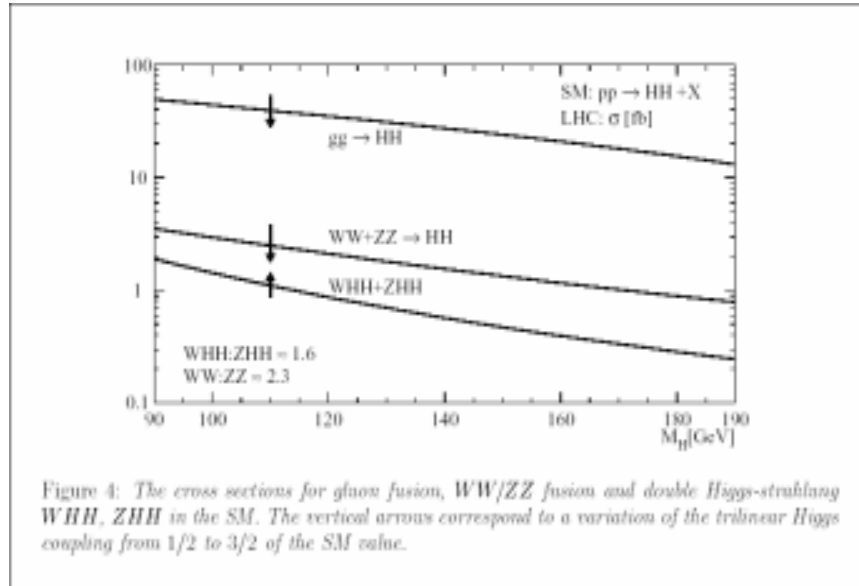


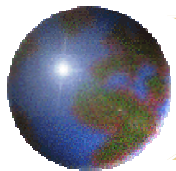
Figure 4. Limits achievable at 95% CL for $\Delta\lambda_{HHH} = \lambda_{HHH} - 1$ ($\lambda_{HHH} = \lambda/\lambda_{SM}$) in $pp \rightarrow \ell^+ \ell'^+ + 4j$ at the LHC. Bounds are shown for integrated luminosities of 300 fb^{-1} (solid lines), 600 fb^{-1} (dashed lines) and 3000 fb^{-1} (dotted lines). The allowed region is between the two lines of equal texture. The Higgs boson self-coupling vanishes for $\Delta\lambda_{HHH} = -1$.



Higgs Self-Coupling



Get a first indication of HHH coupling before the LC?

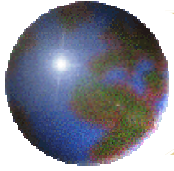


Higgs Self Coupling

Table 8: Expected numbers of signal and background events after all cuts for the $gg \rightarrow HH \rightarrow 4W \rightarrow \ell^+ \ell'^+ 4j$ final state, for $\int \mathcal{L} = 6000 \text{ fb}^{-1}$.

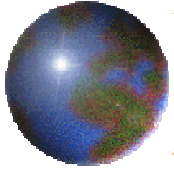
m_H	Signal	$t\bar{t}$	$W^\pm Z$	$W^\pm W^+ W^-$	$t\bar{t} W^\pm$	$t\bar{t} t\bar{t}$	S/\sqrt{B}
170 GeV	350	90	60	2400	1600	30	5.4
200 GeV	220	90	60	1500	1600	30	3.8

A $\sim 20\%$ measurement of the HHH coupling is possible for Higgs masses < 200 GeV. However, to realize these improvements we need to maintain the capabilities of the LHC detectors.



Kinematics

- A particle of mass M produced by partons of momentum fraction x_1, x_2 has the kinematic relations where $s = \text{C.M. energy squared}$. Thus, on average \sqrt{s}
- The rapidity y is , $\frac{\sinh y}{M/p} = \frac{1}{M/p} \sqrt{\frac{E + p_z}{E - p_z}}$. It is related to pseudorapidity, $\eta \sim y$.

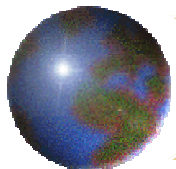


Tracker Front Ends, Links

Table 17: Hadron fluence and radiation dose in different radial layers of the CMS Tracker (barrel part) for an integrated luminosity of 2500 fb^{-1} .

Radius (cm)	Fluence of fast hadrons (10^{14} cm^{-2})	Dose (kGy)	Charged Particle Flux ($\text{cm}^{-2} \text{ s}^{-1}$)
4	160	4200	5×10^8
11	23	940	10^8
22	8	350	3×10^7
75	1.5	35	3.5×10^6
115	1	9.3	1.5×10^6

- **DSM is radiation hard in 0.25 μm**
- **0.13 μm is available – use Moore’s “law”**
- **Development needed – track and adapt commercial processes**
- **Optical links – development needed**

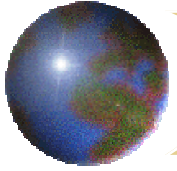


Dose for ECAL and HCAL

Table 18: The neutron fluence and radiation dose at shower maximum at different pseudorapidities for an integrated luminosity of 2500 fb^{-1} .

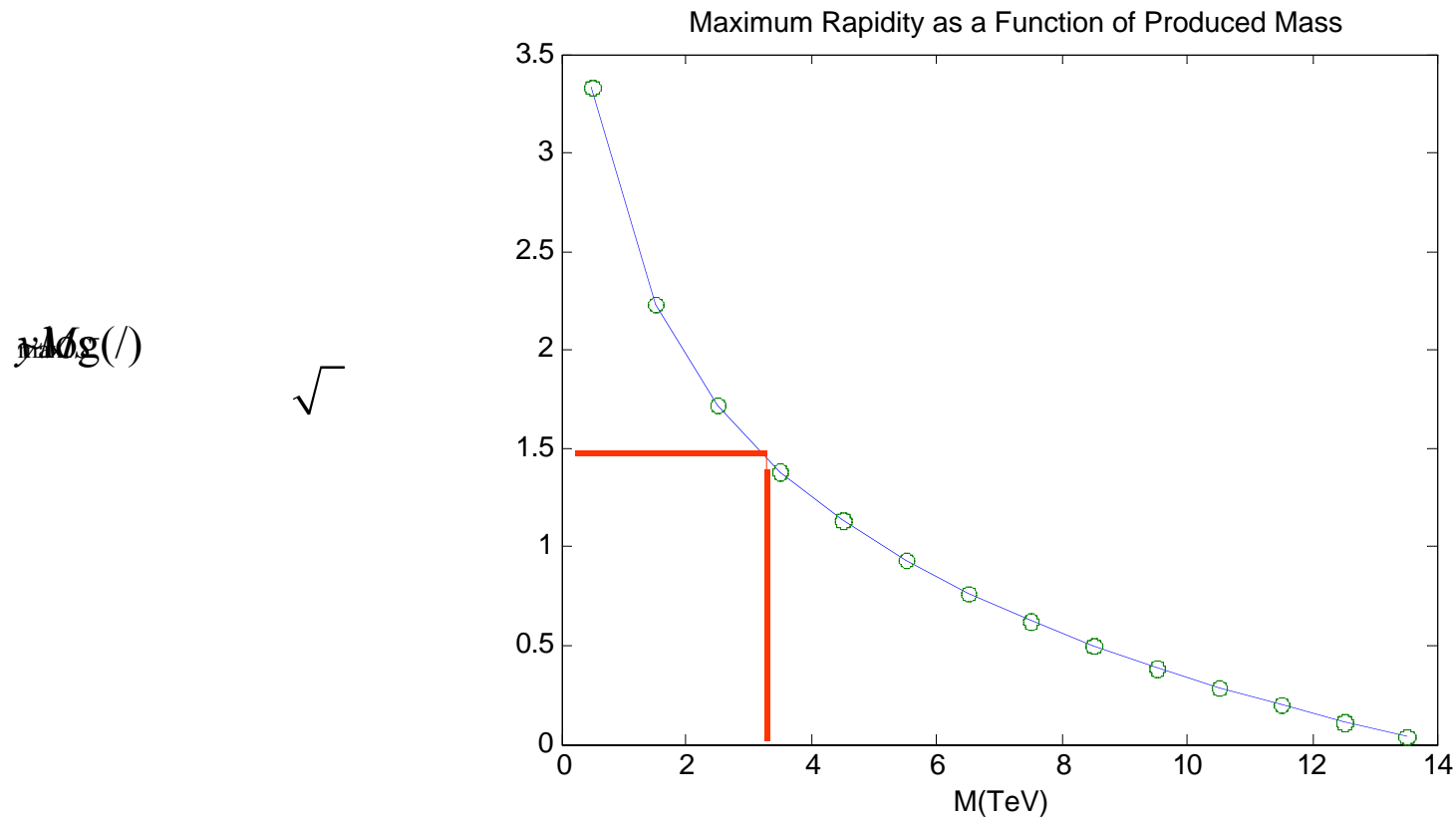
Pseudorapidity η	ECAL Dose (kGy)	HCAL Dose (kGy)	ECAL Dose Rate (Gy/h)
0 - 1.5	15	1	2.5
2.0	100	20	14
2.9	1000	200	140
3.5	-	500	-
5	-	5000	-

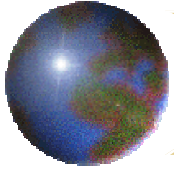
Barrel doses are not a problem. For the endcaps a technology change may be needed for $2 < |y| < 3$. For CMS endcap switch to quartz fiber?



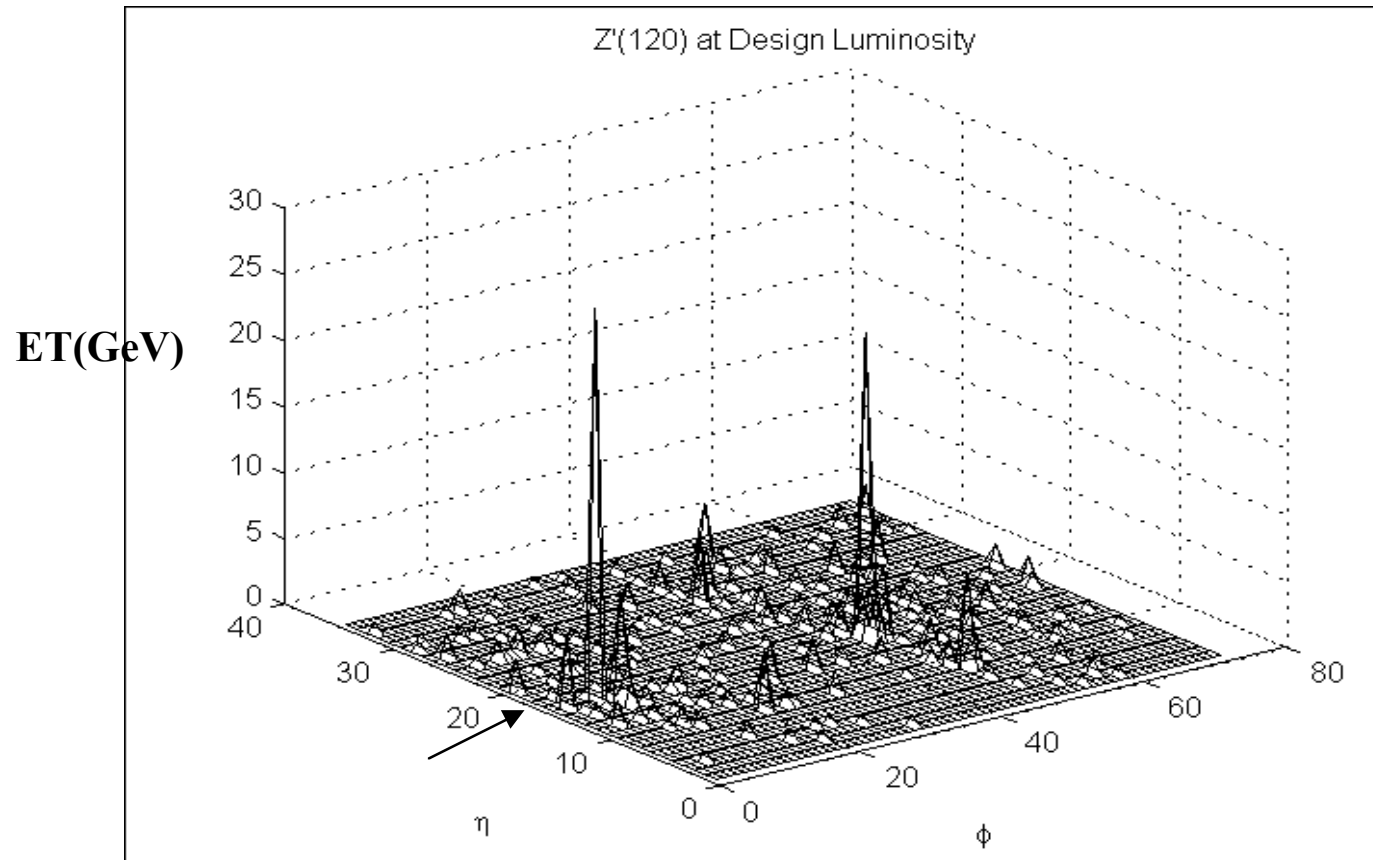
Produced Mass and y

- Heavy states are produced at wide angle to the beams and strike the “barrel”.

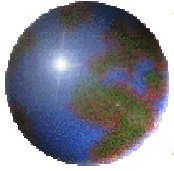




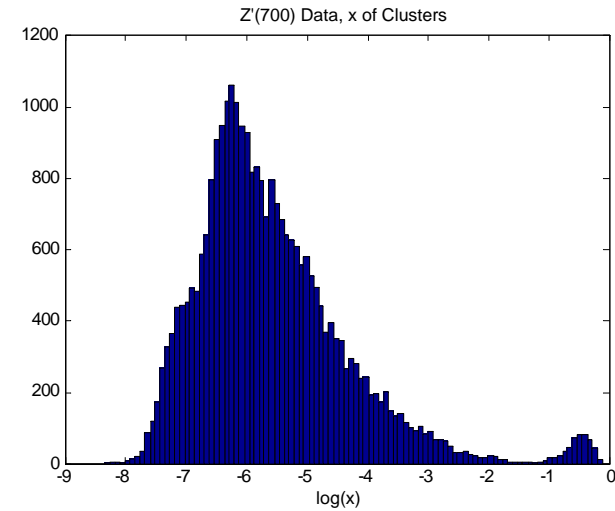
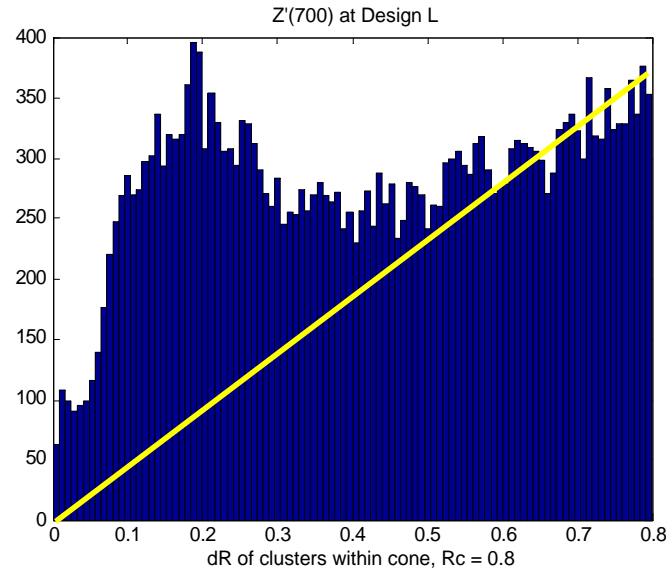
$Z'(120)$ at Design L



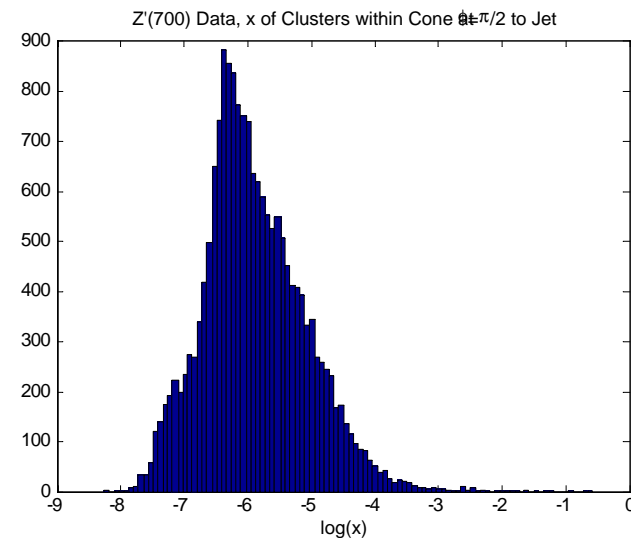
Note that the calorimeter cells are still fairly sparsely populated (granularity ≈ 0.017)

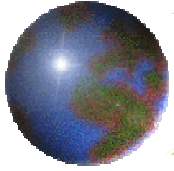


dR and log(x) for Z'(700)



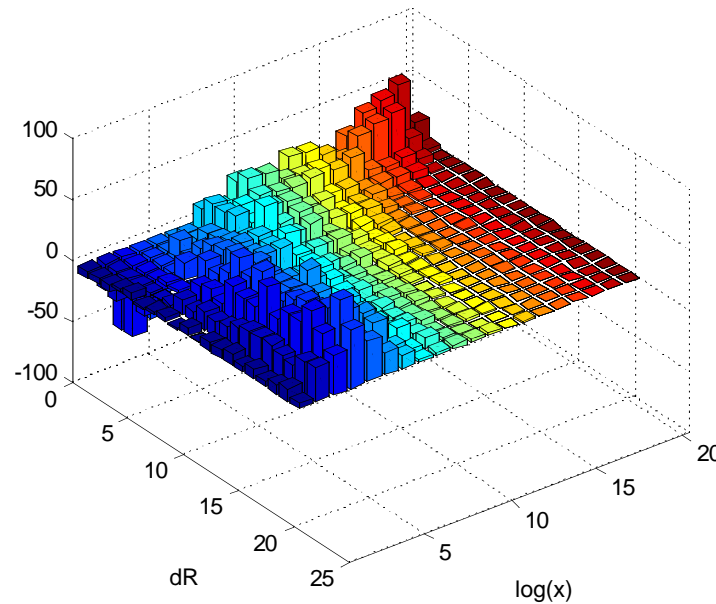
Clearly the pileup has $dN dR \sim R$ (area), while the jet shows a peak at $R \sim 0.2$. The pileup is at low x , while there is a jet "core" at high x . Use tracker to subtract charged pileup? Reduce pileup by 3x.



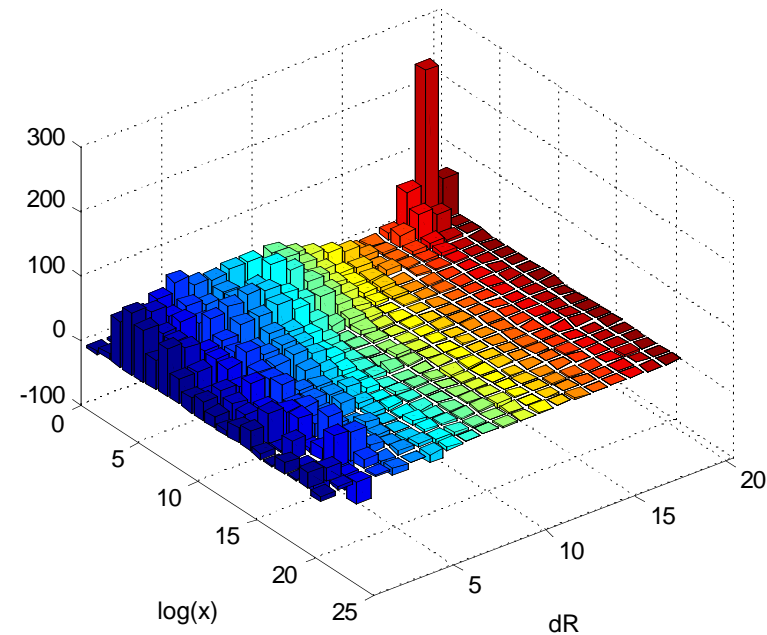


Pileup Subtracted ($\log x$, dR) Jet Contours

Jet With "Underlying Event" Defined $\phi_t = \pi/2$ to Jet Subtracted



$Z'(700)$ Data, ($dR, \log(x)$) $\phi = \pi/2$ "Jet" Subtracted



**For both $Z'(120)$ at low L and for $Z'(700)$ at design L there is a core of the jet at small dR and large x .
Extend to $10 \times L$ using tracker and energy flow inside the jet?**