

# Phenomenology of Extra Dimensions

Eduardo Pontón  
Columbia University

February 16, 2006

Aspen Winter Conference

# What do we expect at the Weak scale?

Higgs mass (hence EW scale) sensitive to high scale physics

Naturality: cutoff at around TeV → new physics!

- Supersymmetry (superpartners)
- Extra Dimensions (Kaluza-Klein modes)
- Higgs as pseudo NGB

Look for new particles at high-energy colliders. Theories predict specific relations between new and observed particles.

- Often new interactions are related to SM ones
- Lorentz transformation properties predicted: spins
- But also lots of "model building freedom" and new parameters

Exciting prospect of seeing extra dimensional nature at colliders  
Maybe a glimpse into even higher energies

# Outline

Overview of selected X-Dim. scenarios

Can we see several KK levels?

A six-dimensional example: UED's



Missing energy signals

Resonances

Conclusions

# Large "Gravitational" Dimensions

(Arkani-Hamed, Dimopoulos & Dvali)




- ADD:   $\delta$  flat, compact extra dimensions
-  Invisible to SM, but accessible to gravity

$$M_{\text{Pl}}^2 = V_\delta M_D^{2+\delta} \quad M_D = \mathcal{O}(\text{TeV})$$

Quantum gravity effects may lurk near the EW scale!

Almost a *continuum* of KK graviton states, with Planck suppressed couplings

Well defined couplings to matter:  $\frac{1}{M_{\text{Pl}}} h_{\mu\nu} T^{\mu\nu}$  or  $\frac{1}{\Lambda_H^2} T^{\mu\nu} T_{\mu\nu}$

-  Modifications of  $f\bar{f}$  production in  $e^+e^-$  and  $pp$  collisions
-  At LHC: Drell-Yan and diphoton production can probe scales up to a few TeV, depending on  $\delta$ .  
(Giudice, Rattazzi & Wells, Mirabelli, Perelstein & Peskin, ...)
-  Graviton emission:  $pp \rightarrow \text{jet} + \cancel{E}_T$  ( $qg \rightarrow qG$ )

# Warped Extra Dimensions

(Randall & Sundrum)

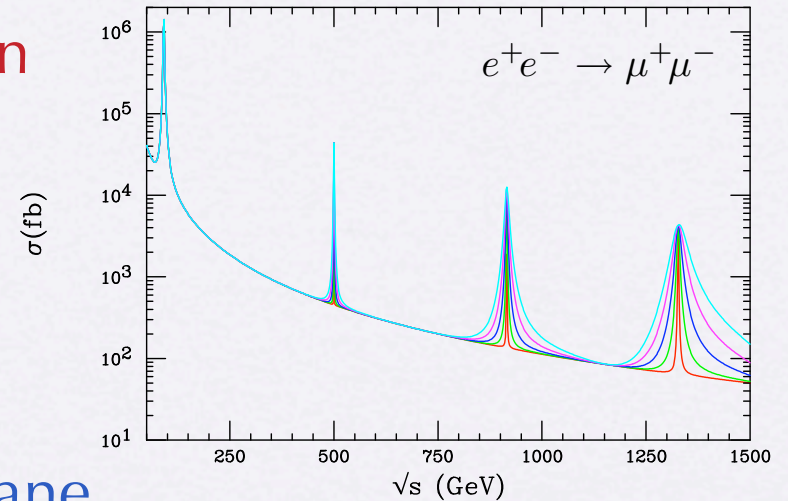
Solves hierarchy problem if Higgs on IR brane. Also  $M_{\text{Pl}}^2 \approx \frac{M^3}{k} e^{2k\pi r_c}$ .

Signal very sensitive to fermion localization

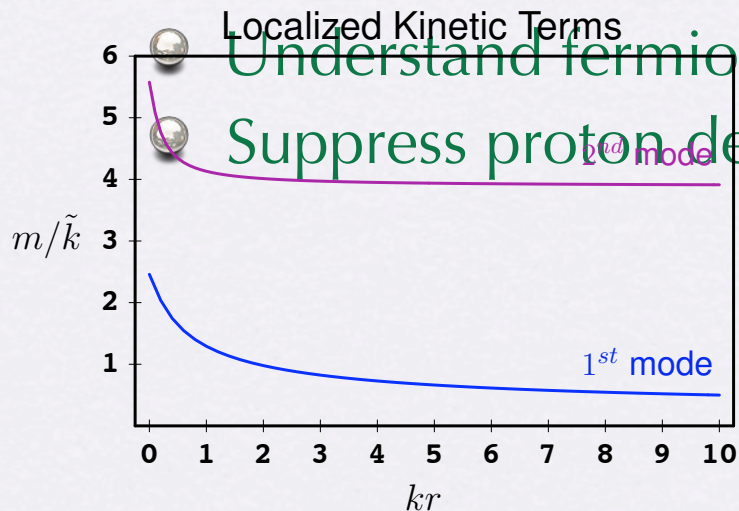
● If SM fermions localized on IR brane:

Graviton resonances:  
EW scale masses and couplings

● Light fermions may be away from IR brane



(from Davoudiasl, Hewett & Rizzo)



(Carena, E.P, Tait & Wagner)

Understand fermion mass hierarchies

Suppress proton decay or dangerous FCNC effects

● But gauge and fermion KK towers

● Important constraints from EWPM  
Mixing due to localized Higgs

● First KK level may be within reach of LHC

# Bulk Gauge Fields in TeV X-Dim

Gauge interactions get strong near the compactification scale

- **Fields localized** → at certain points in extra dimensions
  - **tree-level couplings to gauge fields** breaks down near the weak scale
  - **Hierarchy problem alleviated.** Important constraints from EWPM:  $1/R > 5 \text{ TeV}$
  - **In general, "little hierarchy problem" remains**
- **Combined 7 TeV at 95% C.L. (Cheung & Landsberg)**
- All SM fields delocalized → **Universal Extra Dimensions**
  - Present constraints on the order of a few 100 GeV
  - Potential production of several KK levels!
- **Some similarities with SUSY models with a degenerate spectrum**
  - Superpartners ↔ First KK level states (Cheng, Matchev & Schmaltz)
  - R-parity ↔ KK-parity
  - "Small" EW corrections

# Why 6 Dimensions?

Interesting theoretical constraints:

- Anomalies →
  - generations a multiple of three
  - predicts right-handed neutrinos

(Dobrescu & Poppitz)
- Discrete symmetries of the compactified theory
  - $Z_2^{\text{KK}}$  : KK-Parity → LKP is stable (similar to 5D)
  - $Z_8$  :
    - Nucleon decay suppressed
    - Neutrinos are Dirac fermions

(Appelquist, Dobrescu, E.P. & Yee)

These symmetries are tightly related to the higher dimensional Lorentz invariance

Useful to distinguish between

Tree-level effects

Loop-level effects



# Tree-level Structure

Tree-level spectrum:  $M_{j,k}^2 = \frac{j^2 + k^2}{R^2}$

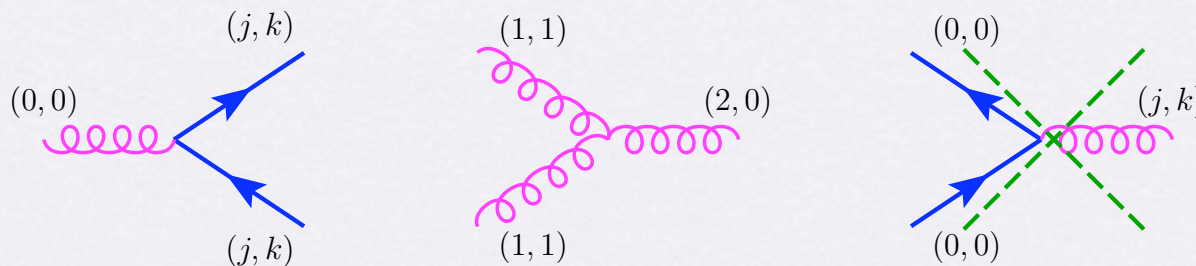
- In detail:  $\frac{1}{R}$ ,  $\frac{\sqrt{2}}{R}$ ,  $\frac{2}{R}$ , ...  
 $(1,0)$     $(1,1)$     $(2,0)$

└ does not exist in 5D

- n-th KK level states degenerate at tree-level

Couplings determined by wavefunction profiles (boundary cond.)

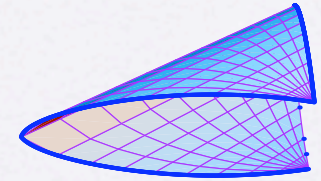
KK number conservation: related to momentum conservation



At tree-level: only pair production of KK modes!

# Radiative corrections

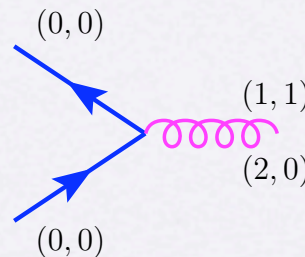
- Low-energy chiral theory requires "singularities"



→ Localized operators and new effects!

- Mass splittings: depend on quantum numbers

- New interactions, e.g.



$$P_{\text{KK}} = (-1)^{j+k}$$

- Theory has a cutoff at  $\Lambda \sim \mathcal{O}(10)/R$  (when QCD gets strong)

- Localized operators receive contributions from

- Unknown UV completion at  $\Lambda$  (assume NDA philosophy)

- "Known" physics below  $\Lambda$  (loop induced and logarithmically enhanced)

# KK Spectrum in 6D

1-loop leading order mass corrections take the form

$$\delta M_{j,k}^F = \frac{1}{16\pi^2} \ln \frac{\Lambda^2}{\mu^2} C_F M_{j,k}$$

One finds:

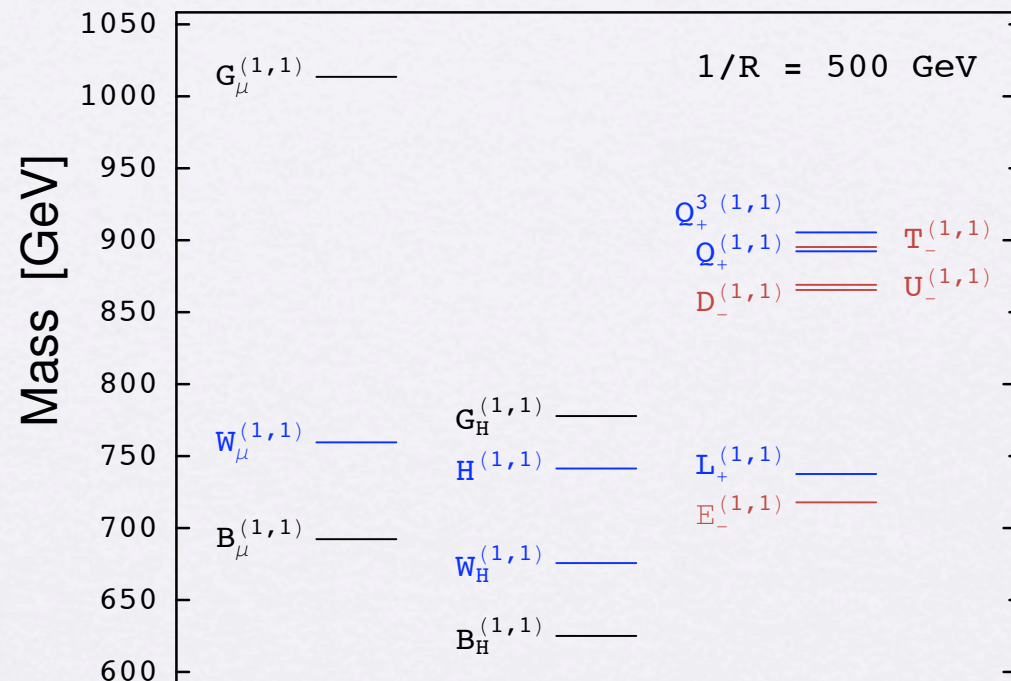
$$C_{A_\mu} = g_4^2 \left( \frac{17}{3} C(G) - \frac{2}{3} \sum_{\Psi} T(r) - \frac{1}{2} \sum_{\Phi} T(r) \right)$$

$$C_{A_H} = g_4^2 \left( 9 C(G) - 4 \sum_{\Psi} T(r) + \frac{7}{2} \sum_{\Phi} T(r) \right)$$

$$C_{\Psi} = 5 \sum_{\text{gauge}} g_4^2 C_2(\Psi) + \frac{3}{4} \sum_i |\lambda_{4,i}|^2$$

Take these as a guide keeping in mind theoretical uncertainties

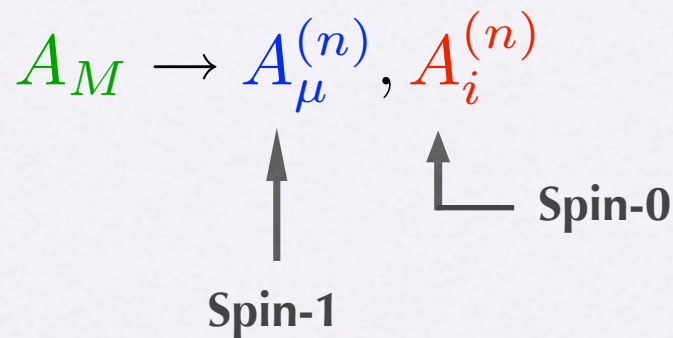
(1,1) level states



(E.P. & Wang)

Quarks and gluons should receive larger corrections than leptons and electroweak gauge bosons.

# Scalars from Gauge fields



- One linear combination eaten by the massive gauge fields  $\rightarrow$  Higgs Mech.

- 5D: nothing left

- 6D and higher: new scalars  $\rightarrow$  "Spinless Adjoints"

- Considering possibility of the Higgs from  $A_M \rightarrow$  enlarge gauge group

$$G_H^a \quad W_H^\pm \quad W_H^3 \quad B_H$$

- Interactions constrained by KK-parity and gauge invariance

- No dimension-4 couplings to fermions or gluons

- Higher dimension operators, e.g.

$$\bar{Q} \Gamma^M \Gamma^N \Gamma^L Q \partial_M G_{NL} \quad \text{or} \quad G_{MN} G^{NL} G_L^M$$

# Scalars from Gauge fields

• The couplings to gluons vanish!  $f^{abc} G_a^{\mu\nu} G_{\mu\nu}^b G_H^{(j,k)c} = 0$

• Couplings to fermions exist if  $(-1)^{j+k} = +1$  :

$$(\bar{q} \gamma^\mu T^a q) \partial_\mu G_H^{(j,k)a} \rightarrow \text{couplings proportional to } m_q$$

$\uparrow$                        $\uparrow$   
 No chirality flip

• The KK-parity even spinless adjoints are hard to produce from  $\bar{q}q$  or  $gg$  initial states

• But they can be produced in decays of KK quarks or fermions

• They decay into top quarks almost always!

• In 6D, the LKP is expected to be  $B_H^{(1,0)}$ , a scalar

# Pair Production: the SUSY-like case

Involves KK-number *conserving* interactions

- 1st level states: cascade decays to LKP

- Strong production of  $Q^1 Q^1$

Soft jets  $\rightarrow$  hard!

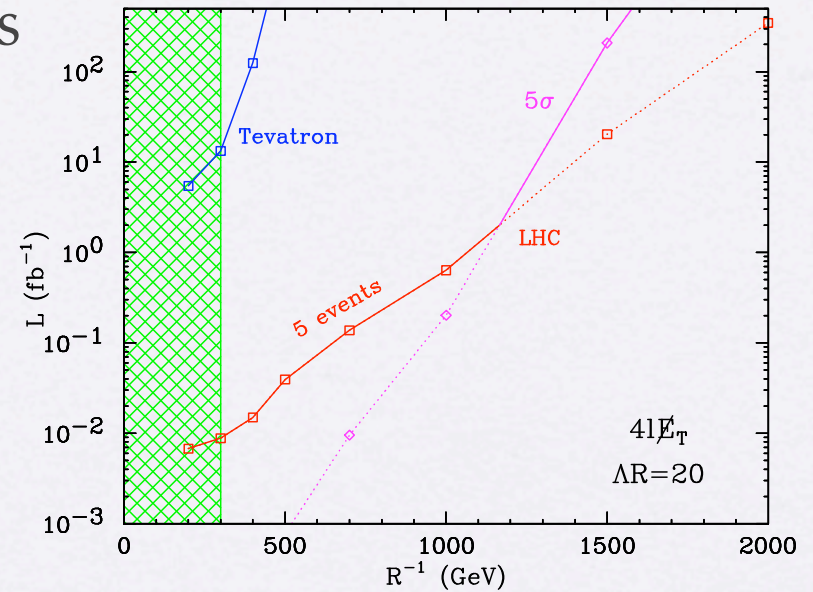
$\rightarrow$  look for 3 or more lepton plus missing energy signal

- Decays into  $W_1^\pm$  and  $Z_1 \rightarrow$  leptons

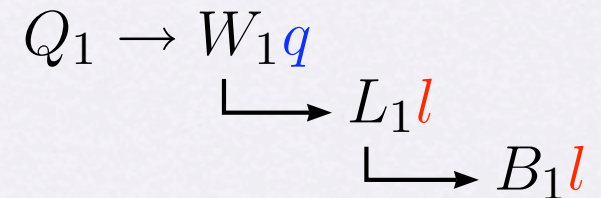
“Gold-plated” mode:  $\cancel{E}_T + 4$  leptons

- Heavier KK states (phase space suppression)

- Spin Measurements at hadron collider are challenging!



(from Cheng, Matchev & Schmaltz)



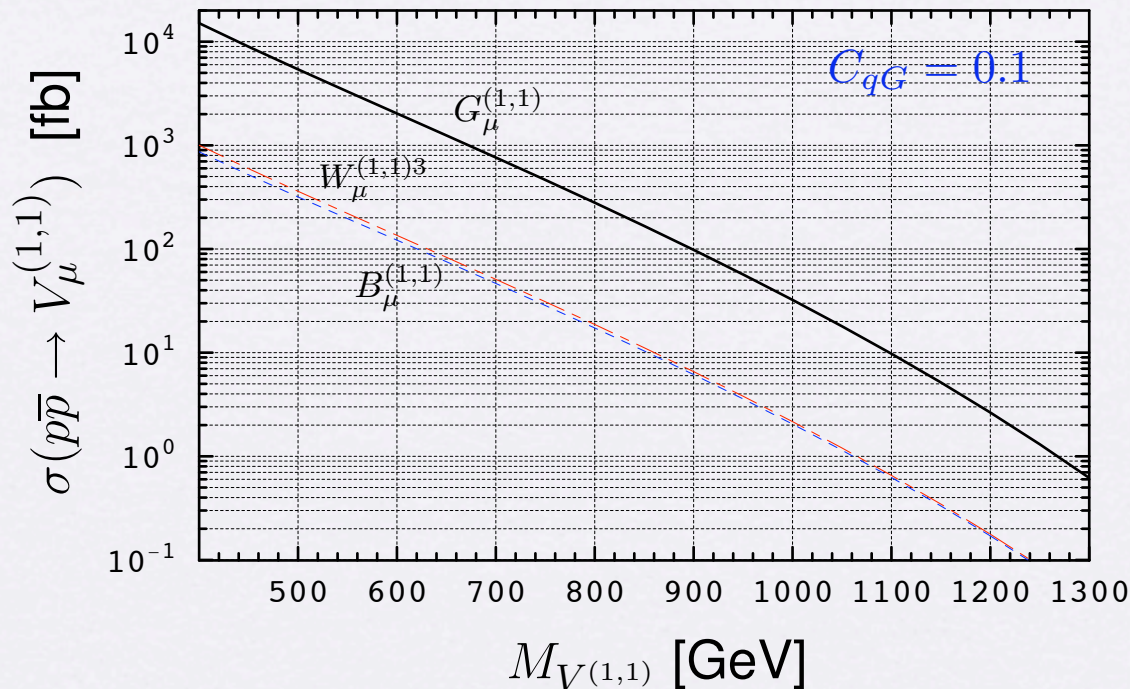
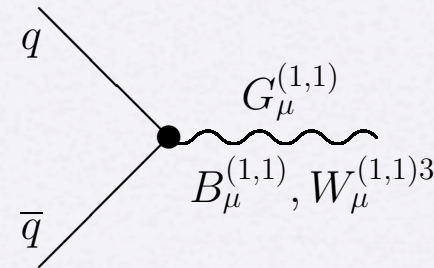
(Smillie and Weber;  
Datta, Kong & Matchev)

# Vector Mode Production at the Tevatron

(Dobrescu, E.P & Burdman)

- Single production (through KK *violating* interactions) of (1,1) and (2,0) states, e.g.

$$q\bar{q} \rightarrow G_\mu^{(1,1)}, W_\mu^{(1,1)3}, B_\mu^{(1,1)}$$



In 6D:

- Couplings to fermions

$$g_s C_{qG} (\bar{q} \gamma^\mu T^a q) G_\mu^{(1,1)a}$$

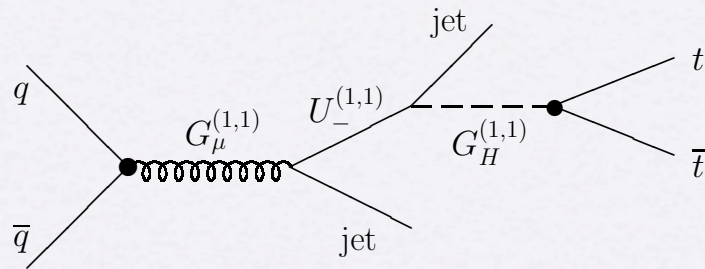
- Mass shifts

$$\delta M_{j,k}^2 / M_{j,k}^2$$

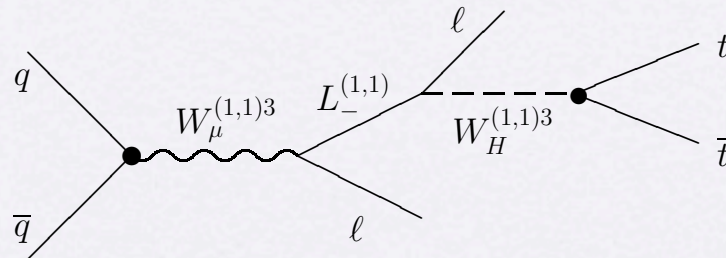
controlled by independent fundamental parameters

# $t\bar{t}$ Resonances at the Tevatron

(1,1) excitations of the gauge bosons have large BR into spinless adjoints, hence into  $t\bar{t}$

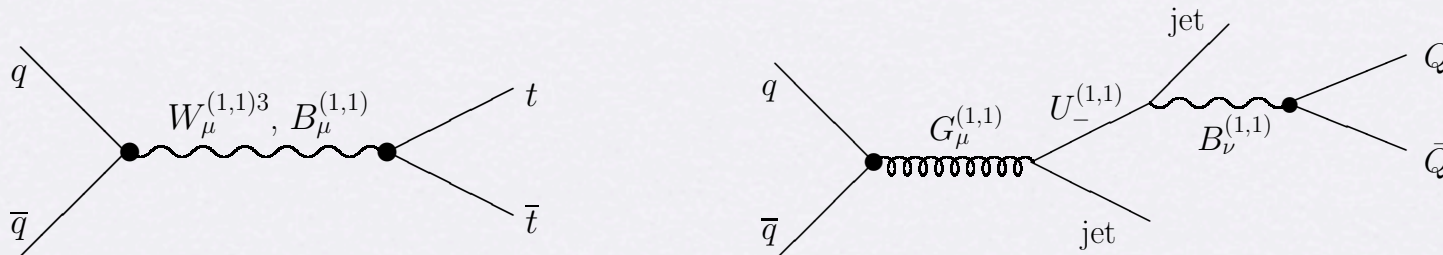


or also



decay modes	$G_\mu^{(1,1)}$
$G_H^{(1,1)} + \text{jets}$	60.5
$W_H^{(1,1)3} + \text{jets}$	3.2
$W_H^{(1,1)\pm} + \text{jets}$	6.1
$B_H^{(1,1)} + \text{jets}$	4.8
$W_\mu^{(1,1)3} + \text{jets}$	4.3
$W_\mu^{(1,1)\pm} + \text{jets}$	7.0
$B_\mu^{(1,1)} + \text{jets}$	9.3
$t\bar{t}$	0.5
$b\bar{b}$	0.8
dijet (no $b\bar{b}$ )	3.3

EW gauge bosons can also decay into top pairs





# $t\bar{t}$ Resonances at the Tevatron

Three (potential) peaks:

$$G_H, W_\mu \rightarrow 1.1 \times M_{1,1}$$

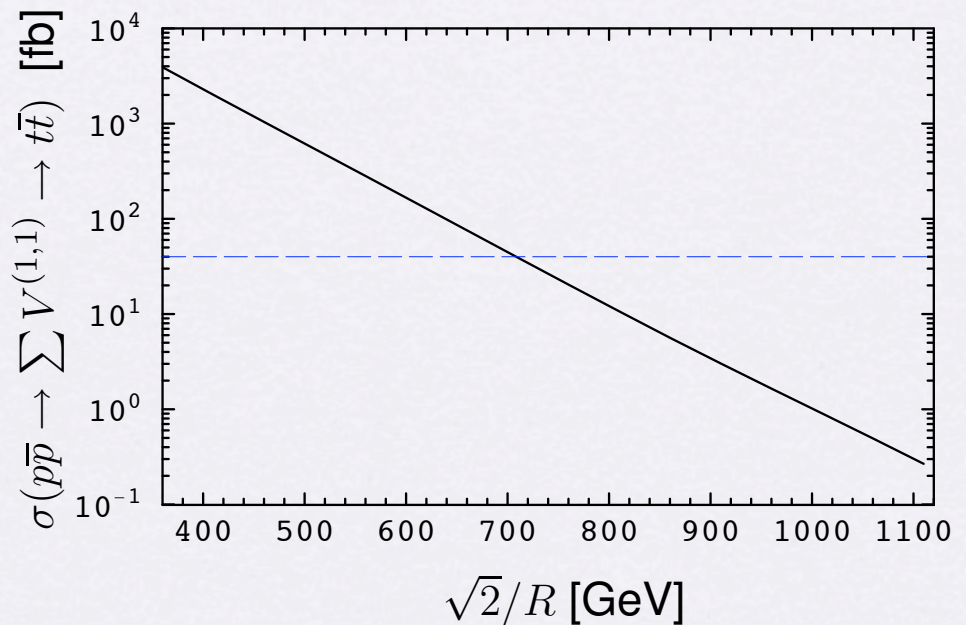
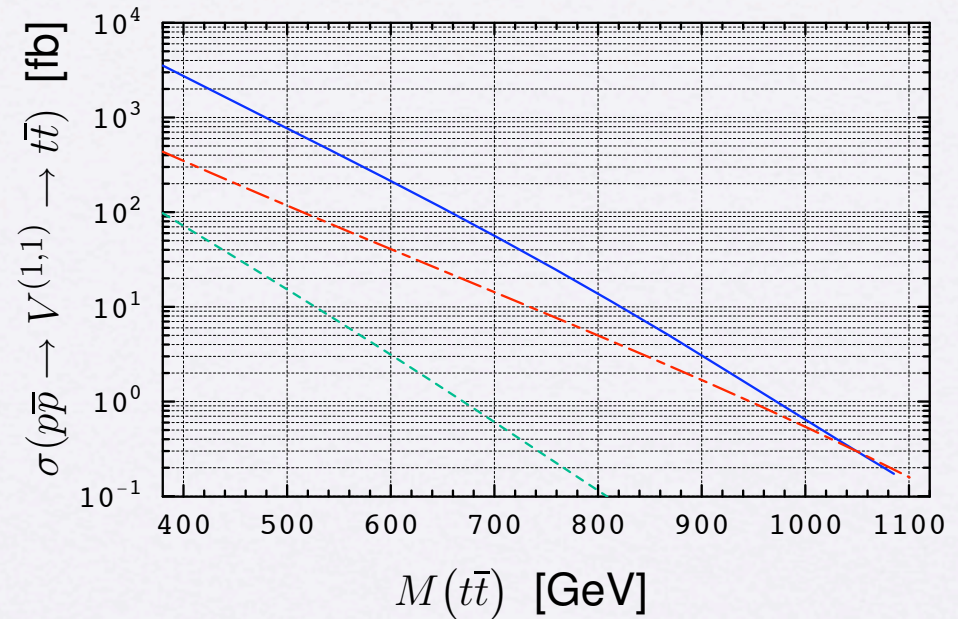
$$W_H, B_\mu \rightarrow 0.97 \times M_{1,1}$$

$$B_H \rightarrow 0.86 \times M_{1,1}$$

Current limit from CDF  
with  $682 \text{ pb}^{-1}$ :

$> 1 - 2 \text{ pb}$  at 95% C.L. for  
 $350 \text{ GeV} < M(t\bar{t}) < 600 \text{ GeV}$

→ No limit on  $1/R$

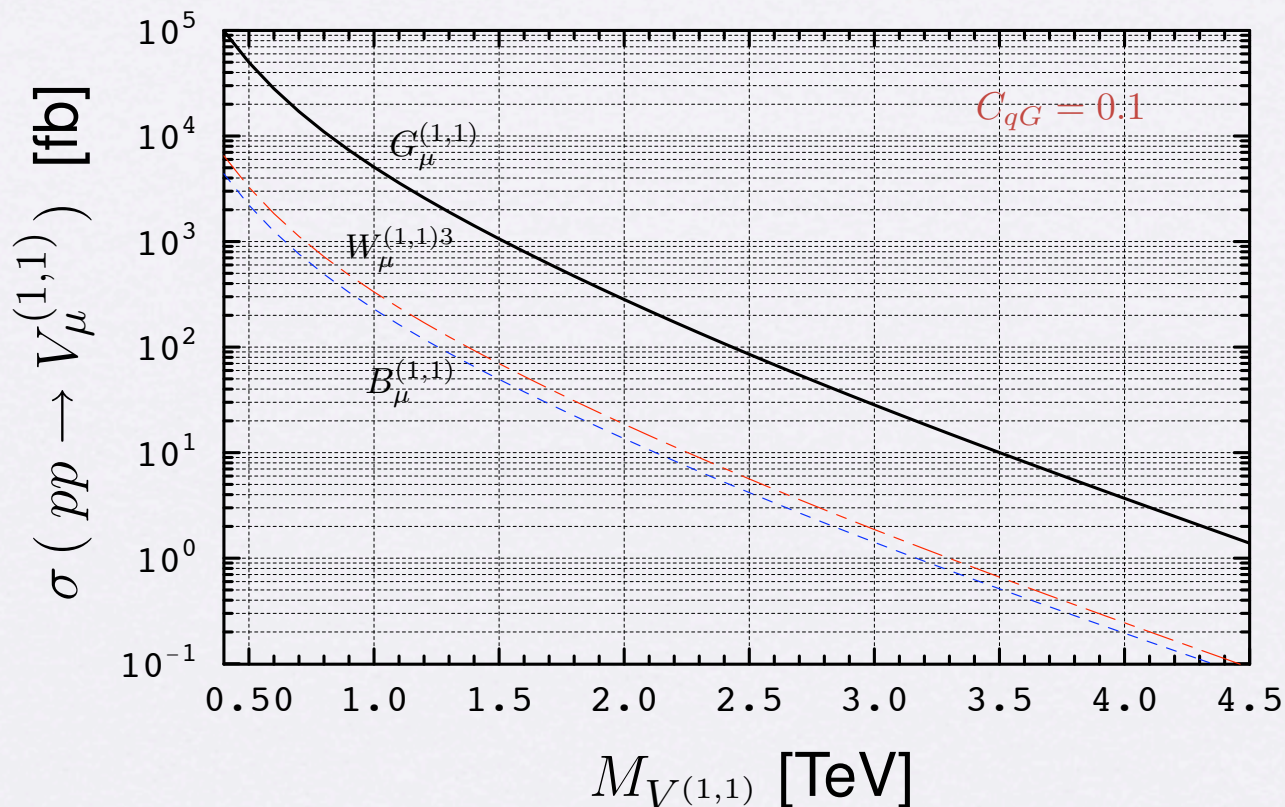


# Vector Mode Production at the LHC

- Single production at the LHC

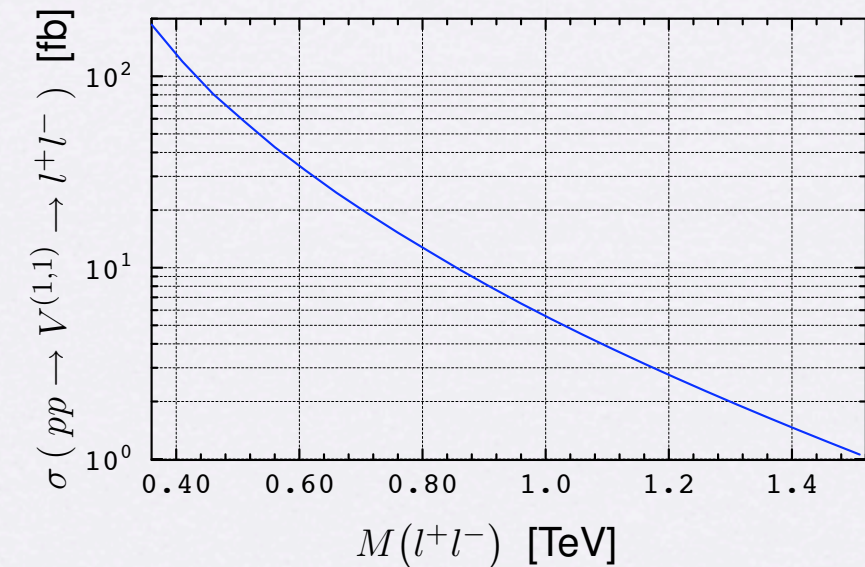
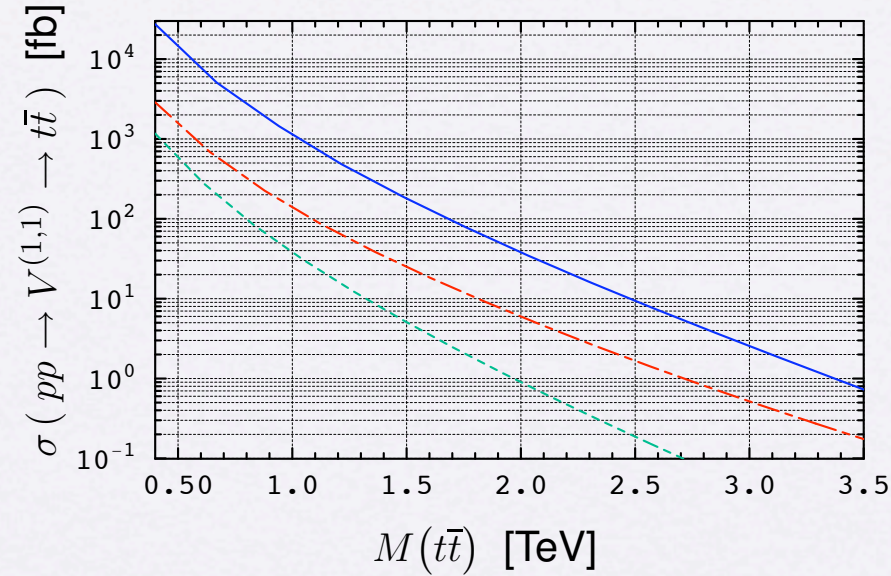
Initial  $gg$  or  $qg \rightarrow G^{(1,1)}$  (not log enhanced)

$\rightarrow q\bar{q} \rightarrow G^{(1,1)}, W^{(1,1)3}, B^{(1,1)}$  probably still dominant

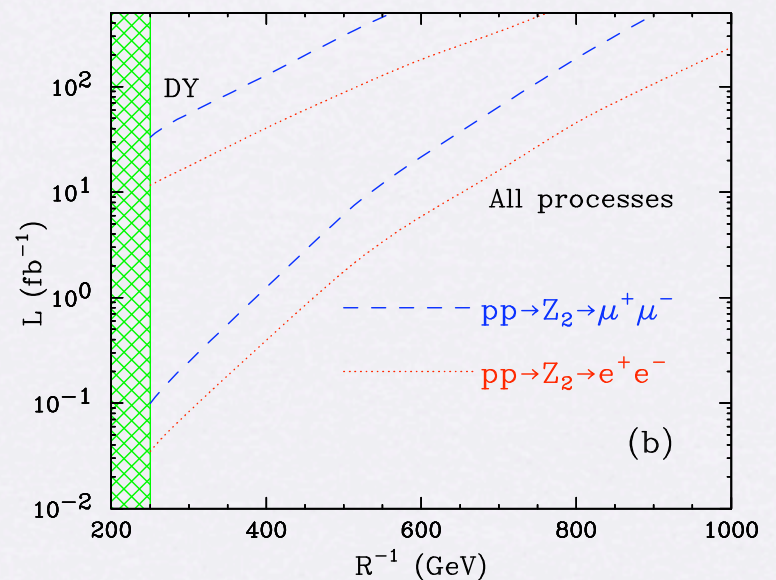
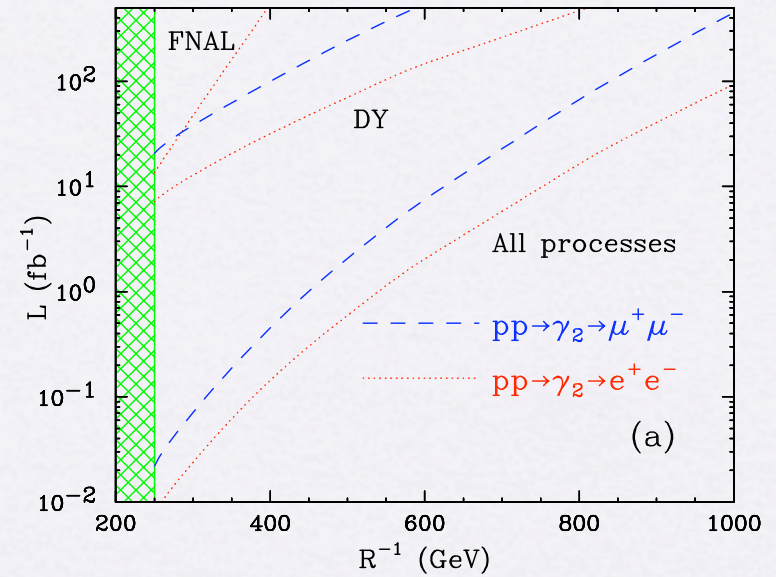


# Resonances at the LHC

In 6D: (Dobrescu, E.P & Burdman)



In 5D: (from Datta, Kong & Matchev)



# Conclusions

- Extra dimensions at the TeV scale are well motivated, and can be observed at high-energy colliders
- We may see several KK levels in UED scenarios!
  - striking manifestation of extra dimensions
- Observation of (1,1) level at  $\sim \sqrt{2}/R$  and (2,0) level at  $\sim 2/R$  allows a clear distinction between 5D and 6D
  - Should hold for  $B_{\mu}^{(1,1)}$  (small corrections)
- New scalars in 6D with large couplings to tops
  - Look for bumps in  $t\bar{t}$  at the Tevatron and LHC!
- $B_{\mu}^{(1,1)}$  resonance in dilepton channel also promising
- The SUSY/UED "confusion" is probably resolvable at hadron colliders, even in the absence of a spin measurement