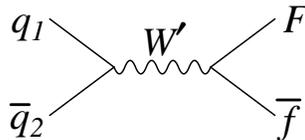


## SEARCHING FOR $W'$ BOSONS AT HADRON COLLIDERS



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Based on:

Z.S., Phys. Rev. D 66, 075011 (2002) [hep-ph/0207290]

Z.S., QCD Moriond 2003 [hep-ph/0306266]

Z.S., in production

Lots of room for further study.

# $W'$ bosons appear in many models

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**Working definition:** A  $W'$  boson is any particle that mediates a flavor-changing charged vector or axial-vector current.

**Left-right symmetric models: Broken  $SU(2)_L \times SU(2)_R$**

- Generic mixing of  $W_L$ – $W_R$   
R. N. Mohapatra, J. Pati, A. Salam, G. Senjanovic, . . .
- Orbifold-breaking — suppressed mixing, enhanced couplings  
Y. Mimura, S. Nandi, . . .
- Supersymmetric  $L$ – $R$  models  
M. Cvetič, J. Pati, . . .

**Models with additional left-handed  $W'$**

- Little Higgs:  $SU(5)/SO(5)$ ,  $SU(6)/SP(6)$ ,  $SU(N)/SU(N-1)$ , . . .  
T. Gregoire, N. Arkani-Hamed, S. Chang, H. C. Cheng, A. Cohen, I. Low, D. E. Kaplan, E. Katz, O. C. Kong, A. Nelson, M. Schmaltz, W. Skiba, D. Smith, J. Terning, J. Wacker, . . .
- Topcolor — topflavor, leptophobic topflavor seesaw, generic mixing  
H. Georgi, H. J. He, E. Jenkins, X. Li, E. Ma, E. Malkawi, D. Muller, S. Nandi, E. Simmons, T. Tait, C. P. Yuan, . . .
- Extra dimensions: Kaluza-Klein modes of the  $W$   
A. Datta, P. O'Donnell, T. Huang, Z. Lin, X. Zhang, . . .
- Non-commuting extended technicolor  
R. Chivukula, E. Simmons, J. Terning, . . .

**+ 1000's more**

# Fully differential NLO $W'$ cross section

Z.S., PRD 66, 075011 (2002) [hep-ph/0207290].

Calculated widths and 2-to-2 cross section for arbitrary couplings:

$$\mathcal{L} = \frac{1}{\sqrt{2}} \bar{f}_i \gamma_\mu \left( g_R e^{i\omega} \cos \zeta V_{f_i f_j}^R P_R + g_L \sin \zeta V_{f_i f_j}^L P_L \right) W' f_j + \text{H.c.}$$

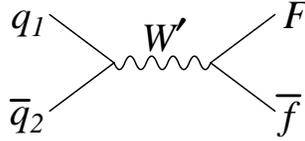
Proved factorization of couplings through NLO in width:

$$\frac{\Gamma(W')}{\Gamma_{SM}} = |g V_{f_i f_j}'|^2 \equiv |g_L \sin \zeta V_{f_i f_j}^L|^2 + |g_R \cos \zeta V_{f_i f_j}^R|^2$$

$$\Gamma_{\text{NLO}}(W' \rightarrow \nu \bar{\ell}) = \frac{g^2}{16\pi} |V_{\nu \ell}'|^2 \frac{M_{W'}}{3} \quad (\text{if left-handed})$$

$$\Gamma_{\text{NLO}}(W' \rightarrow q \bar{q}') = \frac{g^2}{16\pi} |V_{qq'}'|^2 M_{W'} (1 + \alpha_s/\pi)$$

$$\Gamma_{\text{NLO}}(W' \rightarrow t \bar{q}') = \frac{g^2 \beta^2}{16\pi M_{W'}} |V_{tq'}'|^2 (M_{W'}^2 + m_t^2/2) [1 + \alpha_s/\pi F(\beta)]$$



The fully differential LO cross section including width effects is

$$\sigma_0 = \int q_1 \otimes \bar{q}_2 \otimes \frac{1}{2\hat{s}} \frac{2N_c}{3} \frac{R_t t(t - m^2) + R_u u(u - m^2)}{(\hat{s} - M_{W'}^2)^2 + M_{W'}^2 \Gamma_{W'}^2} d\text{PS}_2$$

$$R_{t/u} = (|V_i|^2 + |A_i|^2)(|V_f|^2 + |A_f|^2) \pm 4\text{Re}(V_i A_i^*)\text{Re}(V_f A_f^*)$$

$$|V|^2 + |A|^2 = \frac{|g_R \cos \zeta V_{f_i f_j}^R|^2 + |g_L \sin \zeta V_{f_i f_j}^L|^2}{4}$$

$$2\text{Re}(V_i A_i^*) = \frac{|g_R \cos \zeta V_{f_i f_j}^R|^2 - |g_L \sin \zeta V_{f_i f_j}^L|^2}{4}$$

If  $W' \equiv W'_L, W'_R$ , or small mixing:  $|V| = |A|$

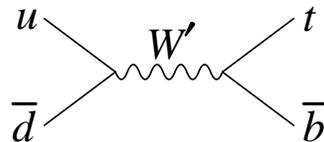
$$\Rightarrow R_u = 0; R_t = \frac{g^4 |V_i'|^2 |V_f'|^2}{8}$$

$$\Rightarrow \sigma_{\text{NLO}} = \sigma_{\text{NLO}}^{\text{SM}} \text{ up to overall factor: } (g'/g_{\text{SM}})^4$$

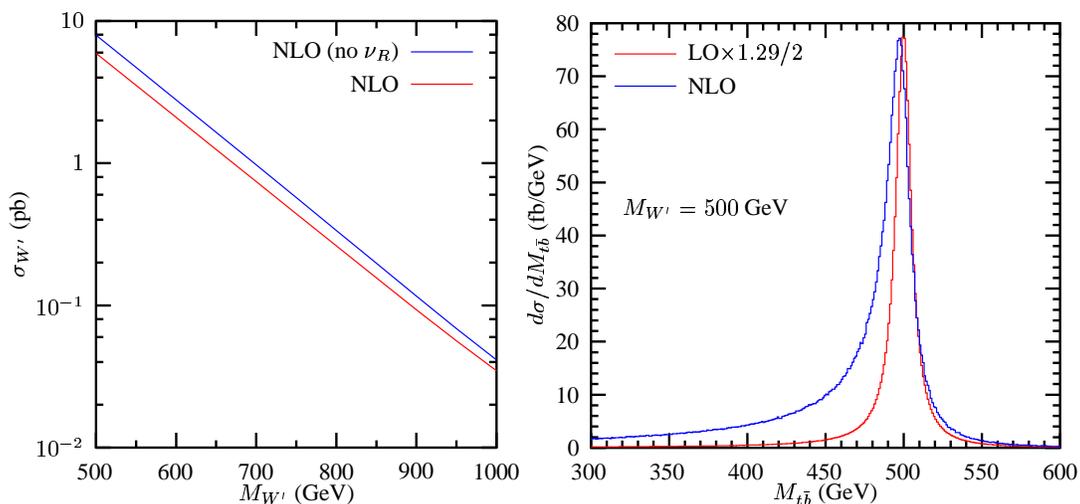
# A model independent search for $W'$

Z.S., PRD 66, 075011 (2002) [hep-ph/0207290].

The single-top-quark cross section is the same whether you have a  $W'_L$ ,  $W'_R$ , or small mixing. Thus this is a great channel for model-independent searches.



$pp\bar{p} \rightarrow W' \rightarrow t\bar{b}$  production at the Tevatron



CDF just set a bound of  $M_{W'} > 536(566)$  GeV. PRL 90, 081802 (03)

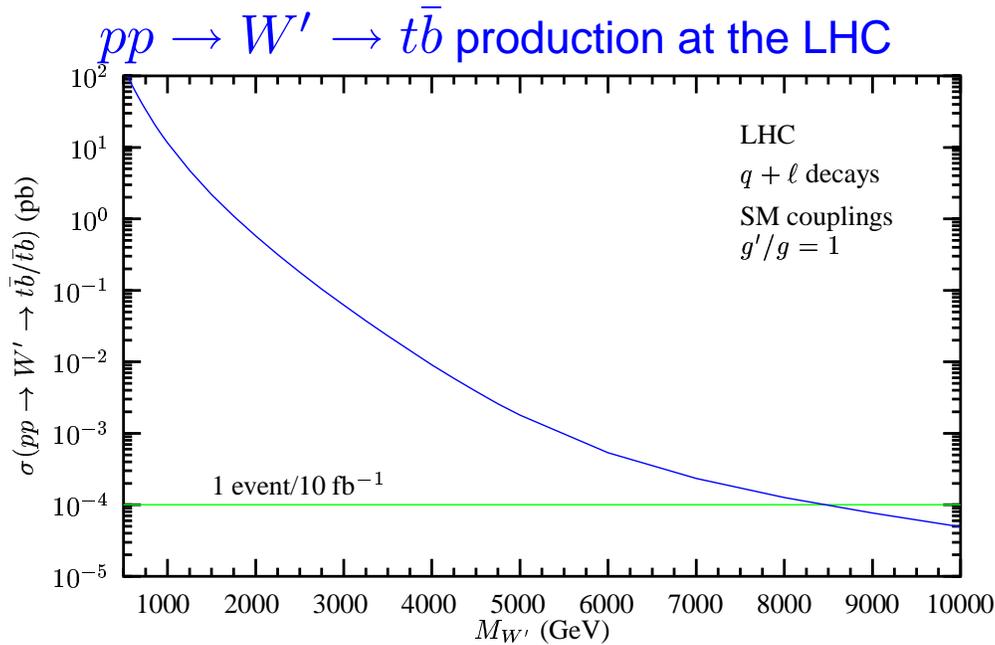
Previous limit for  $W'_R$  was  $\sim 400$  GeV ( $g' = g_{\text{SM}}$ ),  
CDF limit is  $\sim 786$  GeV for  $W'_L \rightarrow \ell\nu$ .

The mass reach should extend to 800–900 GeV in Run IIA.

If a  $W'$  is found we can use angular correlations of decay products to determine whether it looks like a  $W'_L$ ,  $W'_R$ , or mixed state.

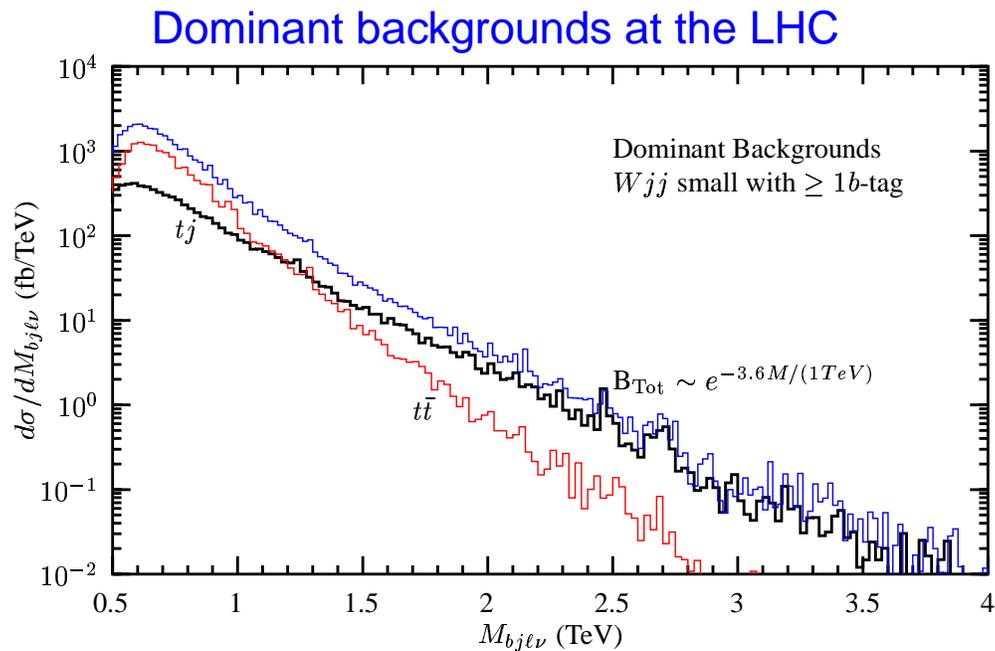
# The search for $W'$ at the LHC

Z.S., QCD Moriond 2003 [hep-ph/0306266].



About 50, 10 TeV  $W'$ 's are produced in 100 fb<sup>-1</sup> of data.

Can we find them?



Total Background includes:  $t\bar{t}$ ,  $t$ -channel single-top,  $Wjj$ ,  $Wc\bar{j}$ ,  $Wb\bar{b}$ ,  $Wc\bar{c}$ ,  $WZ$ ,  $Wt$ , &  $s$ -channel single-top.

## Simulation details

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### Method

Signal (S) and backgrounds (B) used PYTHIA, except for high-mass backgrounds which use a ME calculation fit to differential NLO.

Note: PYTHIA/HERWIG alone fail by 2 orders of magnitude simulating single-top and  $Wjj$  for our cuts. Generally difficult to generate high-mass backgrounds.

Used SHW detector simulation updated to match ATLAS.

### Cuts

Demand at least 2 jets ( $w/\geq 1b$ -tag), a lepton ( $e/\mu$ ), and  $\cancel{E}_T$

$$E_{Tj_1} > \max[200 \text{ GeV}, \min(10\%M_{W'}, 500 \text{ GeV})] \quad |\eta_{j_1}| < 2.5$$

$$E_{Tj_2} > \min(10\%M_{W'}, 150 \text{ GeV}) \quad |\eta_{j_2}| < 2.5$$

$$E_{T\ell} > 30 \text{ GeV} \quad |\eta_\ell| < 2.5$$

$$\cancel{E}_T > 50 \text{ GeV}$$

$k_T$ -jets  $R=1$  (this is like a cone of 0.7)

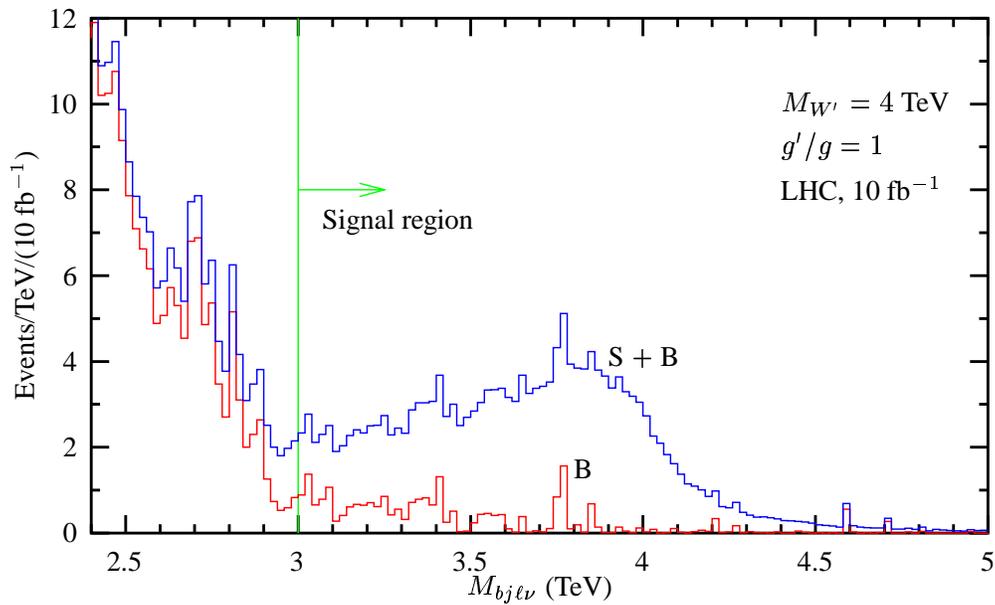
$100 \text{ GeV} < M_{j_2\ell\nu} < 450 \text{ GeV}$  (reduces  $Wjj$ , but can be ignored)

Finally reconstruct  $M_{bj\ell\nu}$ , and place cut just below target mass (up to 3 TeV).

Aside: Triggered on lepton, but a high- $\cancel{E}_T$  jet trigger might have been better.

## What does the signal look like?

Signal+Background vs. Background for  $M_{W'} = 4 \text{ TeV}$

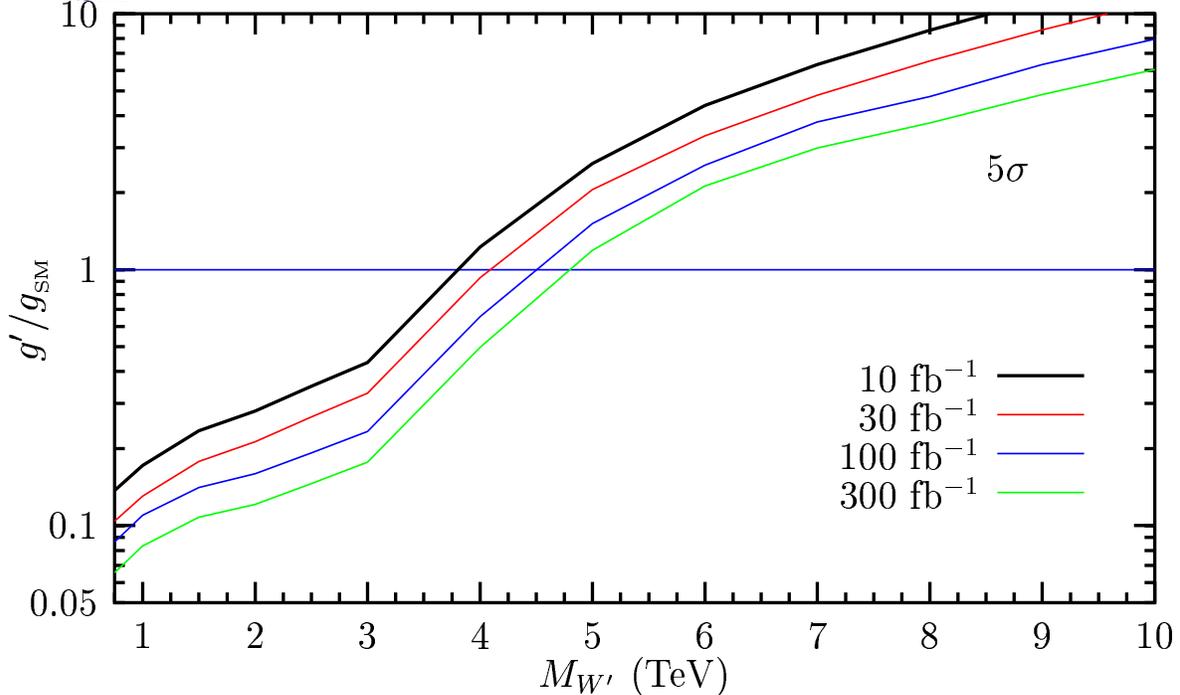


For a  $W'$  with SM-like couplings, we expect  $\sim 4$  events per low-luminosity year over a background of less than 1 event.

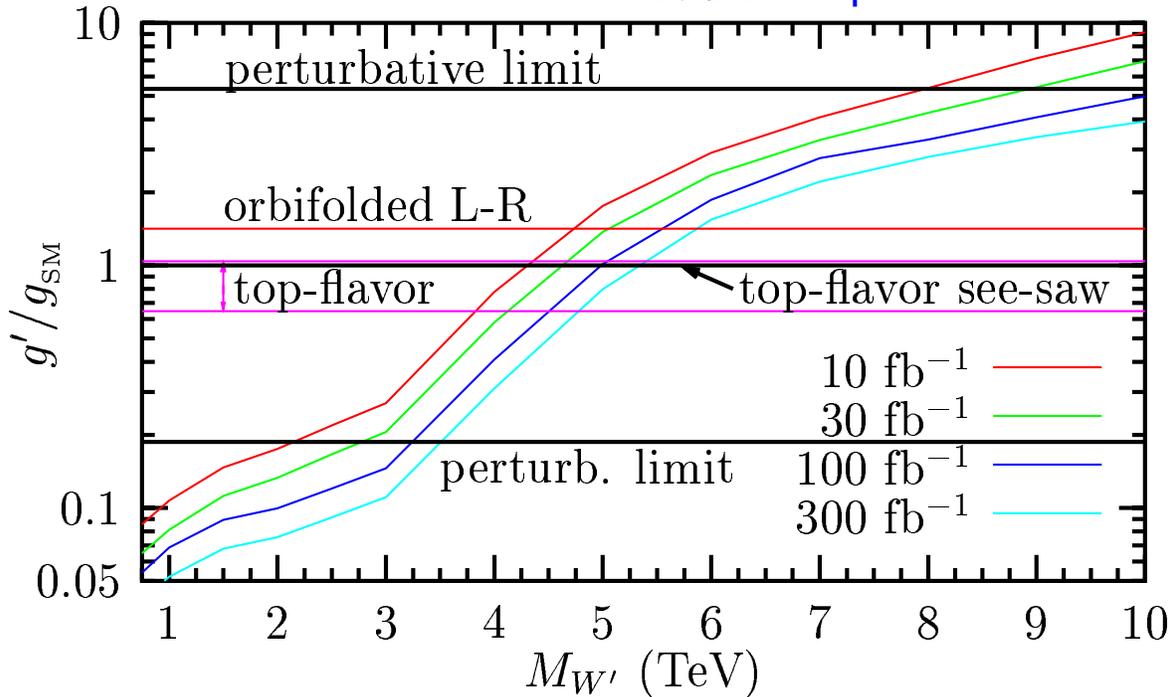
# 5 $\sigma$ discovery or 95% exclusion at the LHC

(using a Tevatron-like analysis)

$W'$ 's can be discovered between 3.7–4.7 TeV for SM-like ( $g'/g_{\text{SM}} \sim 1$ ) couplings.

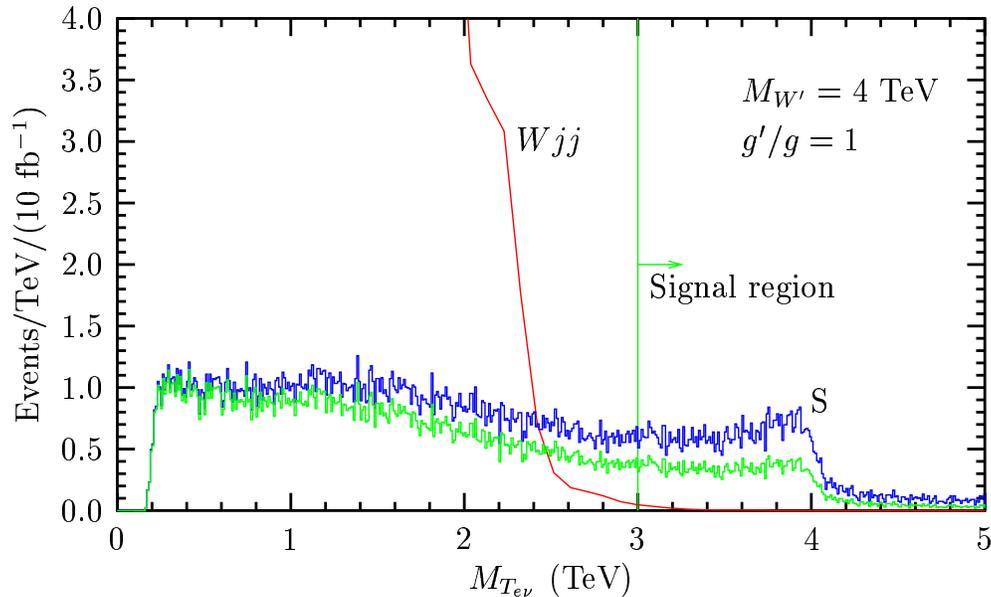


SM-like  $W'$ 's can be excluded at 95% C.L. up to 5.5 TeV!



## What does $W' \rightarrow e\nu$ look like?

Signal vs.  $Wjj$  background for  $M_{W'} = 4$  TeV



For a  $W'$  with SM-like couplings, we expect  $< 1$  event per low-luminosity year.

This is less than  $1/4$  the signal of the  $Wbj$  final state, and does not give a mass peak.

Only the  $Wjj$  background is shown. There are many other large backgrounds.

Can we improve on this?

# Questions for $W'$ physics and more

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- A large fraction of non-supersymmetric models will be strongly constrained by the search for  $W'$  bosons in the single-top-quark channel. All models are now simple to compare. [Z.S., PRD 66, 075011 \(2002\) \[hep-ph/0207290\]](#).
  - Reach: Tevatron  $\sim 900$  GeV, LHC  $\sim 750$  GeV–5.5 TeV. [Z.S., PRD 66, 075011 \(2002\); and hep-ph/0306266](#).
  - Factor of 2 improvements in the LHC cross section reach are attainable, but not fully studied. **Are there other ideas?**
  - **Is there a better analysis for the Tevatron?**
- The lepton final state will be very hard to observe at LHC. **Can we get around this?**
- More NLO calculations may be necessary to get kinematic distributions at the LHC. **What backgrounds are missing?**

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The following applies to both  $W'$  and  $Z'$  measurements.

- **What is the tagging rate for  $> 1$  TeV  $b$ -jets?**  
For the  $W'$  study, I assumed only the  $\sim 200$  GeV  $b$  from the top decay could be tagged. What about the recoil  $b$ ?
- **What is the  $E$  resolution and tagging rate for  $> 1$  TeV  $\mu$ ?**  
This effects  $W' \rightarrow \mu\nu$  and  $Z' \rightarrow \mu^+\mu^-$ . Do the muons bend enough? Do material interactions hide them?
- **What is the  $E$  resolution and tagging rate for  $> 1$  TeV  $e$ ?**  
They shower 100+ GeV photons. How well do you separate  $e-\gamma$ , or even  $e$ -jet? **Using a jet-shape variable  $\Rightarrow \epsilon_e$ ?**

There will be a broadly defined “Z-prime” workshop on Nov. 16 at Northwestern.  $W'$  issues will be addressed there too.