

Neutrino masses from double- β decay and kinematics experiments

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Last decade: the age of ν physics

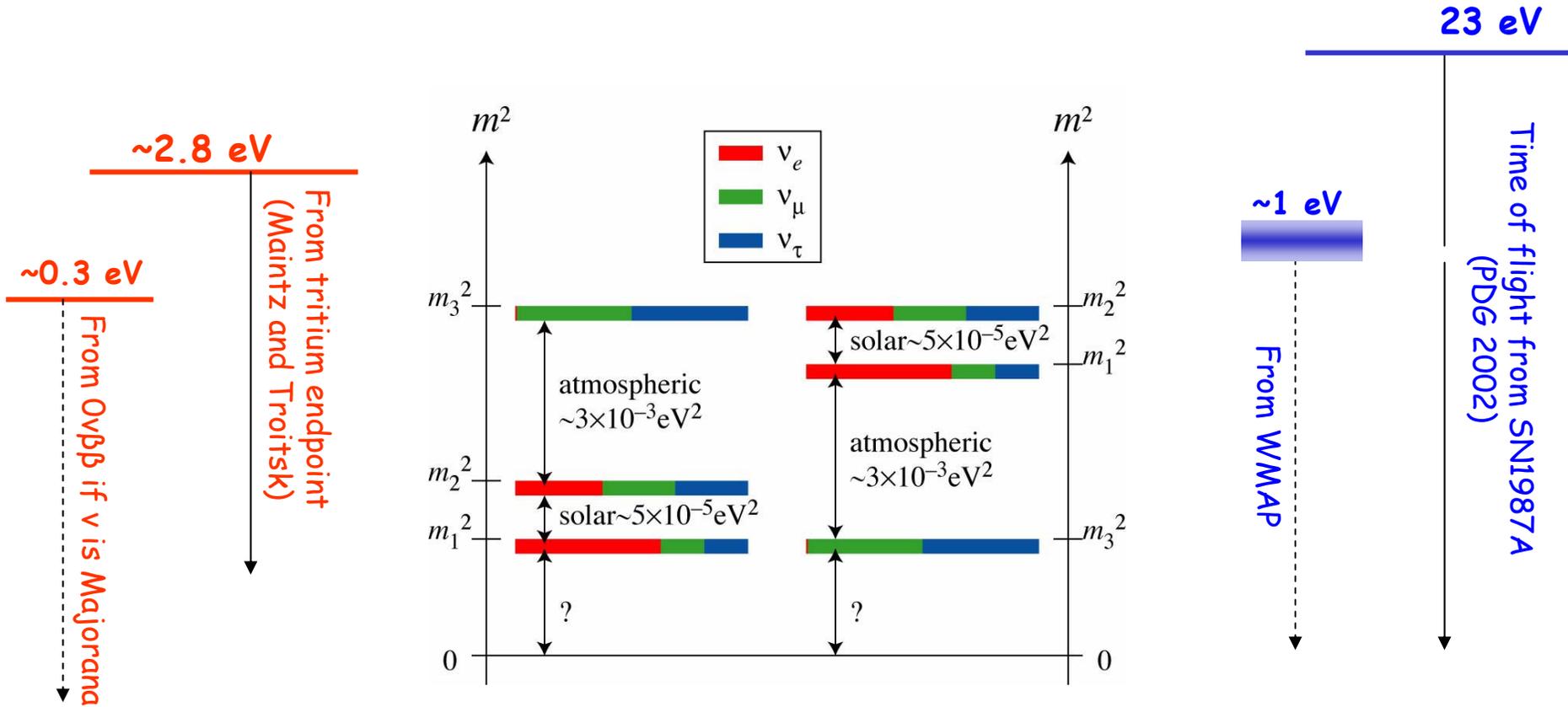
Discovery of ν flavor change

- *Solar neutrinos (MSW effect)*
- *Reactor neutrinos (vacuum oscillation)*
- *Atmospheric neutrinos (vacuum oscillation)*
- *K2K (vacuum oscillation)*
- *Lose ends: LSND/Karmen/miniBoone*

So, assuming miniBoone sees no oscillations,
we know that:

- ν masses are non-zero
- there are 2.981 ± 0.008 ν (Z lineshape)
- 3 ν flavors were active in Big Bang Nucleosynthesis

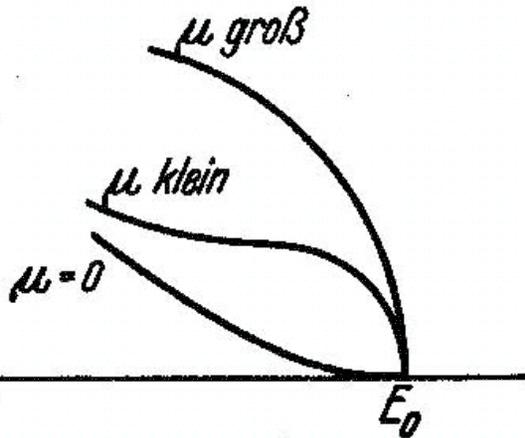
Yet, we still do not know: - the neutrino mass scale
- the choice of mass hierarchy



These *experimental* problems take a central place in the future of Particle Physics

Endpoint mass measurements

Study the spectral shape near the endpoint of a β decay (note that the end-point value is generally *not* known well enough to use its absolute position)



Principle almost as old as neutrino itself:
E. Fermi, Z. Phys. 88 (1934) 161

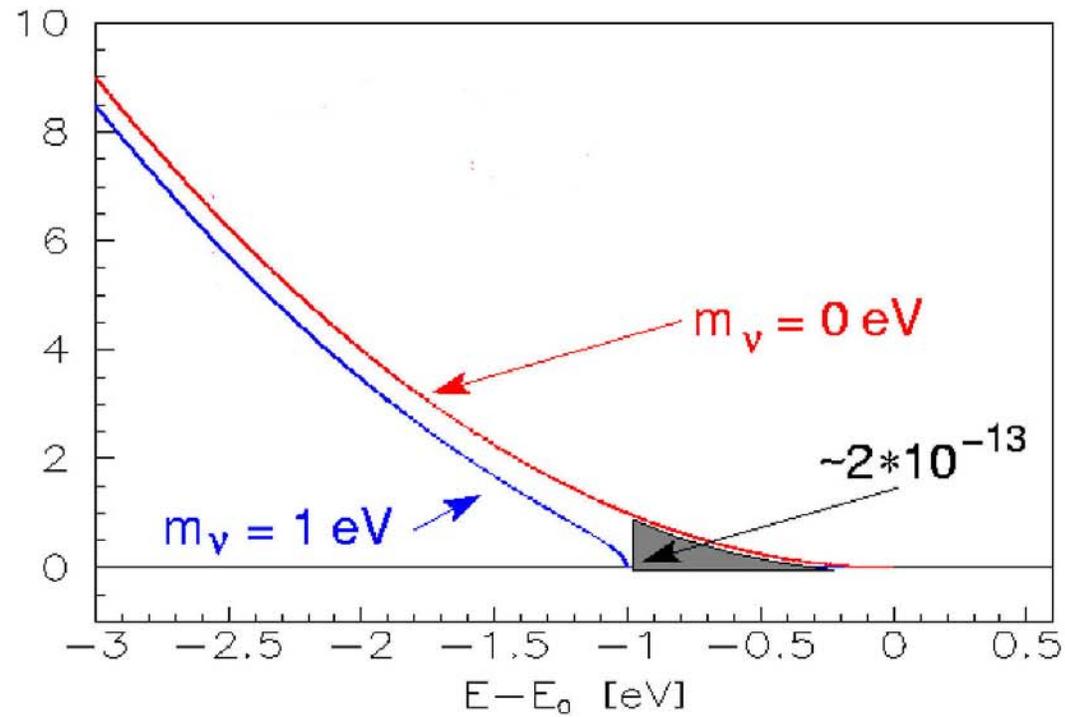
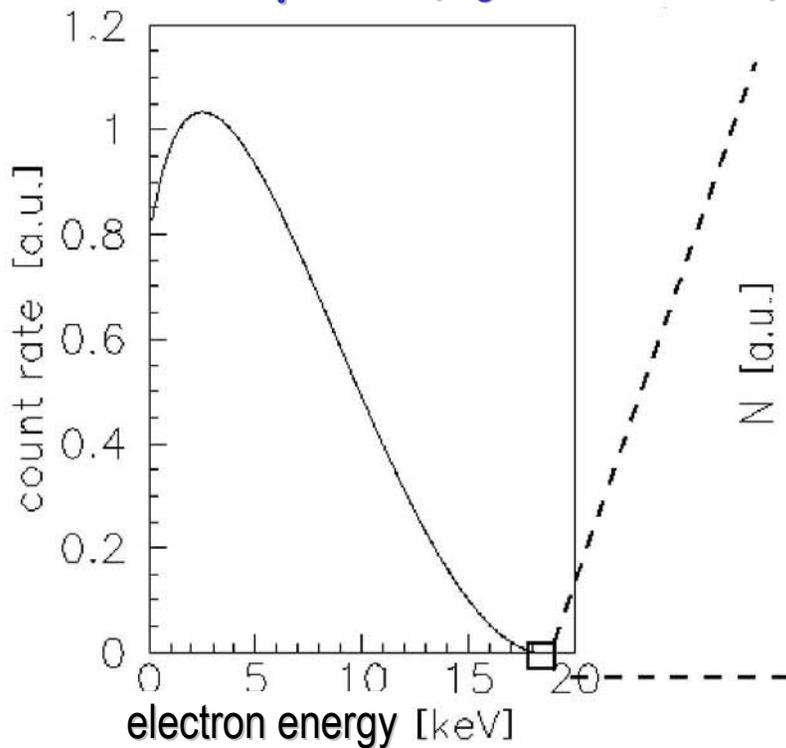
Measure the quantity:

$$m_{\nu_e}^{2(\text{eff})} = \sum_i |U_{ei}|^2 m_i^2$$

If the experimental resolution is smaller than $m_i^2 - m_j^2$ then one should see a separate kink in the spectrum for each of the states i and j

In modern experiments use mainly ${}^3_1T \rightarrow {}^3_2He + e^- + \bar{\nu}_e$

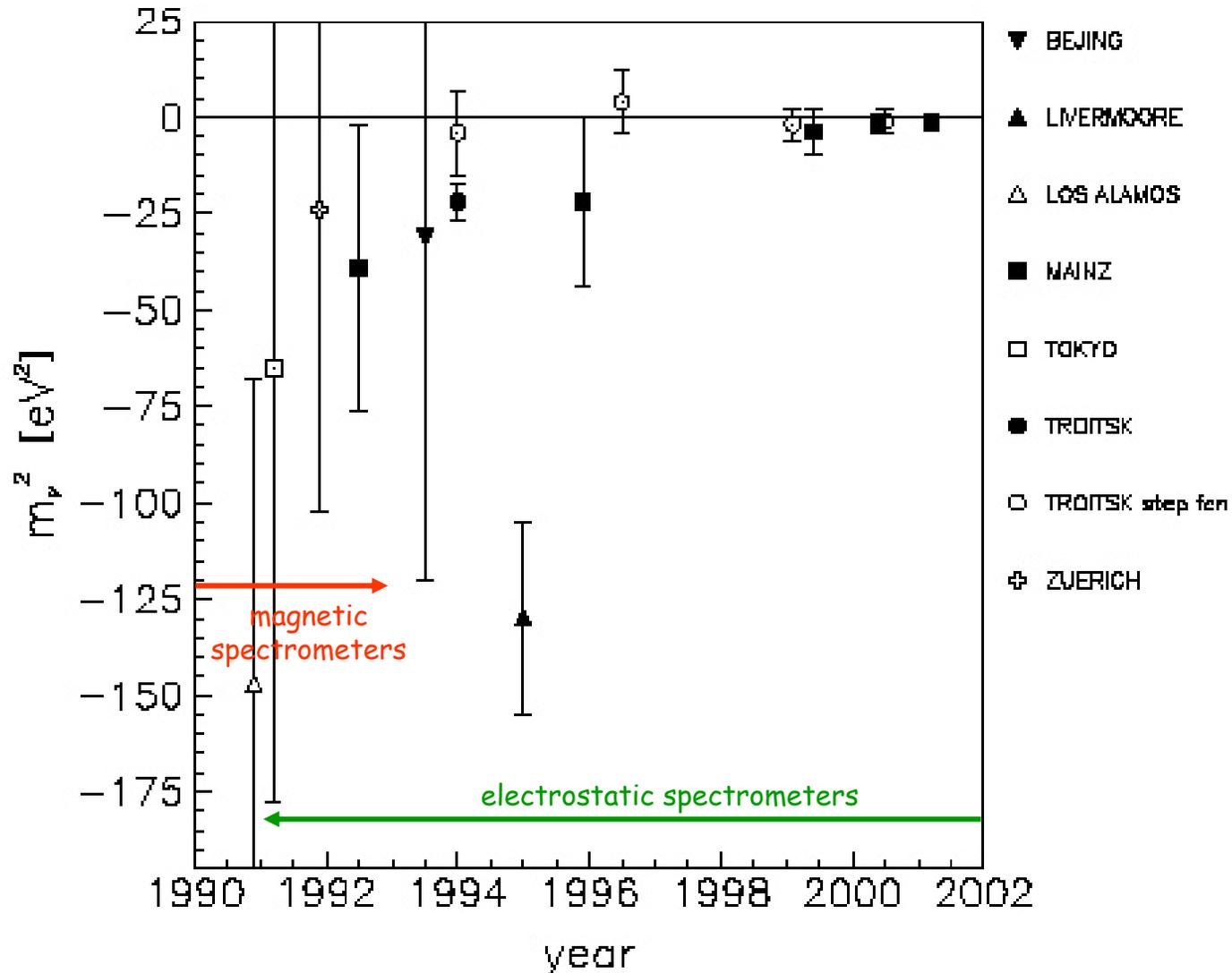
a super-allowed transition with rather good combination of low end point ($E_0=18.6$ keV) and short half life ($T_{1/2}=12.3$ yr)



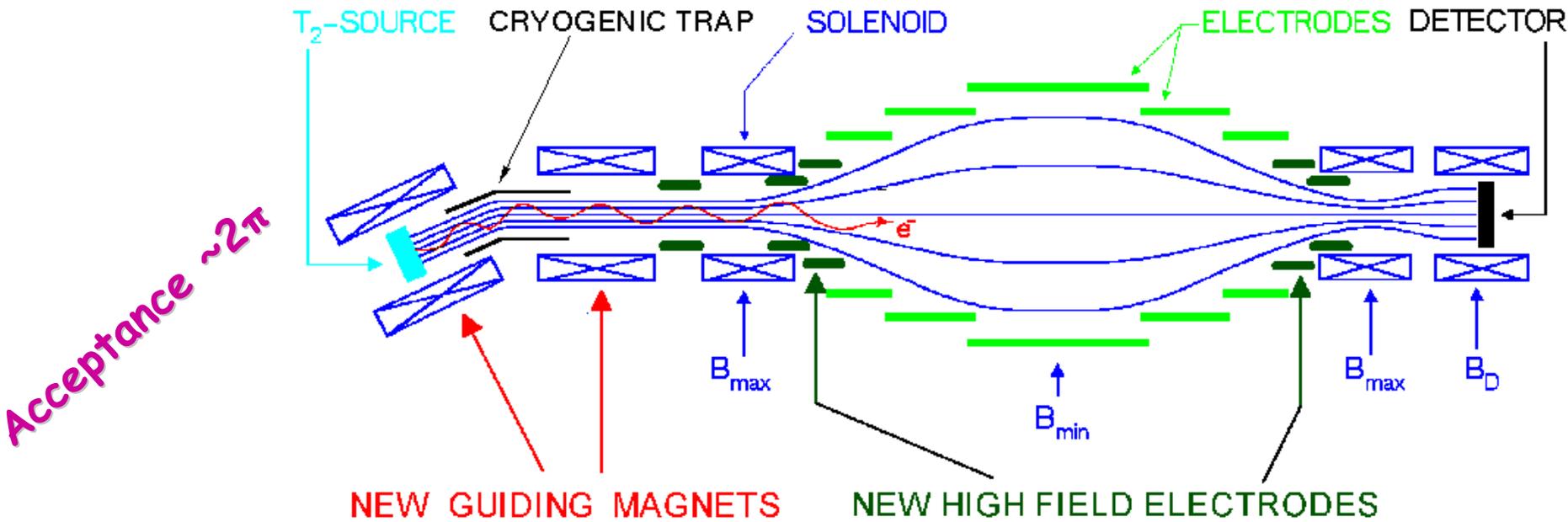
Spectrometer has to have

- 1) very high resolution
- 2) very high luminosity
(most of the statistics in the spectrum is not used...)

Long history of measurements that, for long time, have been plagued by negative central values for m_ν^2 (eff)



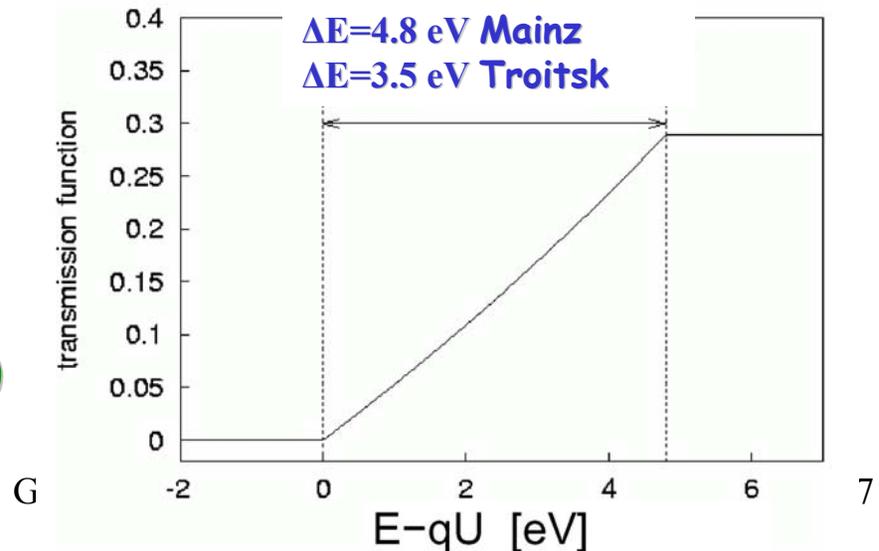
Recent experiments (Mainz and Troitsk) use Magnetic Adiabatic Collimation, Electrostatic Filter (MAC-E) integrating spectrometers



Sharp integrating transmission function with no tails

$$\Delta E/E = B_{\min}/B_{\max}$$

Low background (if vacuum good)



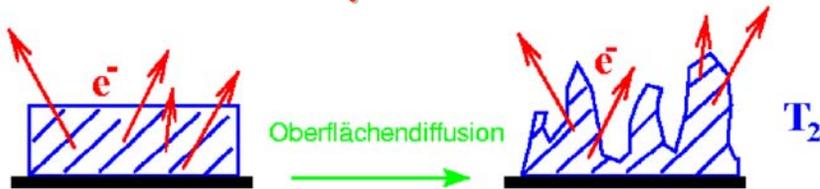
Main difference between exp:

- Mainz: solid (frozen) source
- Troitsk: windowless gaseous source

Any broadening of the spectrometer resolution reduces the apparent value of m^2

A. Saenz et al. Phys. Rev. Lett. 82 (2000) 242

Example: Mainz

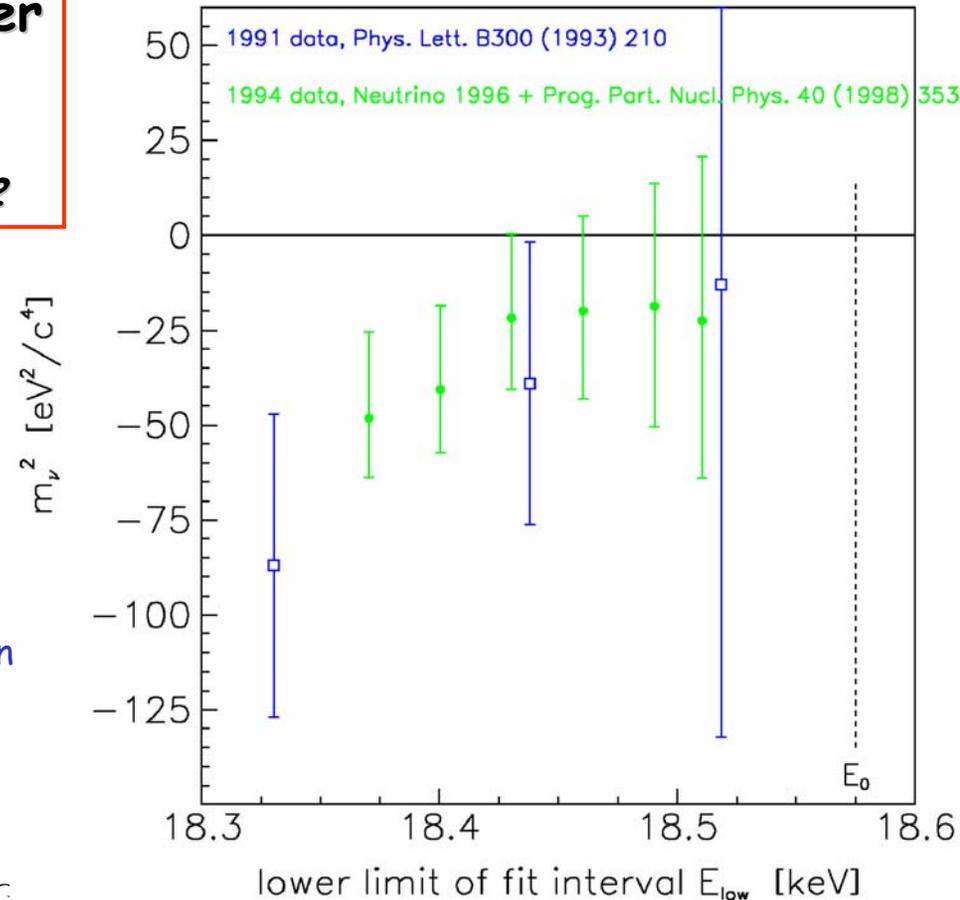


source produced as a thin, smooth layer

"roughening" transition occurs at finite temperature
→ change of energy loss function

At 2K the transition has a time constant of ~10yrs

...still imperfect modeling of the energy loss in the source was the origin of the early negative m^2 effects in all experiments



Mainz results:

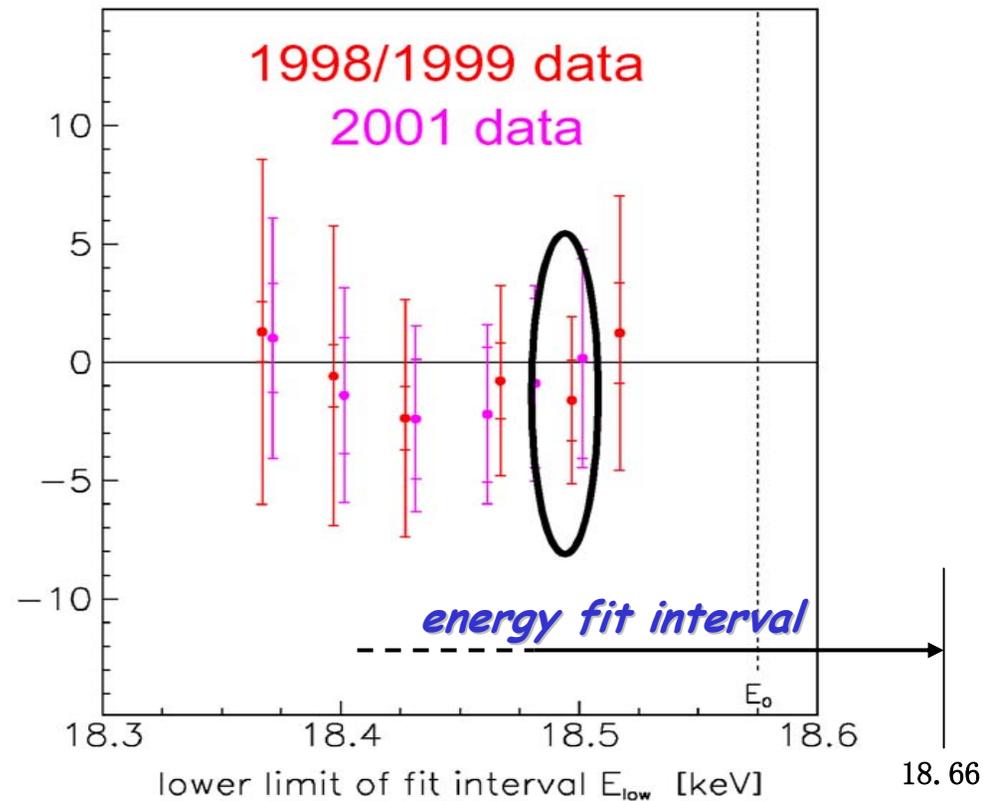
*Use fit range
18.500 - 18.666,
(other ranges give
consistent results)*

1998-99 runs

$$m_\nu^{2(\text{eff})} = -1.6 \pm 2.5 \pm 2.1 \text{ eV}^2/c^4$$

2001 runs

$$m_\nu^{2(\text{eff})} = +0.1 \pm 4.2 \pm 2.0 \text{ eV}^2/c^4$$



Together

$$m_\nu^{2(\text{eff})} = -1.2 \pm 2.2 \pm 2.1 \text{ eV}^2/c^4$$

$$m_\nu^{(\text{eff})} < 2.2 \text{ eV}/c^2 \quad (95\% \text{ CL})$$

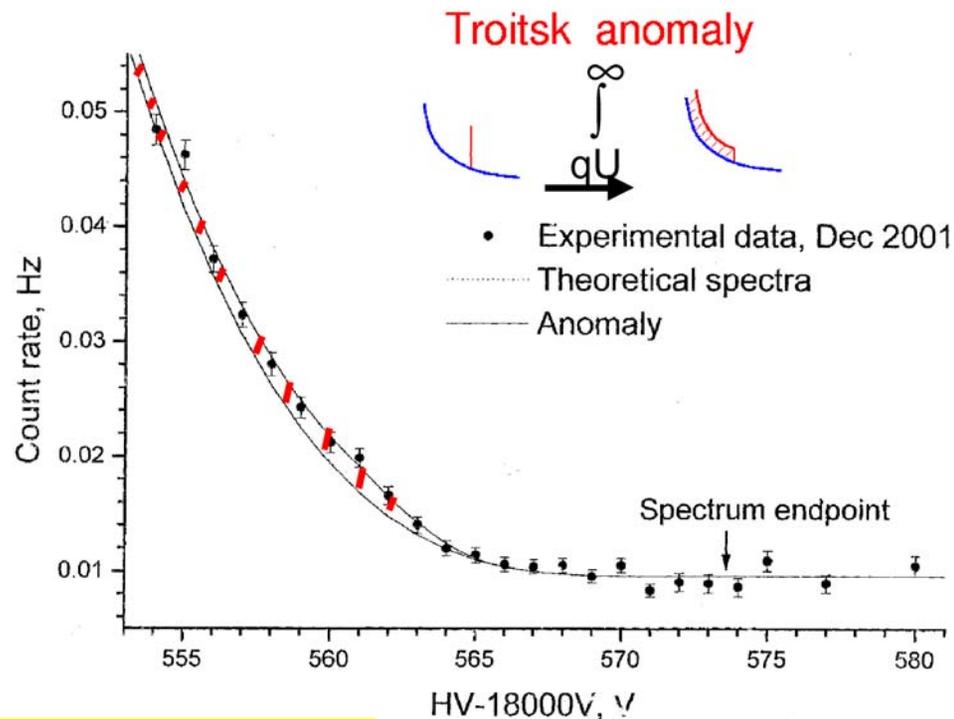
Ch. Kraus et al. Nucl. Phys. B 118 (2003) 482

Ch. Weinheimer Nucl. Phys. B 118 (2003) 279

Troitsk results

A step in the integral spectrum is found. \rightarrow this would imply that there is a line in the energy spectrum of tritium decay!

Position of the "line" seems to change from 0.5 eV to 15 eV with a 6 month period



Not well understood

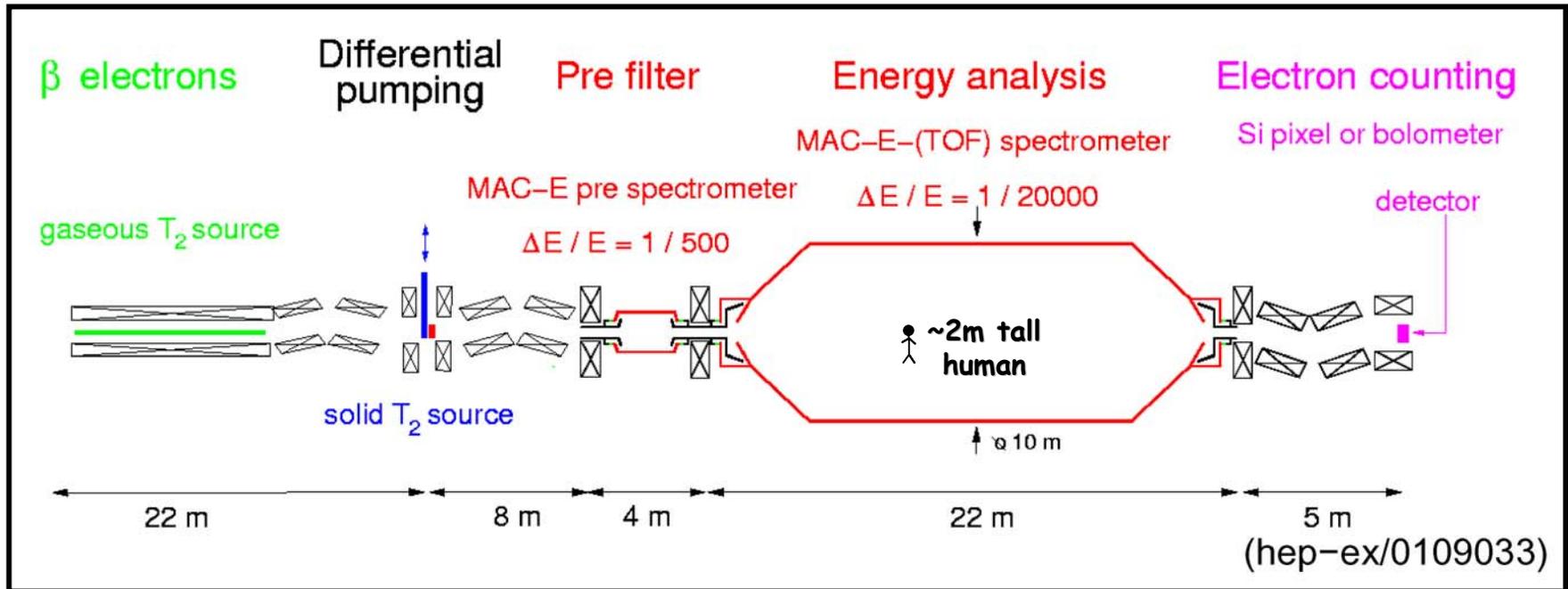
If one ignores the issue and adds a phenomenological peak to the fit (leaving the position free from period to period)

$$m_{\nu}^{2(\text{eff})} = -2.3 \pm 2.5 \pm 2.0 \text{ eV}^2 / c^4$$

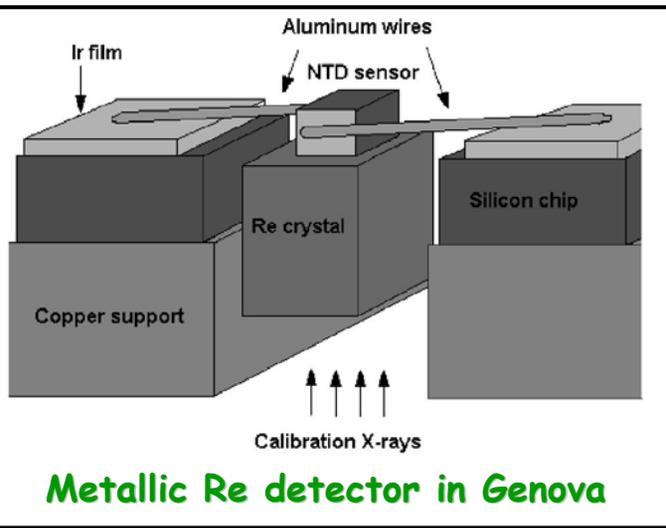
$$m_{\nu}^{(\text{eff})} < 2.05 \text{ eV} / c^2 \quad (95\% \text{ CL})$$

New, very large spectrometer being built in Karlsruhe for a better measurement with tritium: "KATRIN"

Forschungszentrum Karlsruhe (FZK), Universitat Mainz, INR (Troitsk), University of Washington (Seattle), University of Wales (Swansea), Nuclear Physics Institute (Rez/Prague), Fachhochschule Fulda, Universitat Karlsruhe, Universitat Bonn, JUNR (Dubna)



Expected sensitivity 0.20 - 0.25 eV
(assuming systematics are understood)



Calorimetric measurements:

Technique less mature and resolution worse but freedom to select β emitter:

- 1) calorimeter should be less sensitive to condensed matter effects
- 2) thin source (large specific activity or short $\frac{1}{2}$ life) not required

Genova

1.6 mg metallic Re crystal (1.1Bq)

$m_\nu^{(eff)} < 26 \text{ eV } 95\% \text{ CL}$

F. Gatti proceedings Neutrino 2000, p293

$^{187}\text{Re} \rightarrow ^{187}\text{Os} + e^- + \nu$

$E_0 = 2.5 \text{ keV}$ lowest end-point

$T_{1/2} = 4.1 \cdot 10^{10} \text{ yr}$

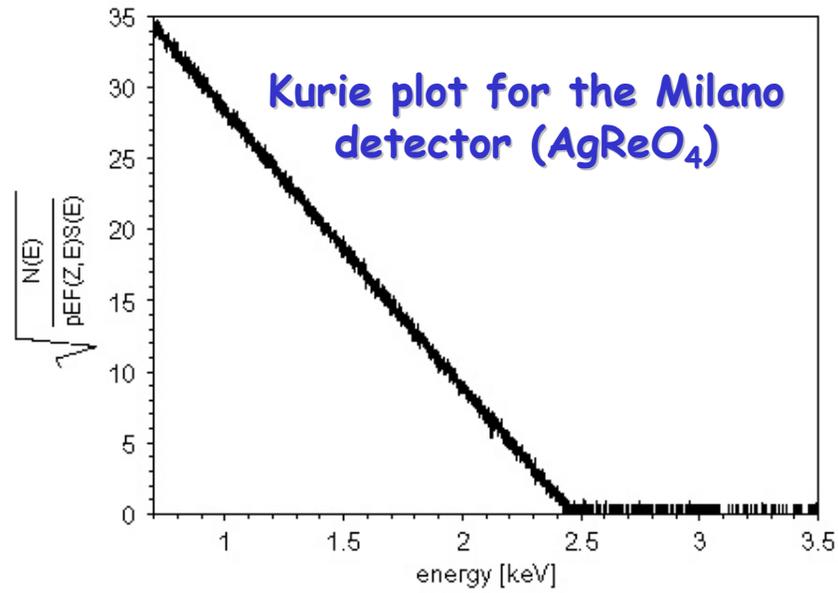
Milano

~10 · 30 μg AgReO_4 crystal

resolution 28 eV FWHM

$m_\nu^{(eff)} < 21.7 \text{ eV } 90\% \text{ CL}$

C. Arnaboldi et al. hep-ex/0302006



Kinematics mass measurements at "high energy"

• $M(\nu_\mu) < 0.19 \text{ MeV}/c^2$ 90% CL

from $\pi^+ \rightarrow \mu^+ \nu$ decays at rest

(K. Assamagan et al. PRD 53 (1996) 6065 + PDG 2002)

BNL E952 proposal expects $\sim 8 \text{ keV}$ sensitivity

• $M(\nu_\tau) < 18.2 \text{ MeV}/c^2$ 95% CL from τ decays
in ALEPH (R. Barate et al. EPJ C2 (1998) 395)

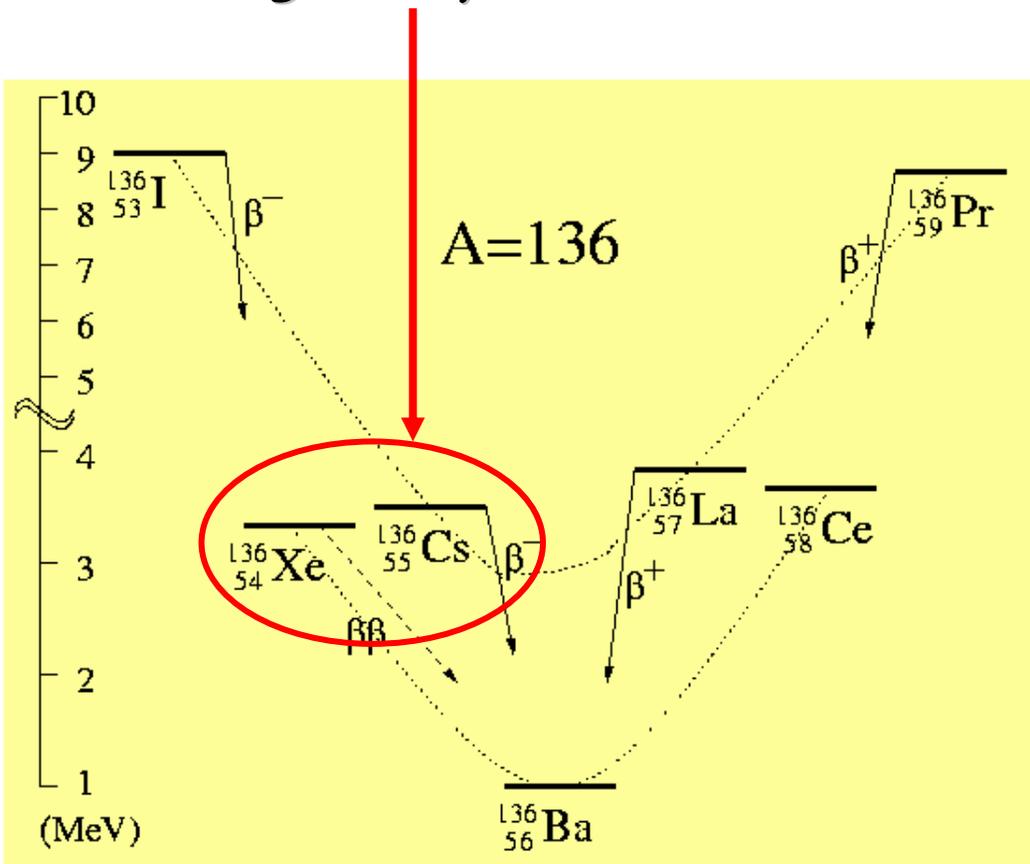
$M_{\nu_\tau} < 15.5 \text{ MeV}/c^2$ 95% CL from combined fit to $\Upsilon(4s)$ and Z^0 data
(J.M. Roney, Neutrino 2000, Sudbury)

$\sim 3 \text{ MeV}$ seems the asymptotic sensitivity of B factories

*Unlikely to reach
the "interesting" region below 1 eV*

Double-beta decay:

*a second-order process
only detectable if first
order beta decay is
energetically forbidden*



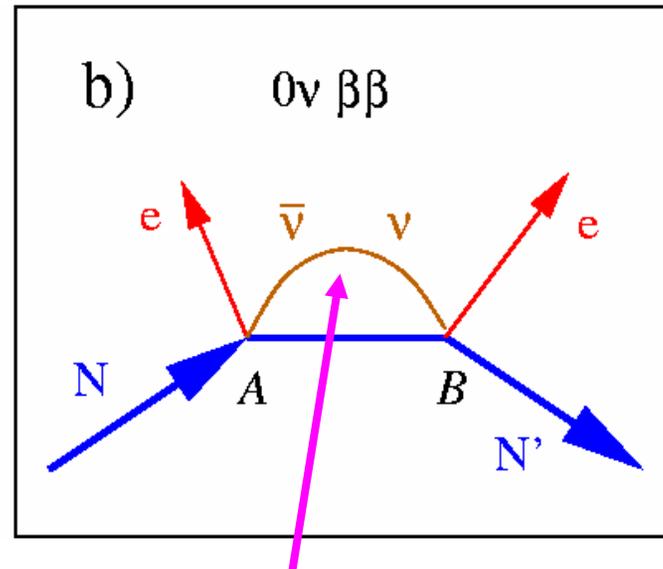
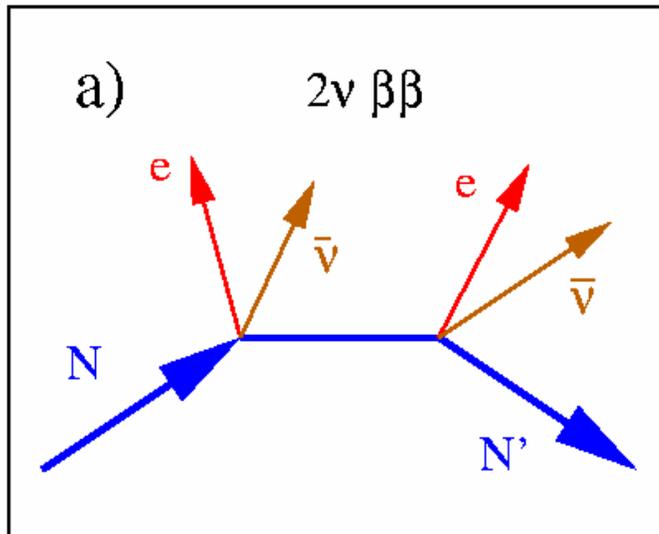
Candidate nuclei with $Q > 2$ MeV

Candidate	Q (MeV)	Abund. (%)
$^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$	4.271	0.187
$^{76}\text{Ge} \rightarrow ^{76}\text{Se}$	2.040	7.8
$^{82}\text{Se} \rightarrow ^{82}\text{Kr}$	2.995	9.2
$^{96}\text{Zr} \rightarrow ^{96}\text{Mo}$	3.350	2.8
$^{100}\text{Mo} \rightarrow ^{100}\text{Ru}$	3.034	9.6
$^{110}\text{Pd} \rightarrow ^{110}\text{Cd}$	2.013	11.8
$^{116}\text{Cd} \rightarrow ^{116}\text{Sn}$	2.802	7.5
$^{124}\text{Sn} \rightarrow ^{124}\text{Te}$	2.228	5.64
$^{130}\text{Te} \rightarrow ^{130}\text{Xe}$	2.533	34.5
$^{136}\text{Xe} \rightarrow ^{136}\text{Ba}$	2.479	8.9
$^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$	3.367	5.6

There are two varieties of $\beta\beta$ decay

2ν mode: a conventional
 2^{nd} order process
in nuclear physics

0ν mode: a hypothetical
process can happen
only if: $\bullet M_\nu \neq 0$ Since helicity
has to "flip"
 $\bullet \mathbf{v} = \bar{\mathbf{v}}$

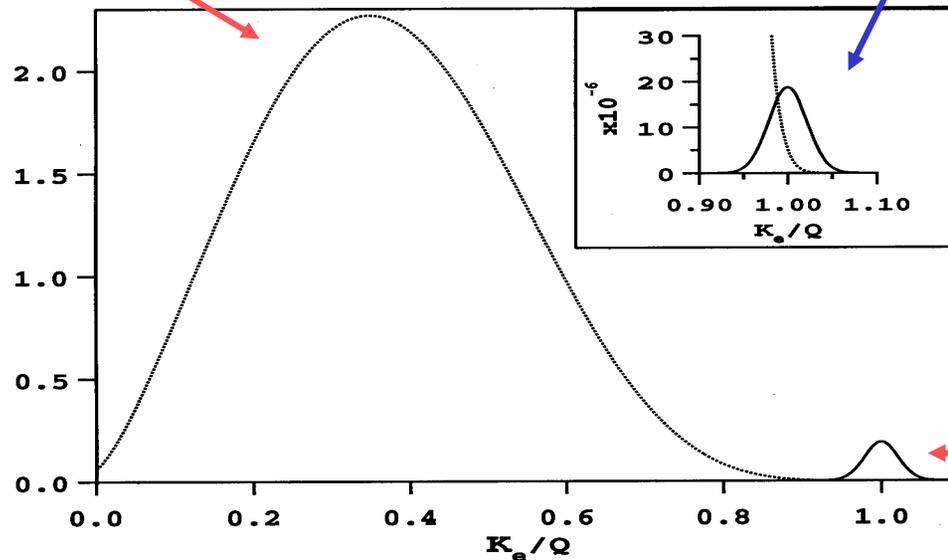


Several new particles can take
the place of the virtual ν
But $0\nu\beta\beta$ decay always implies new physics

Background due to the Standard Model $2\nu\beta\beta$ decay

$2\nu\beta\beta$ spectrum
(normalized to 1)

$0\nu\beta\beta$ peak (5% FWHM)
(normalized to 10^{-6})



$0\nu\beta\beta$ peak (5% FWHM)
(normalized to 10^{-2})

Summed electron energy in units of the kinematic endpoint (Q)

from S.R. Elliott and P. Vogel, Ann.Rev.Nucl.Part.Sci. 52 (2002) 115.

The only effective tool here is energy resolution

$\beta\beta$ decay experiments are at the leading edge of "low background" techniques

- Final state ID: 1) "Geochemical": search for an abnormal abundance of $(A, Z+2)$ in a material containing (A, Z)
2) "Radiochemical": store in a mine some material (A, Z) and after some time try to find $(A, Z+2)$ in it
 - + Very specific signature
 - + Large live times (particularly for 1)
 - + Large masses
 - Possible only for a few isotopes (in the case of 1)
 - No distinction between 0ν , 2ν or other modes
- "Real time": ionization or scintillation is detected in the decay
 - a) "Homogeneous": source=detector
 - b) "Heterogeneous": source \neq detector
 - + Energy/some tracking available (can distinguish modes)
 - + In principle universal (b)
 - Many γ backgrounds can fake signature
 - Exposure is limited by human patience

Real time is needed to discover ν masses, final state ID would be a nice complement !

The Standard Model
 $2\nu\beta\beta$ decay has been
 observed in many isotopes

Isotope	$T_{1/2}^{2\nu}$ (yr)
^{48}Ca	$(4.3 \pm 2.2) \cdot 10^{19}$
^{76}Ge	$(1.77 \pm 0.12) \cdot 10^{21}$
^{82}Se	$(8.3 \pm 1.2) \cdot 10^{19}$
$^{96}\text{Zr}^\dagger$	$(9.4 \pm 3.2) \cdot 10^{18}$ § $(2.1 \pm 0.6) \cdot 10^{19}$ $(3.9 \pm 0.9) \cdot 10^{19}$ §
^{100}Mo	$(9.5 \pm 1.0) \cdot 10^{18}$
^{116}Cd	$(2.6 \pm 0.6) \cdot 10^{19}$
^{128}Te	$(7.2 \pm 0.4) \cdot 10^{24}$ §
$^{130}\text{Te}^\dagger$	$(2.7 \pm 0.1) \cdot 10^{21}$ § $(7.9 \pm 1.0) \cdot 10^{20}$ § $(6.1 \pm 3.5) \cdot 10^{20}$
$^{136}\text{Xe}^\$$	$> 1.1 \cdot 10^{22}$ 90% CL
^{150}Nd	$(6.7 \pm 0.8) \cdot 10^{18}$
^{238}U	$(2.0 \pm 0.6) \cdot 10^{21}$ *

Table *arbitrarily simplified*
 from PDG 2003

†Results not in good agreement
 §Geochemical experiment
 *Radiochemical experiment
 \$Decay NOT observed, lower
 limit reported

If $0\nu\beta\beta$ is due to light ν Majorana masses

$$\langle m_\nu \rangle^2 = \left(T_{1/2}^{0\nu\beta\beta} G^{0\nu\beta\beta}(E_0, Z) \left| M_{GT}^{0\nu\beta\beta} - \frac{g_V^2}{g_A^2} M_F^{0\nu\beta\beta} \right|^2 \right)^{-1}$$

$$M_F^{0\nu\beta\beta} \text{ and } M_{GT}^{0\nu\beta\beta}$$

can be calculated within particular nuclear models

$$G^{0\nu\beta\beta}$$

a known phase space factor

$$T_{1/2}^{0\nu\beta\beta}$$

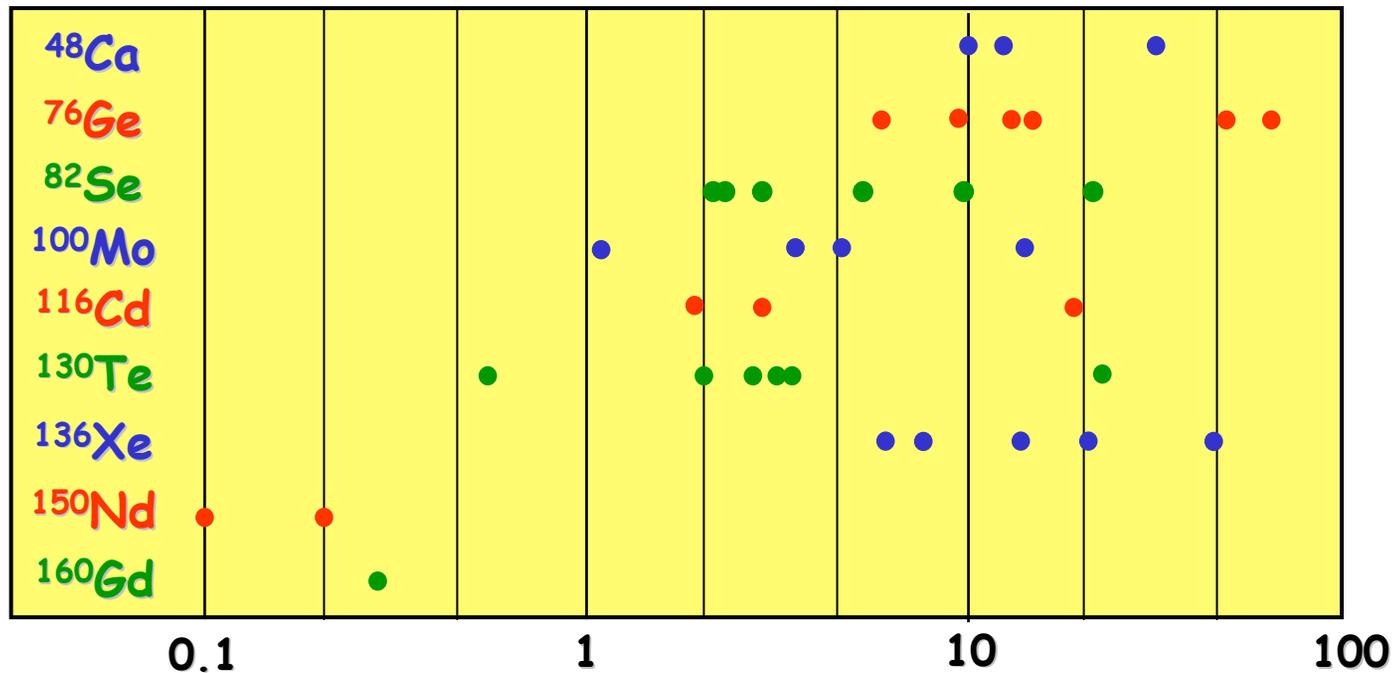
is the quantity to be measured

$$\langle m_\nu \rangle = \left| \sum_{i=1}^3 U_{e,i}^2 m_i \varepsilon_i \right|$$

effective Majorana ν mass
($\varepsilon_i = \pm 1$ if CP is conserved)

Cancellations are possible...

$0\nu\beta\beta$ decay half lives in 10^{26} yr units for $\langle m_\nu \rangle = 50$ meV according to different nuclear matrix element calculations



*[adapted from S.R.Elliott & P.Vogel
Ann. Rev. Nucl. Part. Sci. 52 (2002) 115]*

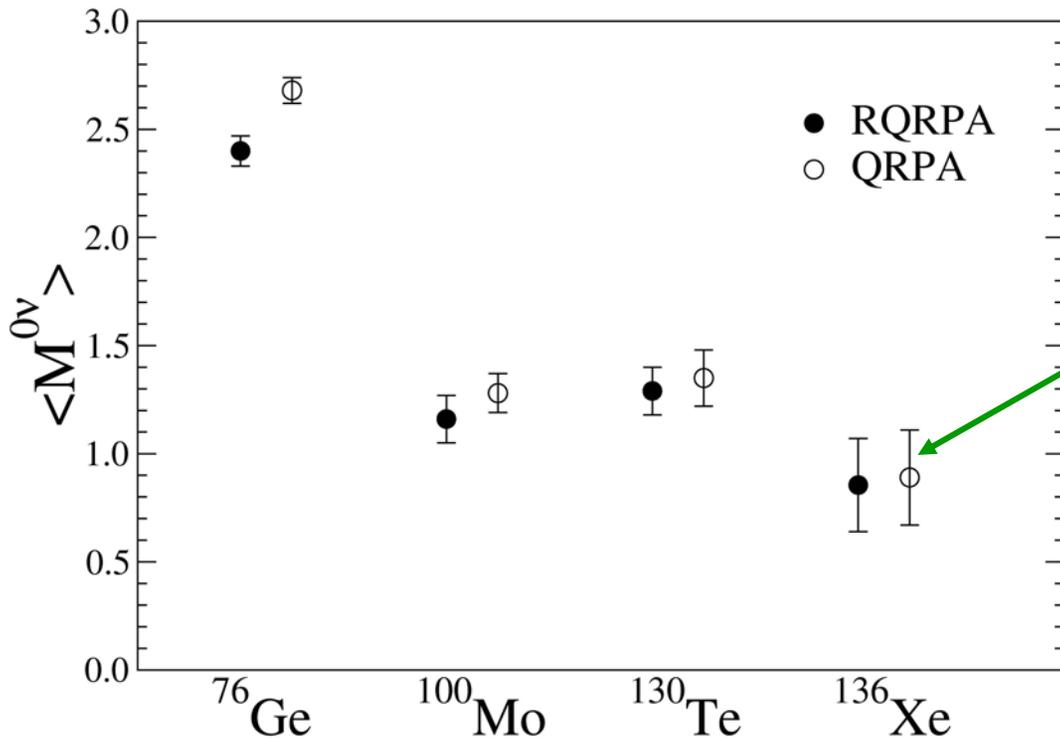
Unfortunately it is not trivial to use the 2ν matrix element to normalize the 0ν one:

- $|M_{2\nu}|$ - has stronger dependence on intermediate states
- $|M_{0\nu}|$ - all multipoles contribute
 - ν propagator results in long range potential

However it was recently found that main uncertainty in (R)QRPA calculations comes from the single particle space around the Fermi surface. This should be the same for $0\nu\beta\beta$ and for $2\nu\beta\beta$.

Use the measured $2\nu\beta\beta$ experimental $T_{1/2}$ to make a correction.

V.A.Rodin et al. nucl-th/0305005



Lower bound on $T_{1/2}$
used for ^{136}Xe

Can one get agreement
from Nuclear Shell
Models ?

Still, if/once $0\nu\beta\beta$ decay is discovered, the $T_{1/2}$ in more than one nucleus will be needed to pin down neutrino masses

Present Limits for 0ν double beta decay

Candidate nucleus	Detector type	(kg yr)	Present $T_{1/2}^{0\nu\beta\beta}$ (yr)	$\langle m \rangle$ (eV)
^{48}Ca	Ge diode	~30	$>9.5 \cdot 10^{21}$ (76%CL)	$<0.39^{+0.17}_{-0.28}$
^{76}Ge			$>1.9 \cdot 10^{25}$ (90%CL)	
^{82}Se			$>9.5 \cdot 10^{21}$ (90%CL)	
^{100}Mo			$>5.5 \cdot 10^{22}$ (90%CL)	
^{116}Cd			$>7.0 \cdot 10^{22}$ (90%CL)	
^{128}Te	TeO ₂ cryo	~3	$>1.1 \cdot 10^{23}$ (90%CL)	$<1.1 - 2.6$
^{130}Te	TeO ₂ cryo	~3	$>2.1 \cdot 10^{23}$ (90%CL)	
^{136}Xe	Xe scint	~10	$>1.2 \cdot 10^{24}$ (90%CL)	
^{150}Nd			$>1.2 \cdot 10^{21}$ (90%CL)	
^{160}Gd			$>1.3 \cdot 10^{21}$ (90%CL)	

Adapted from the Particle Data Group 2003

Has $0\nu\beta\beta$ decay been already discovered ??

EVIDENCE FOR NEUTRINOLESS DOUBLE BETA DECAY

H.V. KLAPDOR-KLEINGROTHAUS^{1,3},
A. DIETZ¹, H.L. HARNEY¹, I.V. KRIVOSHEINA^{1,2}

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Germany

²Radiophysical-Research Institute, Nishnii-Novgorod, Russia

³Spokesman of the GENIUS and HEIDELBERG-MOSCOW Collaborations,

(Part of the Heidelberg-Moscow collaboration)

Mod. Phys Lett. A27 (2001) 2409

...most likely not

...see details in

C.A.Aalseth Mod. Phys. Lett. A17 (2002) 1475

F.Feruglio et al. Nucl.Phys. B637 (2002) 345-377

Addendum-ibid. B659 (2003) 359-362

Yu.Zdesenko et al. Phys.Lett. B 546 (2002) 206

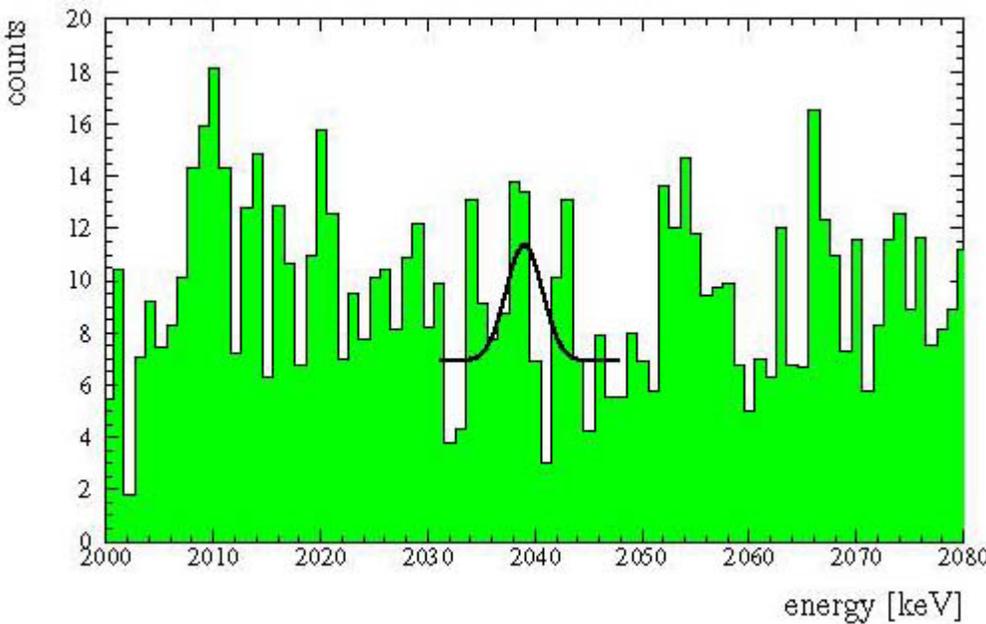
H.L.Harney Mod.Phys.Lett. A16 (2001) 2409

H.V.Klapdor-Kleingrouthaus hep-ph/0205228

*A.M.Bakalyarov et al. ("Moscow" of Heidelberg-Moscow) to appear
in proceedings of NANP 2003, June 2003, Dubna, Russia*

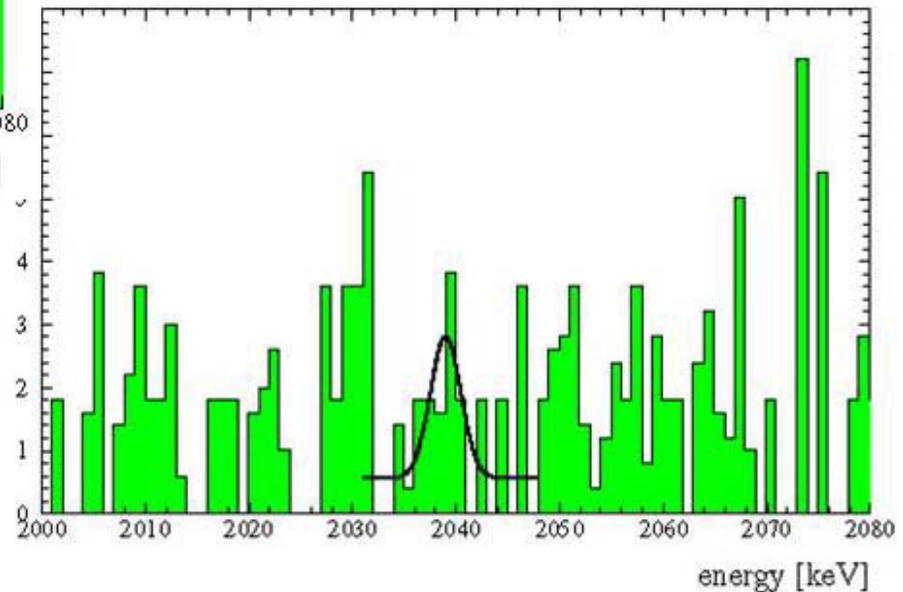
Paper's bottom line is $T_{1/2} = [0.8 - 18.3] \cdot 10^{25}$ yr at 95% CL
 best value is $T_{1/2} = 1.5 \cdot 10^{25}$ yr corresponding to 0.39 eV

Allegedly this is a 2 to 3 sigma effect depending on the analysis



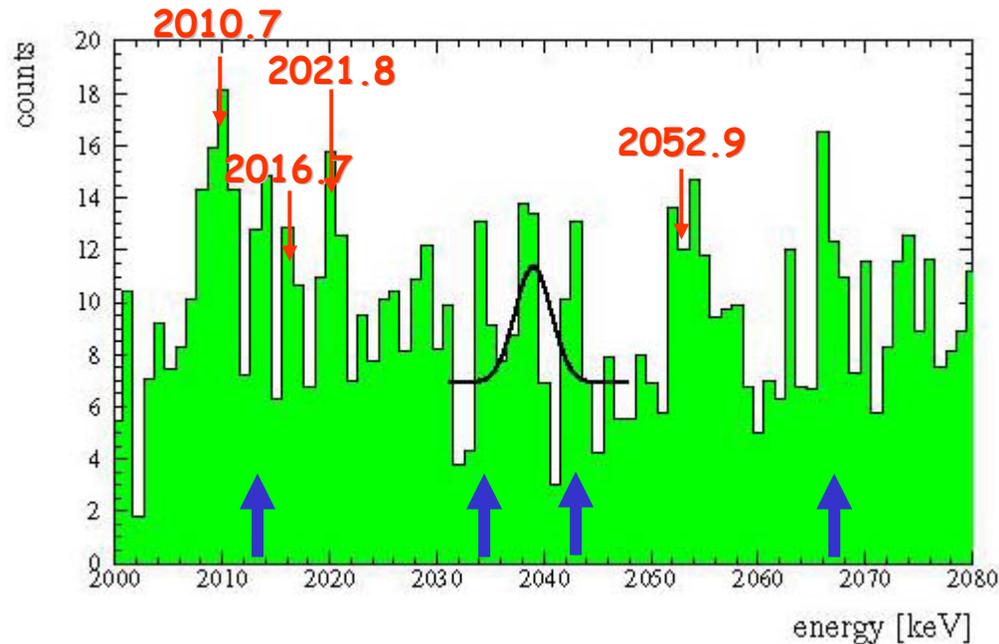
"Evidence" from the search of a peak in the energy spectrum observed in a set of low activity Ge detectors inside the Gran Sasso Lab.

Spectrum can be somewhat cleaned-up by applying pulshape discrimination to remove γ ray events ...still lots of peaks besides the 2039 keV



The fit to the "signal" peak at 2039.006 keV is done *AFTER* the subtraction of 4 peaks that are claimed to be *UNDERSTOOD* background from *IDENTIFIED* lines of ^{214}Bi

Without this subtraction the significance of the 2039 peak is even less than 2 sigma, as it is evident by just staring at the spectrum



But then, what about the other peaks !
There are more that are not understood !

Note that the data used is the same that was earlier interpreted as an upper limit
 $T_{1/2} > 1.9 \cdot 10^{25}$ eV at 90% CL

"The claim of discovery...is considered critically and firm conclusion about, at least, prematurely of such claim is derived on the basis of simple statistical analysis..."

Yu.Zdesenko et al. Phys Lett B546 (2002) 206

The latest 2 experiments to start operation:

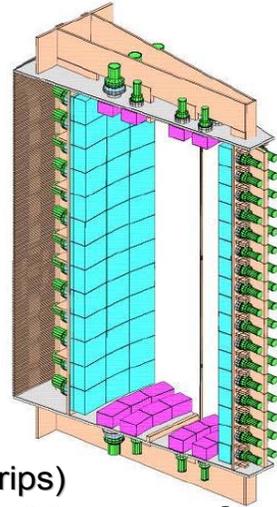
NEMO III

Neutrinoless Experiment with MOlibdenum III or Neutrino Ettore Majorana Observatory

Large Collaboration: 13 groups from Europe, USA and Japan

Passive source - Spectroscopic approach

$0\nu 2\beta$ sensitivity:
 $T \sim 10^{24}$ y
 $\langle m\nu \rangle \sim 0.1$ eV



Detector structure: 20 sectors

1 Source:

up to 10 kg of $\beta\beta$ isotopes
(metal film or powder glued to mylar strips)
cylindrical surface: $20 \text{ m}^2 \times 40\text{-}60 \text{ mg/cm}^2$

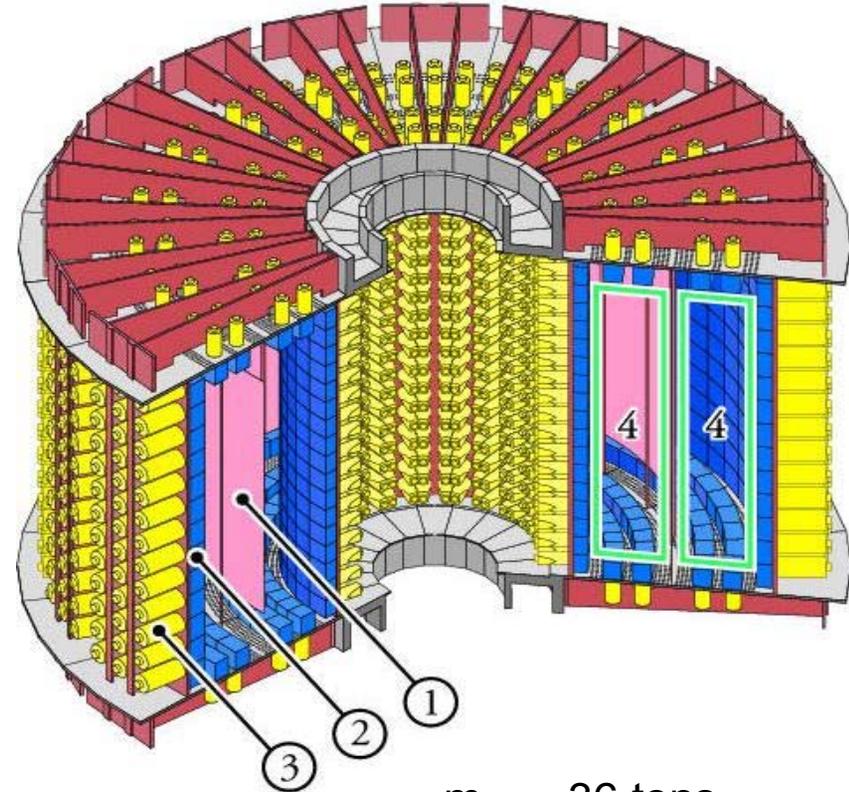
2 Tracking volume:

open octagonal drift cells (6180)
operated in Geiger mode
($\sigma_r = 0.5 \text{ mm}$, $\sigma_z = 1 \text{ cm}$)

3 Calorimeter:

1940 plastic scintillators coupled to low activity PMs:
FWHM(1 MeV) $\sim 11\text{-}14.5 \%$

Magnetic Field (30 G) + Iron Shield (20 cm) + Neutron Shield (30 cm H_2O)



$m_{\text{tot}} \sim 36$ tons
Low activity materials

"Cuoricino" (small CUORE)

Mostly natural TeO_2

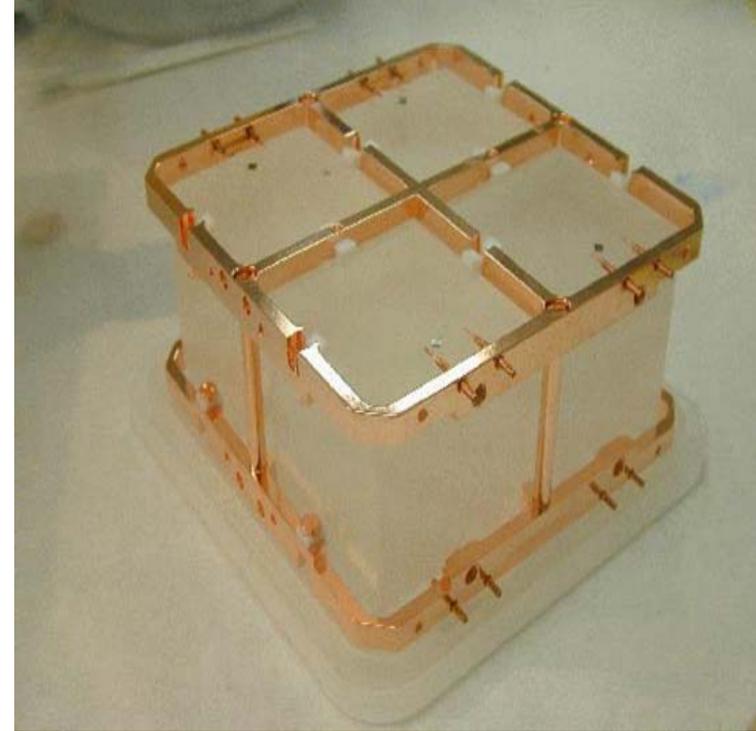
44 $(5 \times 5 \times 5) \text{ cm}^3$ crystals ($44 \times 780\text{g}$)

18 $(3 \times 3 \times 6) \text{ cm}^3$ crystals ($18 \times 340\text{g}$)

Total $\sim 40 \text{ kg}$

Tower structure prototype for much larger CUORE

Running at Gran Sasso in a dilution refrigerator at $\sim 10 \text{ mK}$



NTD thermistor readout:
 $1 \text{ MeV} = \Delta T = 300 \mu\text{V}$

$0\nu\beta\beta$ sensitivity
 $T_{1/2} \sim 4 \times 10^{23} \text{ yr}$
 $\langle m_\nu \rangle \sim 0.7 - 1.6 \text{ eV}$

A (probably incomplete) list of the different ideas discussed by various groups

Experiment	Nucleus	Detector	$T^{0\nu}$ (y)	$\langle m_\nu \rangle$ eV
CUORE	^{130}Te	.77 t of TeO_2 bolometers (nat)	7×10^{26}	.014-.091
EXO	^{136}Xe	10 t Xe TPC + Ba tagging	1×10^{28}	.013-.037
GENIUS	^{76}Ge	1 t Ge diodes in LN	1×10^{28}	.013-.050
Majorana	^{76}Ge	1 t Ge diodes	4×10^{27}	.021-.070
MOON	^{100}Mo	34 t nat.Mo sheets/plastic sc.	1×10^{27}	.014-.057
DCBA	^{150}Nd	20 kg Nd-tracking	2×10^{25}	.035-.055
CAMEO	^{116}Cd	1 t CdWO_4 in liquid scintillator	$> 10^{26}$.053-.24
COBRA	^{116}Cd , ^{130}Te	10 kg of CdTe semiconductors	1×10^{24}	.5-2.
Candles	^{48}Ca	Tons of CaF_2 in liq. scint.	1×10^{26}	.15-.26
GSO	^{116}Cd	2 t $\text{Gd}_2\text{SiO}_5:\text{Ce}$ scint in liq scint	2×10^{26}	.038-.172
Xmass	^{136}Xe	1 t of liquid Xe	3×10^{26}	.086-.252

Note that the sensitivity numbers are somewhat arbitrary, as they depend on the author's guesstimate of the background levels they will achieve



*Alabama, Caltech, Colorado State,
Irvine, ITEP, Neuchatel, Stanford collaboration*

An exotic approach to deal with the main experimental problems

1. To reach $\langle m_\nu \rangle \sim 10$ meV very large fiducial mass (tons) (except for Te) need massive isotopic enrichment
2. Reduce and control backgrounds in qualitatively new ways
bkgnd for Ge ~ 0.3 ev/kg yr FWHM

For no bkgnd

$$\langle m_\nu \rangle \propto 1 / \sqrt{T_{1/2}^{0\nu\beta\beta}} \propto 1 / \sqrt{Nt}$$

Scaling with bkgd
goes like Nt

$$\langle m_\nu \rangle \propto 1 / \sqrt{T_{1/2}^{0\nu\beta\beta}} \propto 1 / (Nt)^{1/4}$$

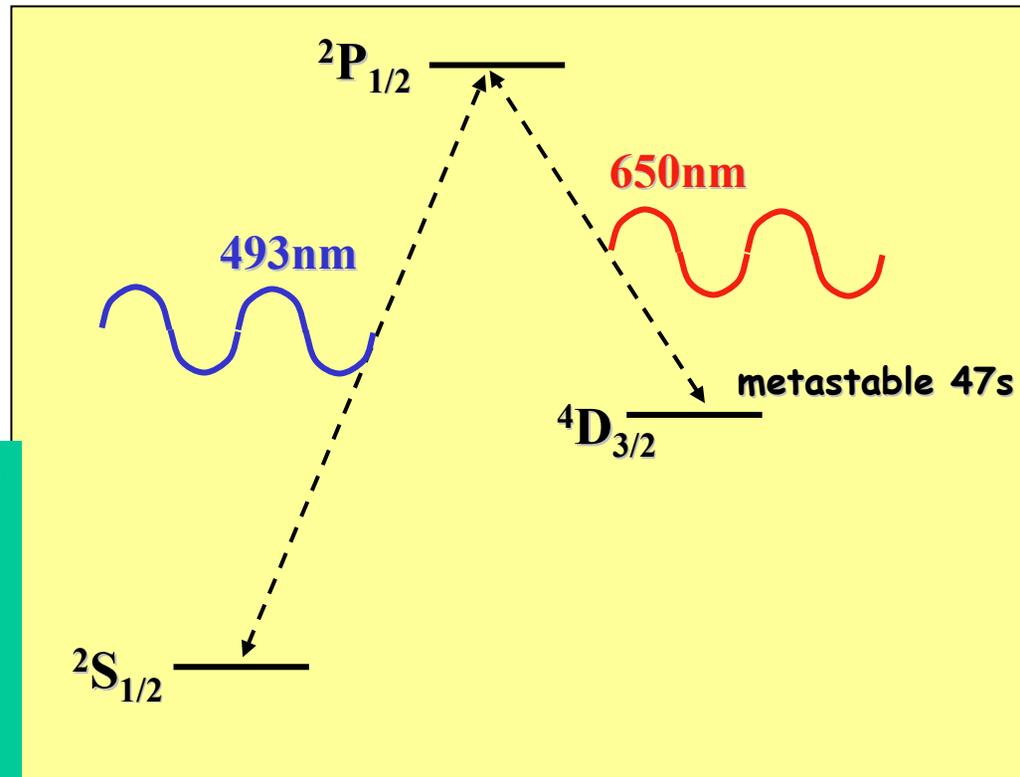
In addition would like a multi-parameter experiment,
→ possible discovery can be backed-up by cross checks with more than one single variable

Xe offers a qualitatively new tool against background:
 $^{136}\text{Xe} \rightarrow ^{136}\text{Ba}^{++} e^- e^-$ final state can be identified
using optical spectroscopy (M.Moe PRC44 (1991) 931)

Ba⁺ system best studied
(Neuhauser, Hohenstatt,
Toshek, Dehmelt 1980)
Very specific signature
"shelving"

Single ions can be detected
from a photon rate of $10^7/\text{s}$

- Important additional constraint
- Huge background reduction



The Ba-tagging, added to a conventional Xe TPC rejection power provides the tools to develop a background-free next-generation $\beta\beta$ experiment

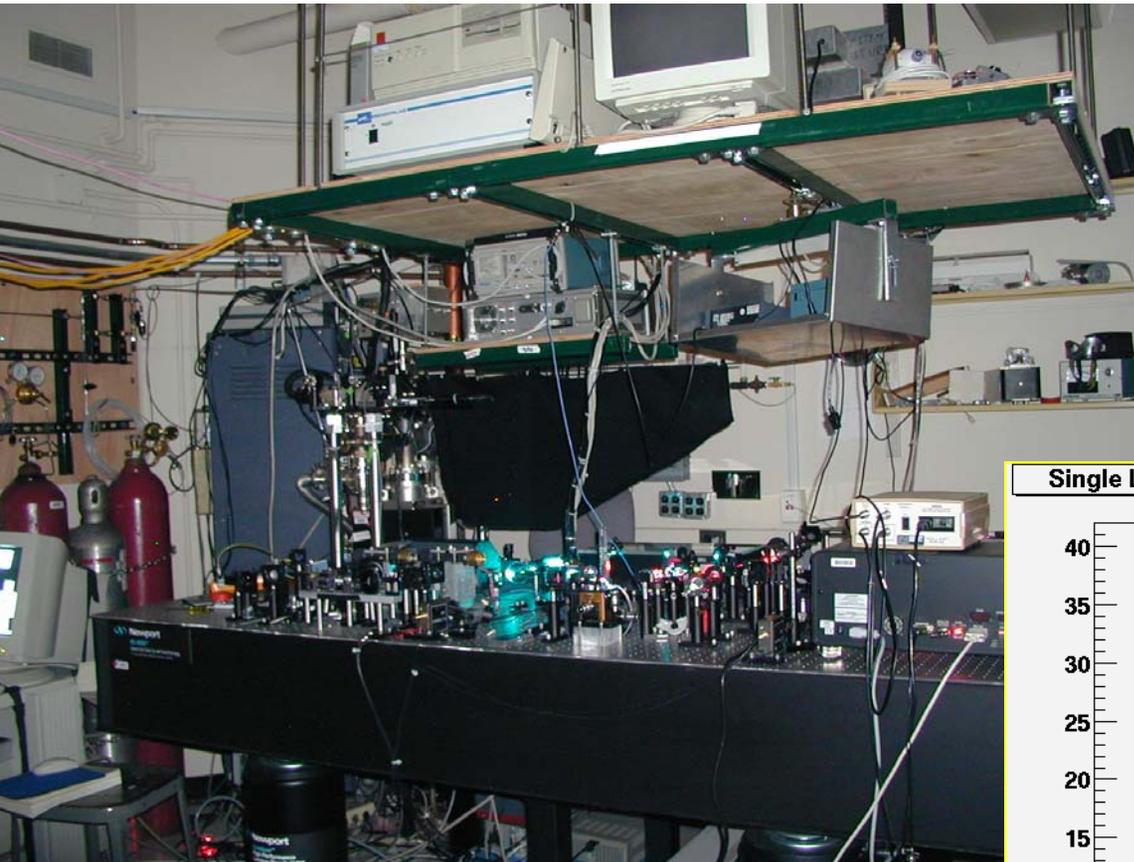
Assume an “asymptotic” fiducial mass of 10 tons of ^{136}Xe at 80%

R&D program focused on:

- Single Ba^+ tagging in Xe background
- Energy resolution in xenon (liquid and gas)
- Transfer of single Ba ions out of LXe
- 200kg prototype detector construction (no Ba tagging yet) to study detector performance, backgrounds and measure $2\nu\beta\beta$ mode
- Isotopic enrichment of large quantities of ^{136}Xe

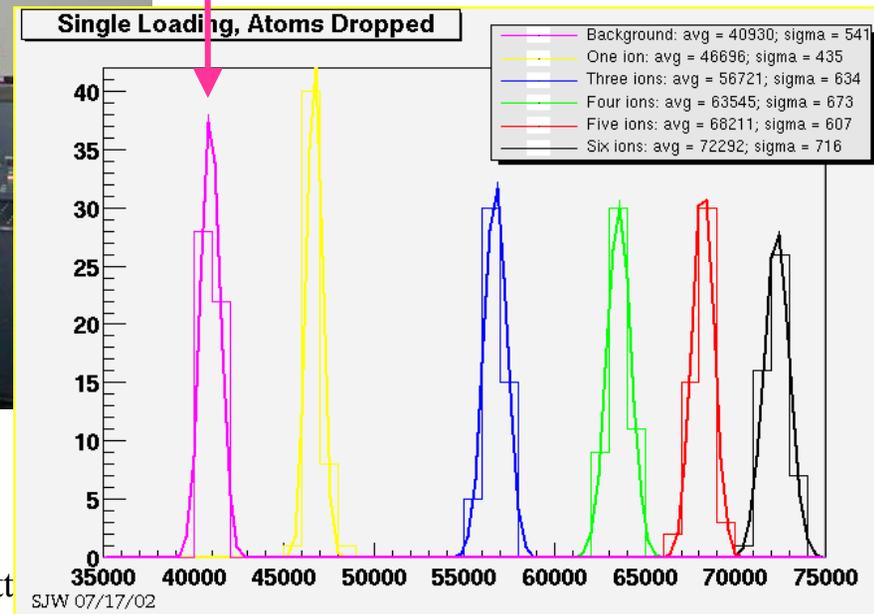
*Already have in hand 200kg of enriched Xe (80% 136 isotope)
→ the largest stockpile of highly enriched isotope ever produced for pure science !*

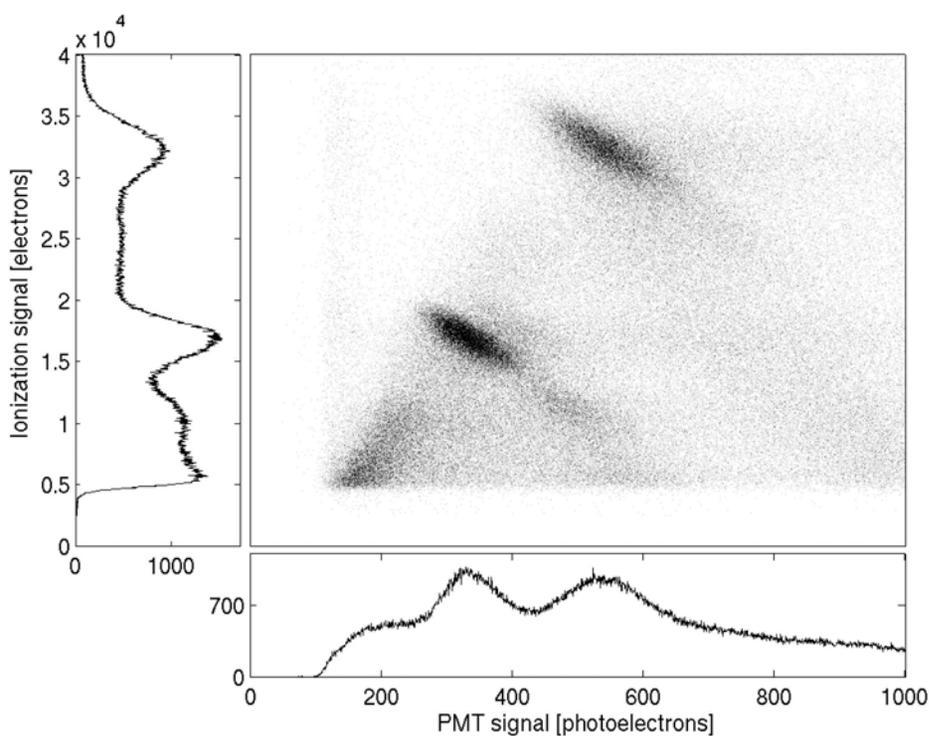
Laser spectroscopy R&D



CCD image of a Ba⁺ ion in vacuum

Zero ion background

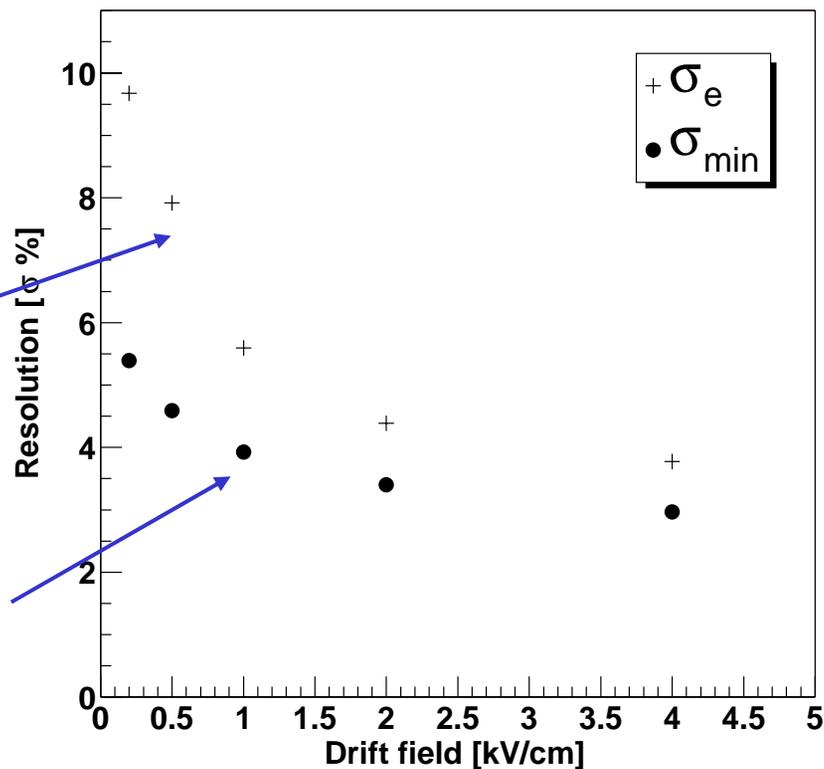




Sufficient improvement in energy resolution already demonstrated using anti-correlation between scintillation and ionization ($\sigma=2\%$ @ 2.5 MeV)

Ionization only

Ionization combined with scintillation



Conclusions

- Welcome to the era of massive neutrinos !
- After 75 years of neutrinos we now know that neutrinos are massive
- For the first time there is a good chance that the mass scale and the Dirac/Majorana structure of the neutrino sector will be measured in the lab
- Lots of fun physics and interesting techniques !