



Advanced Accelerator Concepts

ICFA

Jonathan Wurtele, UCB/LBNL

October 8, 1999

- Overview
- High Gradient Acceleration at 100GHz and Beyond
- Role of Plasma in Advanced Accelerators
 - ➔ Accelerating Structures
 - ➔ Focusing Elements
- Laser Acceleration
- Perspectives



Overview

- Employing NLC technology for high energy ($> 3\text{ TeV}$) colliders leads to extremely large machines and average powers. Some alternatives are being investigated:

Compact linear colliders $\rightarrow$ high-gradient $\rightarrow$ problems:

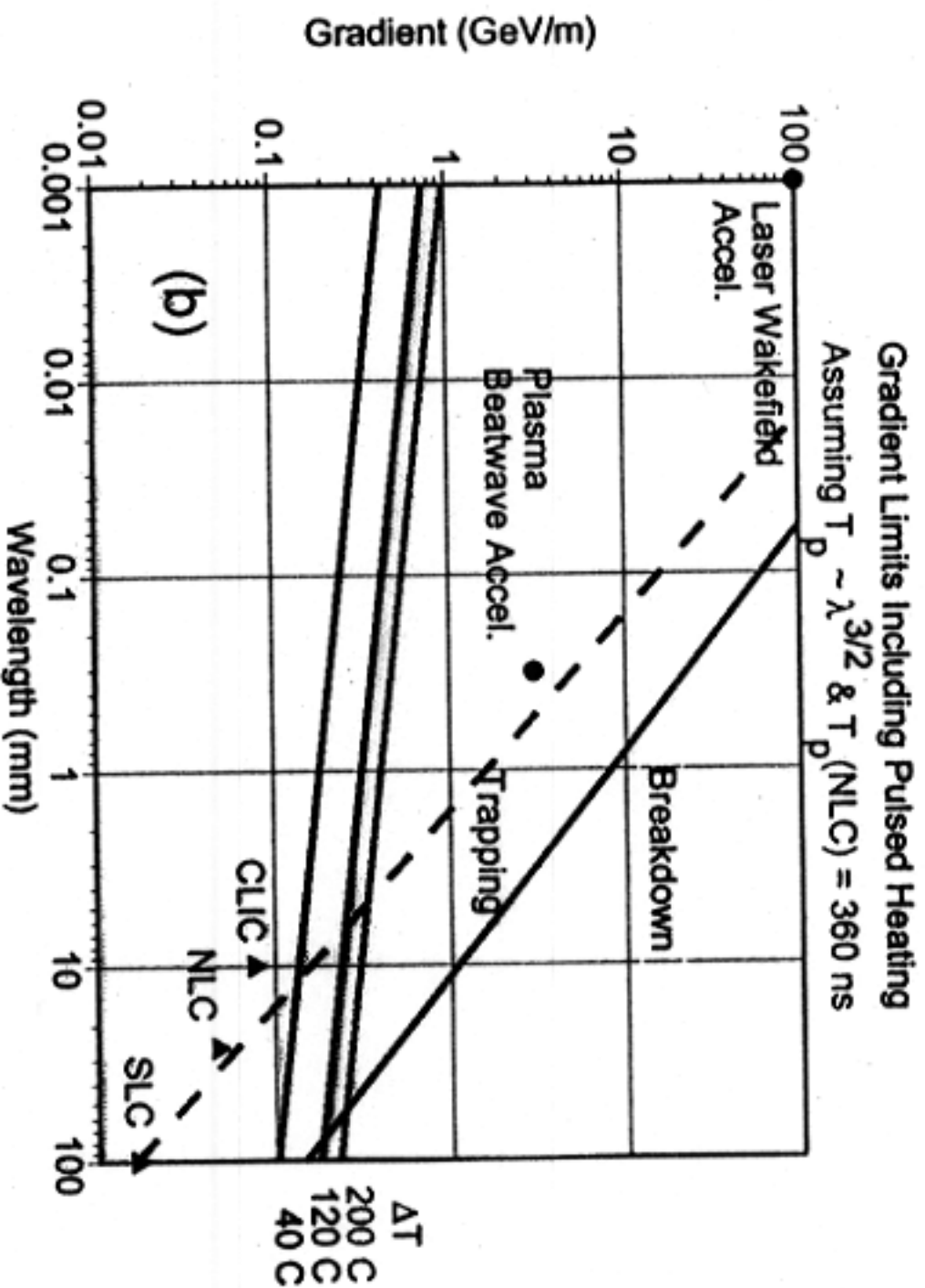
- $\rightarrow$ New structures : fabrication, breakdown, wakefields, jitter, alignment
- $\rightarrow$ New power sources: efficiency, pulse length, coupling
- $\rightarrow$ New final focus : pinch effect, beamstrahlung, length, & backgrounds

luminosity for sensible emittance

AND

Muon colliders $\rightarrow$ muons $\rightarrow$ problems:

- $\rightarrow$ Production
- $\rightarrow$ Cooling
- $\rightarrow$ Collision
- $\rightarrow$ Decay



Pulse heating is thought to be the most serious
limit to scaled solid-state structures

(Courtesy D. Whittum)

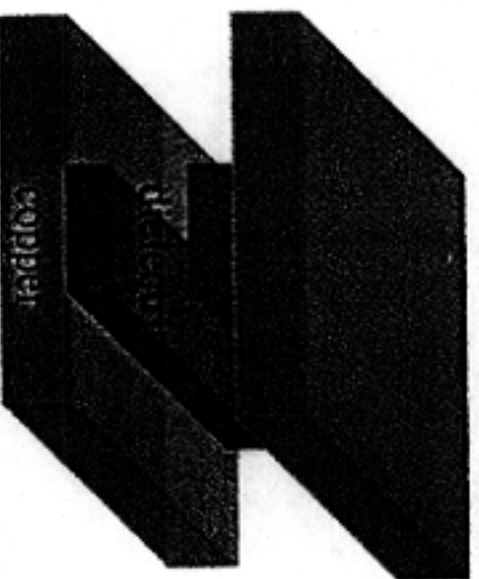


Composite Accelerator Structure

The traditional disk loaded waveguide, like those used at SLAC, use the periodic iris to slow down the phase velocity. The Composite Accelerator Structure (CAS) utilizes a dielectric to slow down the phase velocity to c. M. Hill, et al., submitted to IEEE/MTT.

Advantages:

- The structure is simple to manufacture--no bonding, tolerant to machining errors.
- The planar geometry also makes tuning easier.
- CAS has inherently lower surface electric field.



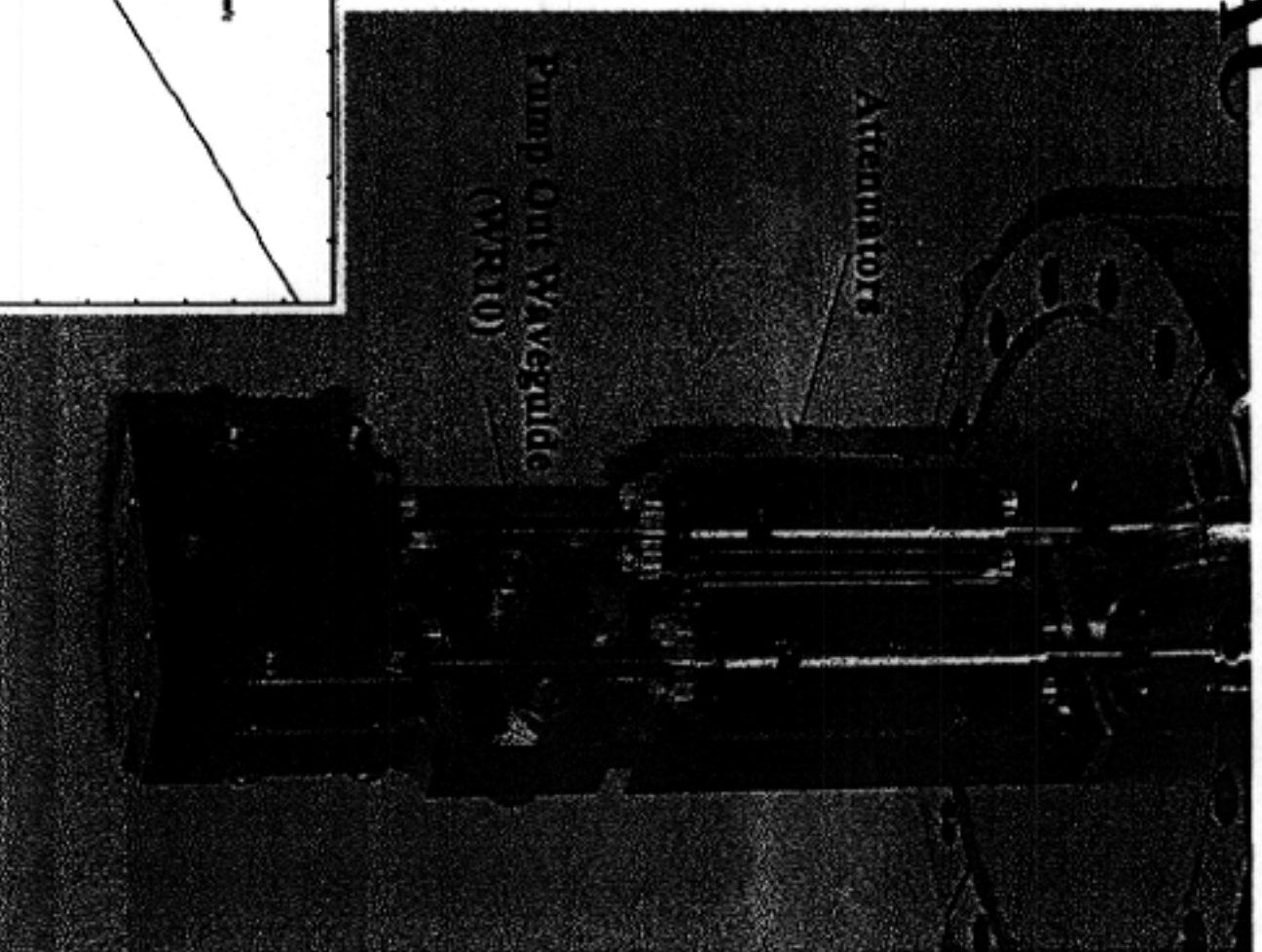
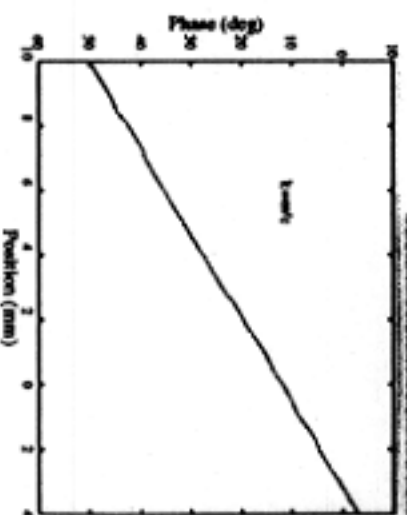
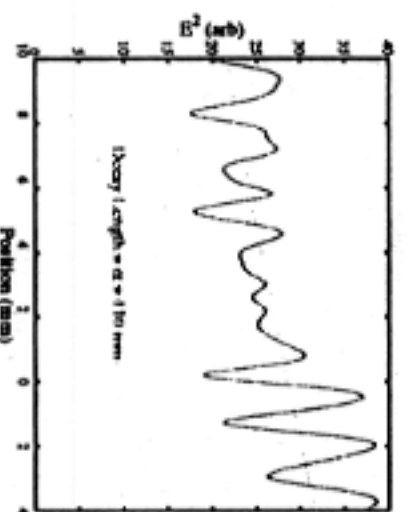
- Powered by beam 300MeV, .4A, 100ns beam of NLCTA
- Bunched at 11.4GHz (8'th subharmonic).
- Beam was focused to ~0.3mm with full transmission
- Resonant interaction yields power of 8kW output from structure. Fields in excess of SLAC linac. M. Hill, et al., in preparation.

W-Band at SLAC

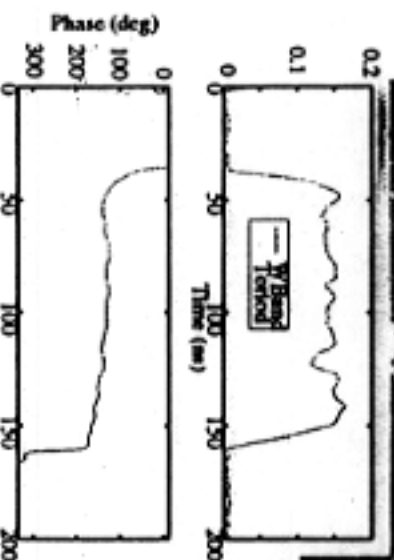
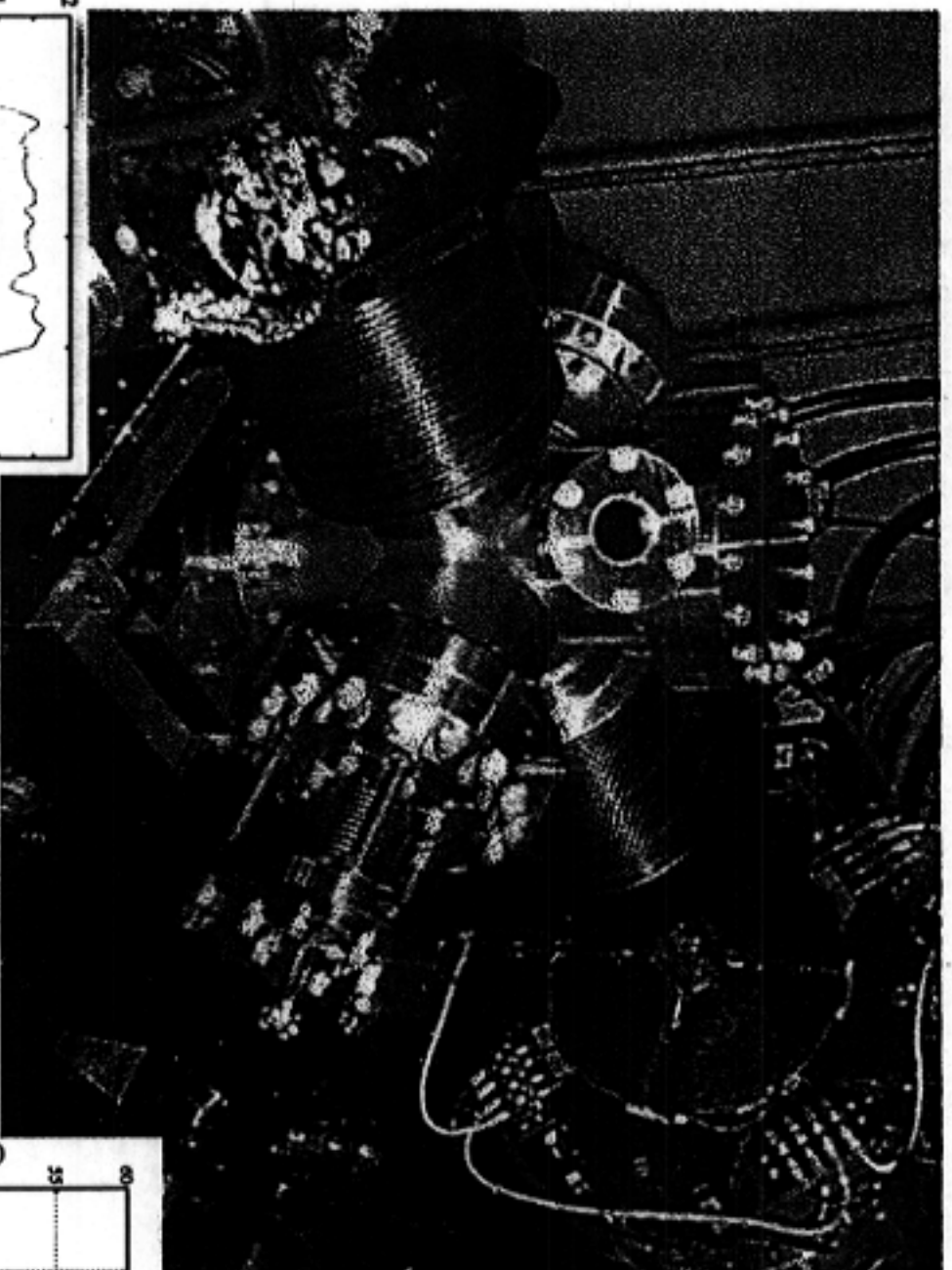
We have fabricated and tested
a W-Band Alumina Dielectric
Structure.

Tests include bench
measurements under low
power and beam tests
producing high gradients.

Beadpull measurements verify proper
phase velocity and Q of the system.

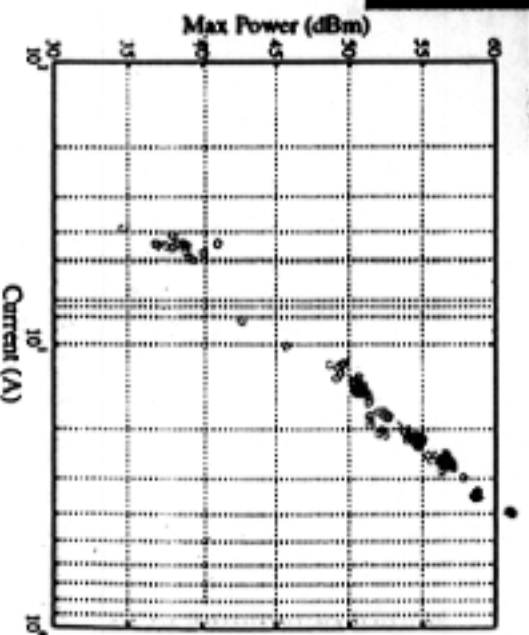


W-Band at SLAC



Tests performed in January '99 produced 1 kW of W-Band power using klystron output cavity tuned to the 6th harmonic of the X-Band beam.

HARPP
HARmonic Power
Production experiments
testing 91.4 GHz
accelerator structures
with a 0.5 A, 11.4 GHz
bunched beam.



250kW

PLASMA ACCELERATORS

Accelerated Bunch

Plasma Oscillation

Beatwave

Drive Pulse

Accelerated Bunch

Laser Wake

Drive Bunch

Accelerated Bunch

Laser Wake





Recent Experimental Advances

- High Gradients demonstrated with unstable laser pulse breaking up in plasma (no injection of electrons), $>100\text{GeV}/\text{m}$. (e.g., Umstadter, et al, Science 273,472 (1996); Gordon et al, PRL 80, 2133 (1998))
 - ➔ Disadvantages for accelerators: Based on instability (phase problems), requires low ω_p/ω —rapid dephasing of particles and wake
 - ➔ Preferable to have “Laser Wakefield Accelerator”—pulse length $\sim c/\omega_p$
- First acceleration of injected electrons in Laser Wakefield Accelerator: $1.5\text{GeV}/\text{m}$, 1.5MeV acceleration. (Dorchies et al, Phys. Plas. 6, 2903 (1999)).
- Guiding of high laser powers ($>10^{17}\text{W}/\text{cm}^2$) in channel for ~ 10 diffraction lengths, low intensity for 100 diffraction lengths ($< 8\text{MW}$; 1m)
- Acceleration in channel is critical
- Recent review: Esarey, et al, IEEE Trans. Pl. Sci. 24,252 (96); Proc. of the AAAAC98 workshop.

SUMMARY OF EXPERIMENTS

	γ_{drive} (MeV)	Q_{drive} (nC)	L_{drive} (mm)	n_p (cm^{-3})	$\Delta\gamma$ (MeV)	E (MV/m)	
ANL (USA)	21	4.0	2.1	10^{12}	0.2	5.0	→ 60 MV/m [99]
KEK (Japan)	500	10	3.0	10^{12}	30	30	
KhFTI (Ukraine)	2	0.4	17	10^{11}	0.5	0.25	
SLAC (USA)	3×10^4	3.2	0.6	10^{16}	-	-	E157 IN PROGRESS

Table 1: Electron Beam Driven Plasma Accelerator Experiments (PWFA)

	I (W/cm ²)	τ_L (ps)	n_p (cm ³)	$\Delta\gamma$ (MeV)	E (GV/m)		
PBWA:							
BEAT WAVE {	ILE (Japan)	10 ¹³	1000	10 ¹⁷	10	1.5	→ 2.8 GV/m Injected Beam
	UCLA (USA)	10 ¹⁴	300	10 ¹⁶	28	2.8	
	LULI (France)	10 ¹⁷	90	10 ¹⁷	1.4	0.6	
SM-LWFA:							
SELF-MODULATED {	KEK (Japan)	10 ¹⁷	1.0	10 ¹⁹	17	30	
	LLNL (USA)	10 ¹⁸	0.6	10 ¹⁹	2	-	
	RAL (UK)	10 ¹⁹	0.8	10 ¹⁹	44	100	
	NRL (USA)	5 × 10 ¹⁸	0.4	1.4 × 10 ¹⁹	> 1	-	
	CUOS (USA)	4 × 10 ¹⁸	0.4	3.6 × 10 ¹⁹	> 1	-	
LWFA:							
WAKEFIELD {	KEK (Japan)	10 ¹⁷	1.0	10 ¹⁵	5	0.7	→ 1.5 GV/m Injected Beam
	LULI (France)	4 × 10 ¹⁷	0.4	2 × 10 ¹⁶	1.6	1.5	

Table 2: Laser-Driven Plasma-Based Accelerator Experiments

Table 2: Laser-Driven Plasma-Based Accelerator Experiments

PLASMA

Accelerators

WAKE



$$\lambda_p \equiv \frac{2\pi c}{\omega_p}$$

$$E_z \sim \left(\frac{\xi n}{n}\right) \frac{mc\omega_p}{e} \sim 100 \sqrt{n \text{ cm}^{-3}} \text{ V/m}$$

$$E_\perp \sim k_{\perp} r E_z$$

IN GENERAL, $\frac{\partial E_z}{\partial r} \neq 0$, $E_\perp$ has non linearity



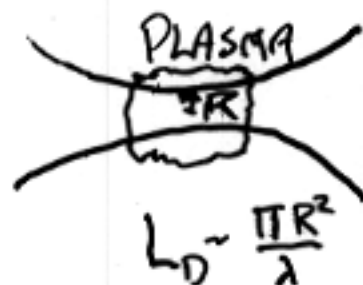
The Creation of Plasma Structures



Mechanisms for Channel Creation

- Heating center of plasma
 - single pulse for ionization and heating (Milchberg et al 93-98)
 - dual pulse scheme (Leemans and Volfbeyn, 98)
- Capillary discharge (Ziegler et al 96-98)
- Laser Nonlinearity
 - relativistic guiding (Max et al 74)
 - Self-channeling (plasma blowout) (Umstadter)
 - Electromagnetic channels (Shvets et al, 98)

LASER WAKEFIELD REQUIRES PLASMA CHANNEL



Graph showing the refractive index $\tilde{n}$ versus radius r . The curve starts at $\tilde{n} = 1$ on the y-axis and decreases as r increases, intersecting a horizontal line at n_p .

$$\tilde{n}(r) \approx 1 - \frac{n_p(r)}{2n_c}$$

- Channels overcome diffraction of the drive laser pulse.

- Guiding in a pre-formed channel gives the option to operate in a linear regime.
- Transverse field profiles determined by the driver profile: $\nabla_{\perp} E \sim \nabla_{\perp} a$
- Field structure similar to the homogeneous plasma case.

Parabolic channels originally studied by Sprangle and co-workers [IEEE Trans. Plasma Sci. PS-15, 145, 1987]. Pulse propagation instabilities in channels Shvets and Wurtele [Phys. Rev. Lett. 73, 3540, 1994] and by Antonsen and Mora [Phys. Rev. Lett. 74, 4440, 1995].

- Hollow channels have additional desirable properties.

- Electromagnetic mode.
- Transversely uniform accelerating gradient: $E_z \sim \text{const.}$
- Linear focusing fields: $dE_r/dr \sim \text{const.}$

Originally examined by Shvets *et al.* [IEEE Trans. Plasma Sci. 24, 351, 1996].



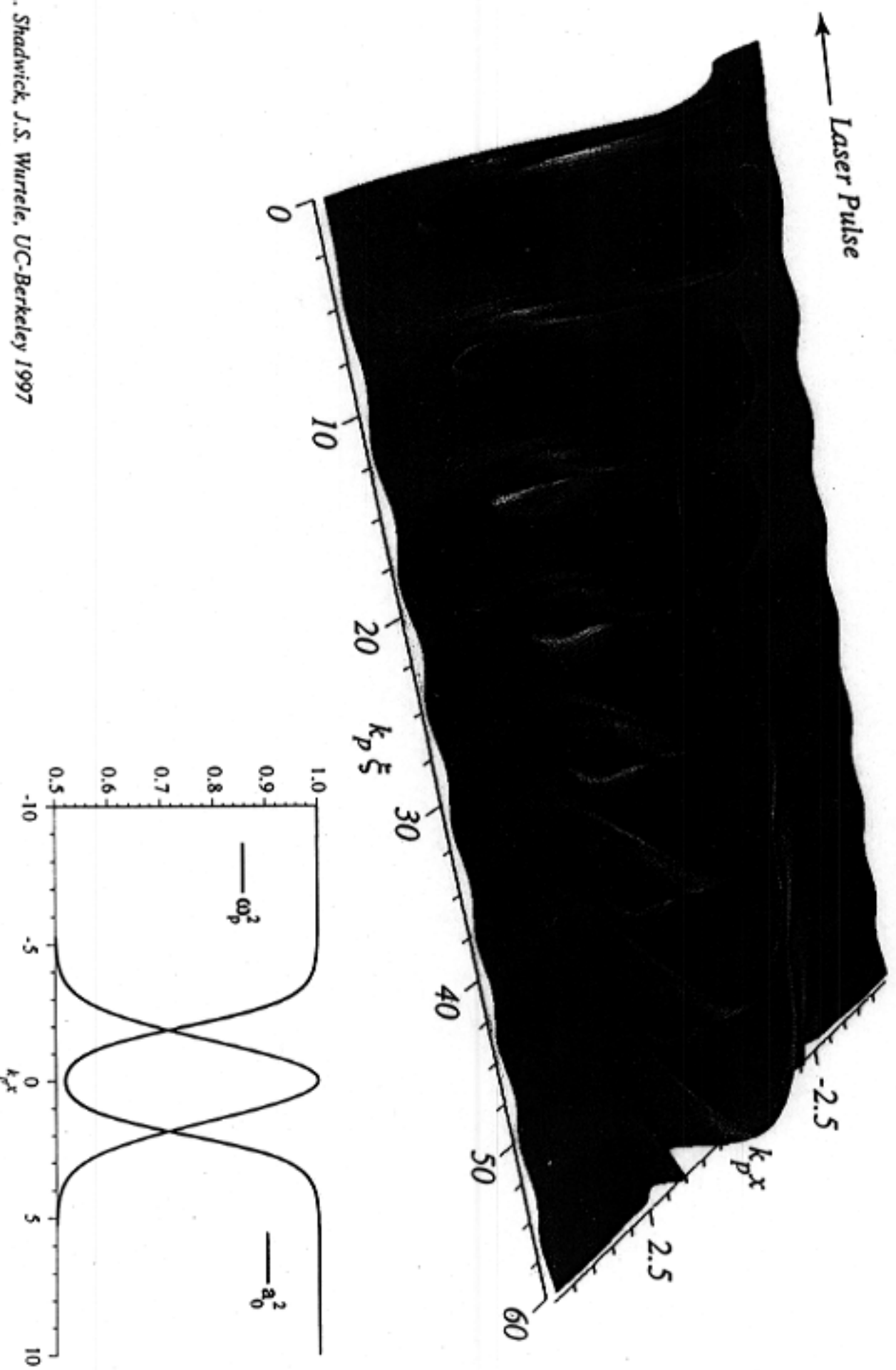
The Quality Factor of a

Plasma Structures

Dissipation in a collisionless regime

Longitudinal Wake Field

Numerical Solution of the Linearized Cold Fluid Equations



Hollow Channels

- Ideal Case

- Analytic solution of linearized equations possible.
- Mode is electromagnetic with $\omega_{ch} < \omega_{p0}$.
- Mode is undamped ($Q = \infty$).



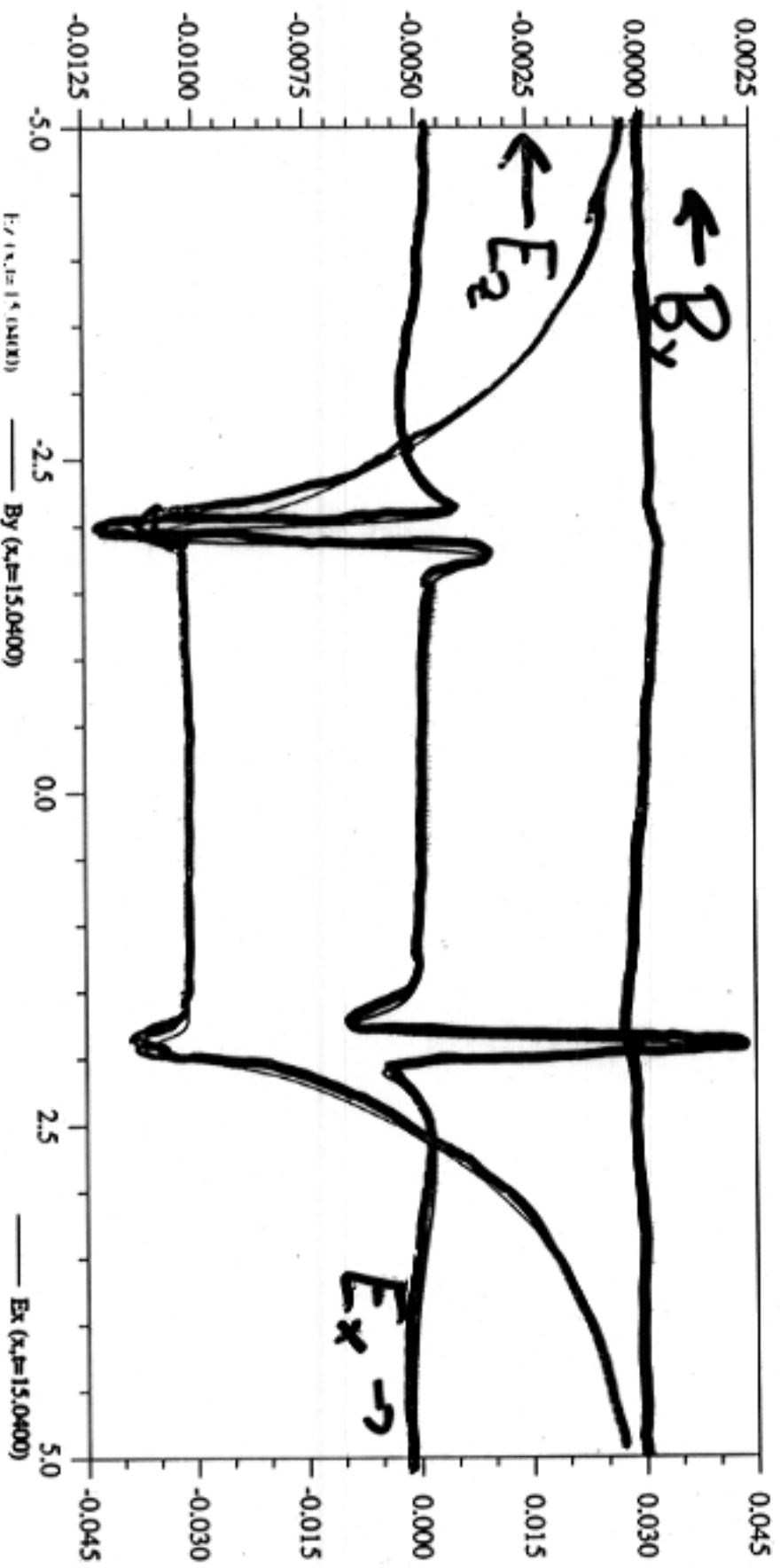
- Non-ideal Case

- Linear equations analytically challenging.
- Dielectric function is spatially dependent: $\epsilon = 1 - \omega_p^2(x)/\omega^2$.
- Every $\omega < \omega_{p0}$ is resonant with the local plasma frequency at some location in the wall.
- Resonant layer leads to absorption of the wake.
- Low Q means a large spread around ω_{ch} resonantly exciting much of the wall.



$$B_y'' - \underbrace{\frac{\epsilon'}{\epsilon}}_{\text{resonance}} B_y' - \frac{\omega_p^2(x)}{c^2} B_y = 0.$$

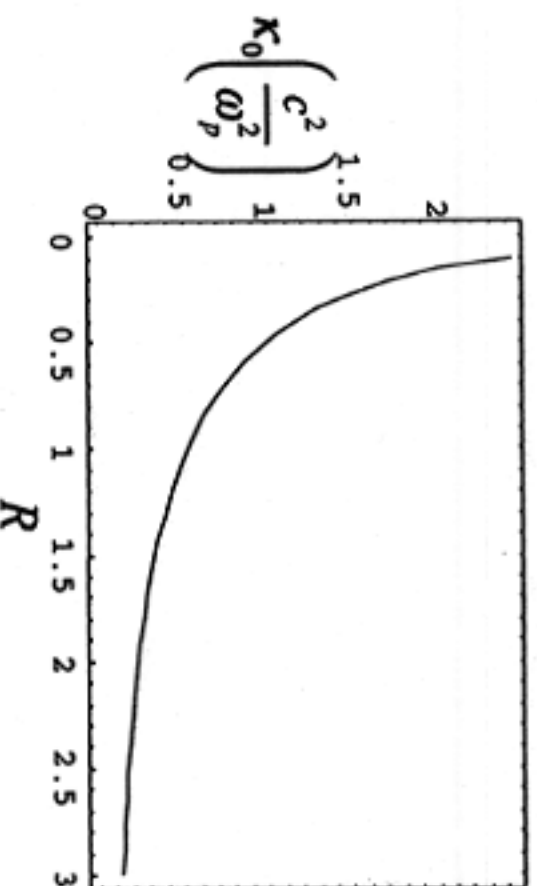
B_y : WELL BEHAVED
 E_x : LARGE @ $x = x_c$



Loss Factor:

Definition: $\kappa_m = E_z^2 / 4U_m$
independent of excitation mechanism, geometrical factor of structure

$$\kappa_m = \frac{\omega_p^2}{c^2} \left[\frac{K_m(R)}{RK_{m+1}(R)} \right] \left[1 + \frac{RK_m(R)}{2(m+1)K_{m+1}(R)} \right]^{-1}$$



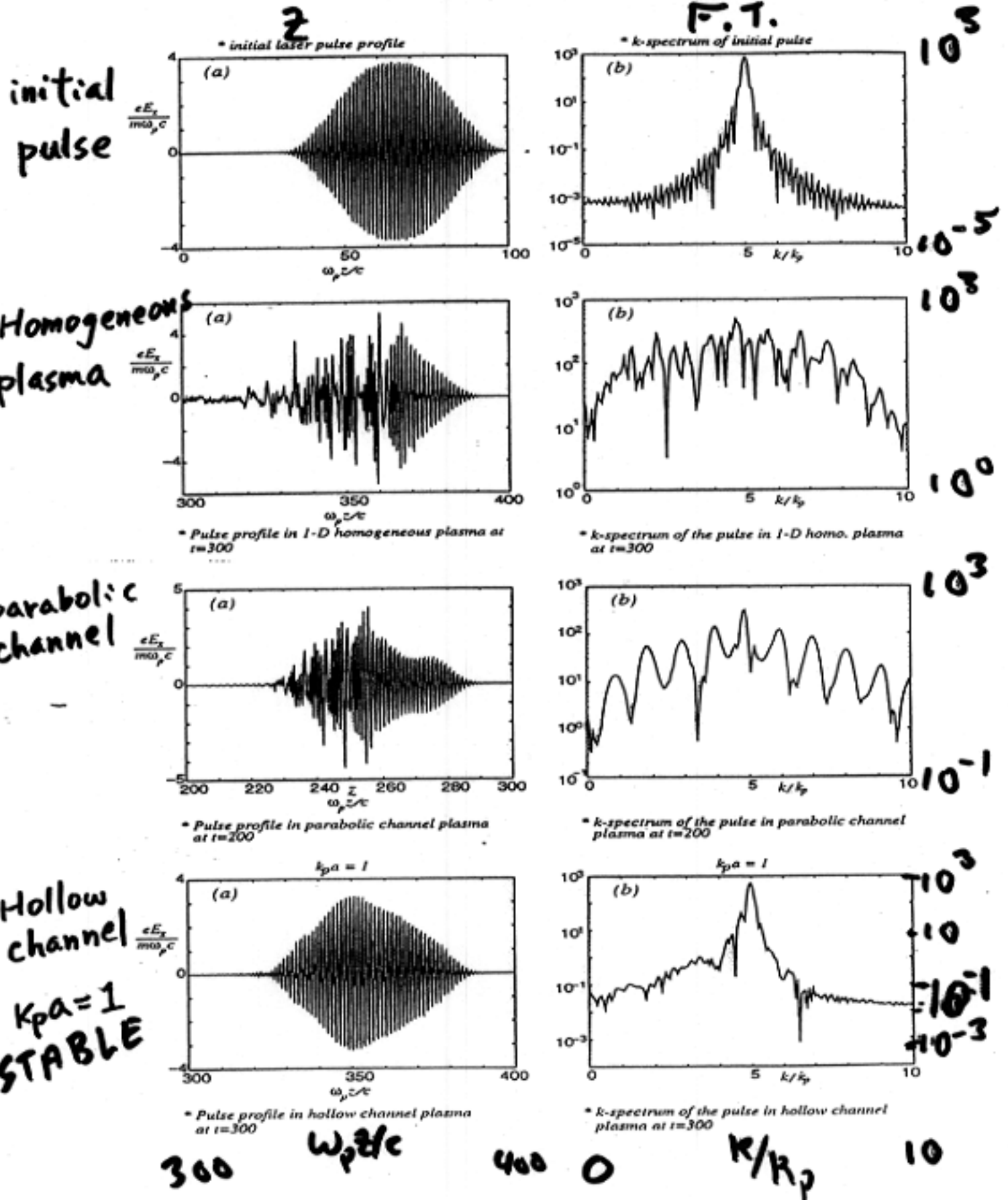
Linearity requires: $\omega_p N_b / c \ll R^2$

Comparison: conventional conducting structure,

$$\kappa_0 = \frac{3.55 \times 10^3}{\lambda^2 [cm]} \frac{V}{pC - m} \left[\frac{K_0(R)}{RK_1(R)} \right]$$

$$\kappa_{SLAC} = \frac{2.1 \times 10^3}{\lambda^2 [cm]} \frac{V}{pC - m}$$

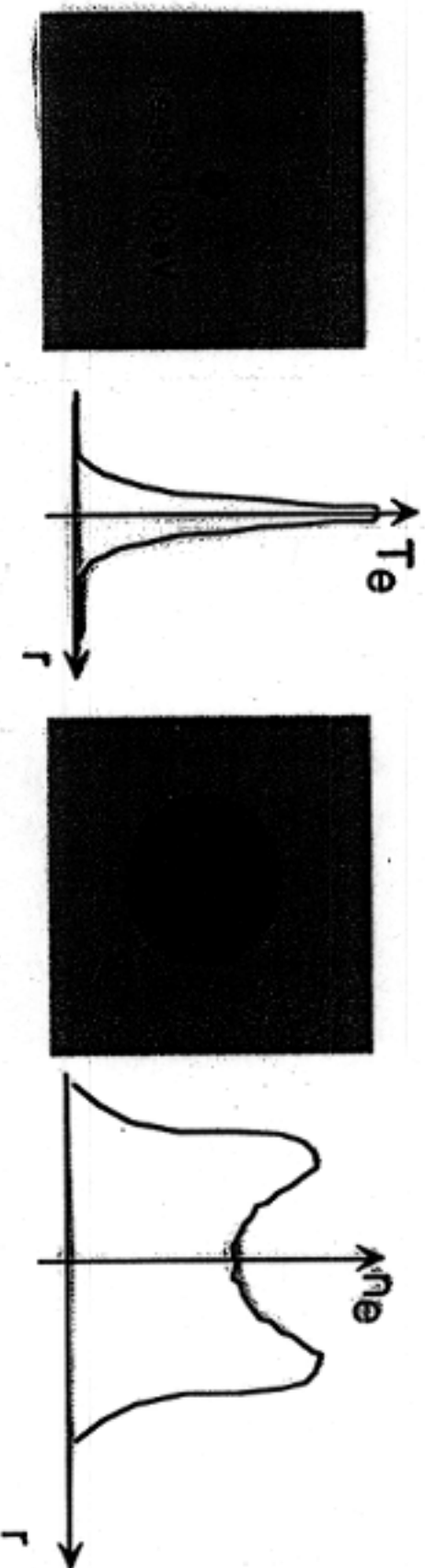
Comparison of Laser Instabilities





Channel Creation*

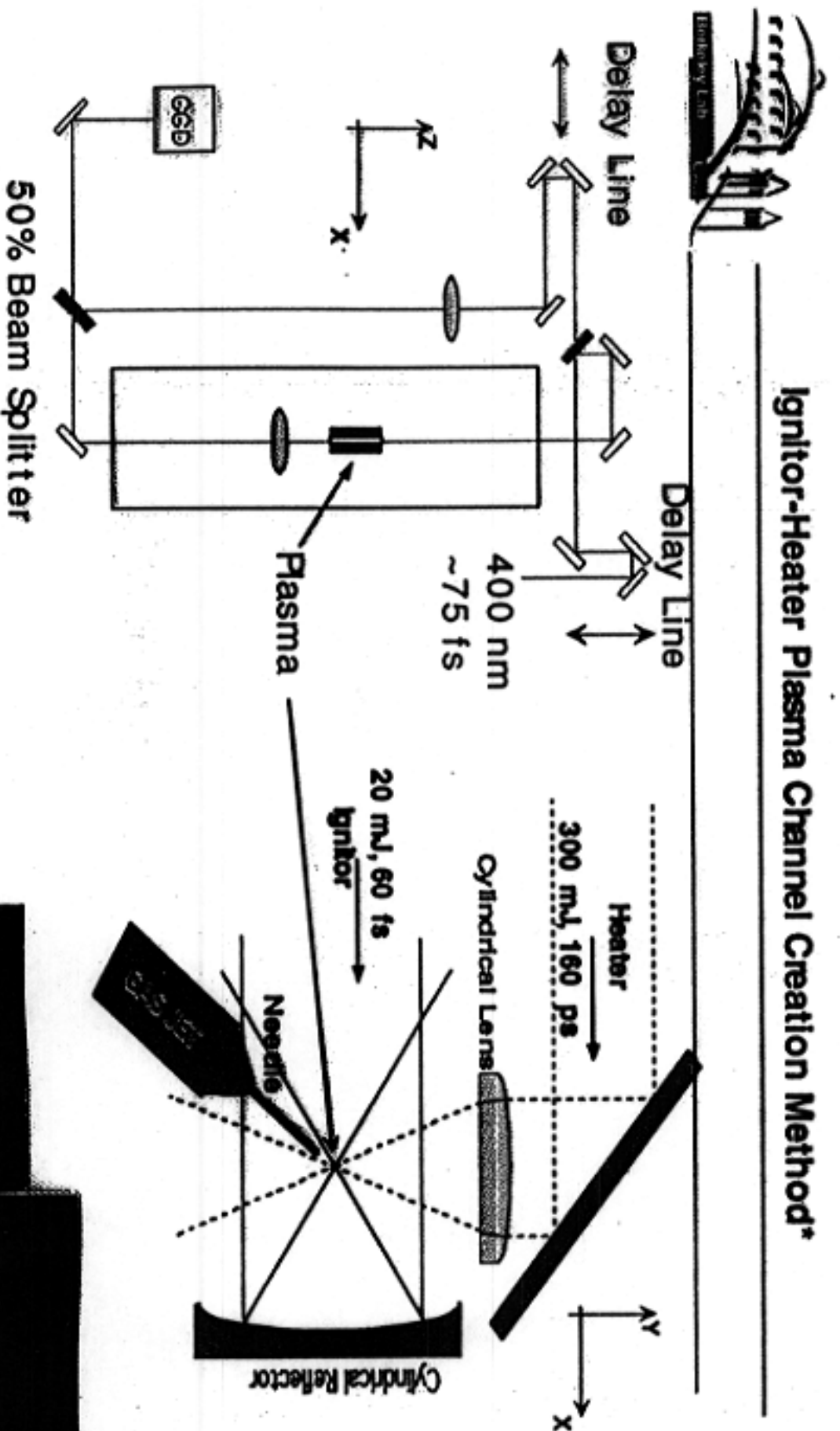
Ionizing and Inverse Bremsstrahlung Heating a thin cylinder of gas leads to hydrodynamic expansion and a Plasma Channel with a density minimum on axis is formed.



Heated Plasma Spark creates a Hydrodynamic shock
that leaves a density minimum on-axis

*G. J. Dufree et al. Phys. Rev. E, v. 51, no. 3, 2368 (1995)
P. Völfbeyn et al. Phys. Plasmas, v. 6, no. 5, 2269 (1999)

Ignitor-Heater Plasma Channel Creation Method*



Transverse Interferometry With Femtosecond Resolution

*P. Volfbeyn et al. Phys. Plasmas, v. 6, no. 5, 2269 (1999)

Accelerator and Fusion Research Division

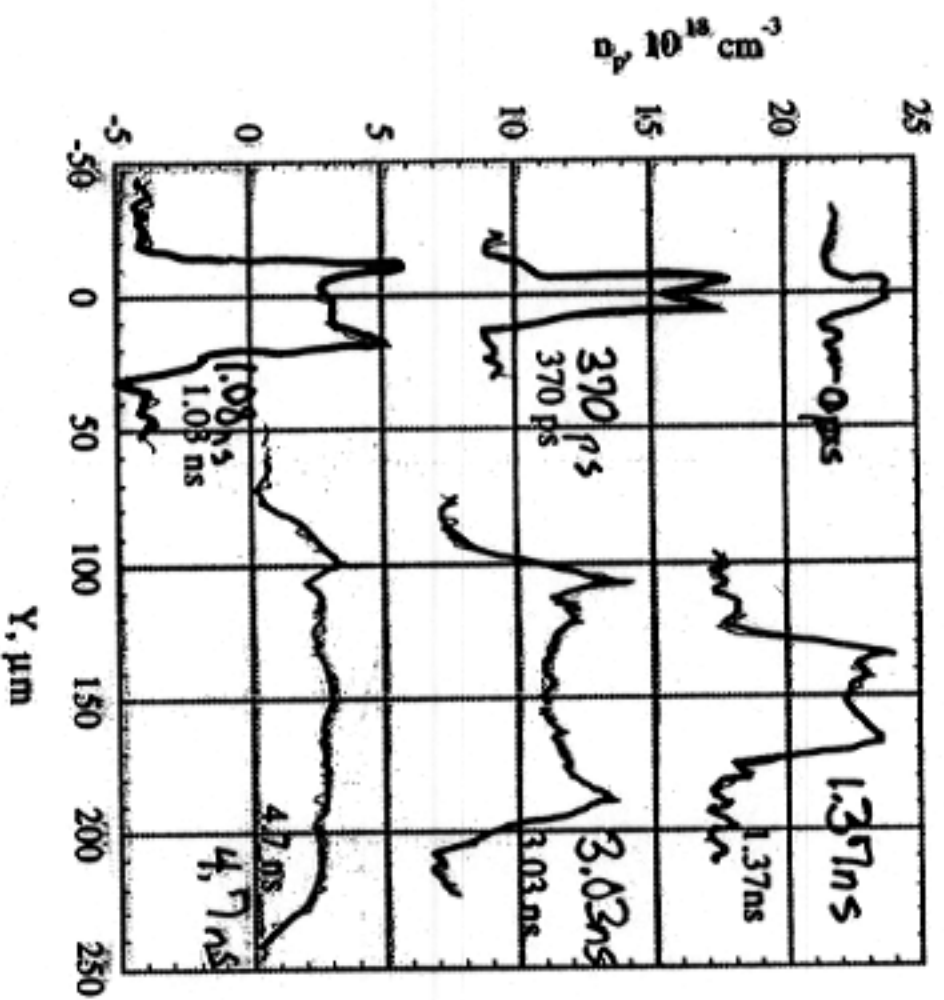
CENTER FOR BEAM PHYSICS

666 pgs

1 fig. no.

2 fig. no.

Channel Formation in N₂ Gas Jet



$E_{\text{shot}} = 50 \text{ mJ}$, $E_{\text{avg}} = 280 \text{ mJ}$
 1000 psi, 0.450 mm from nozzle

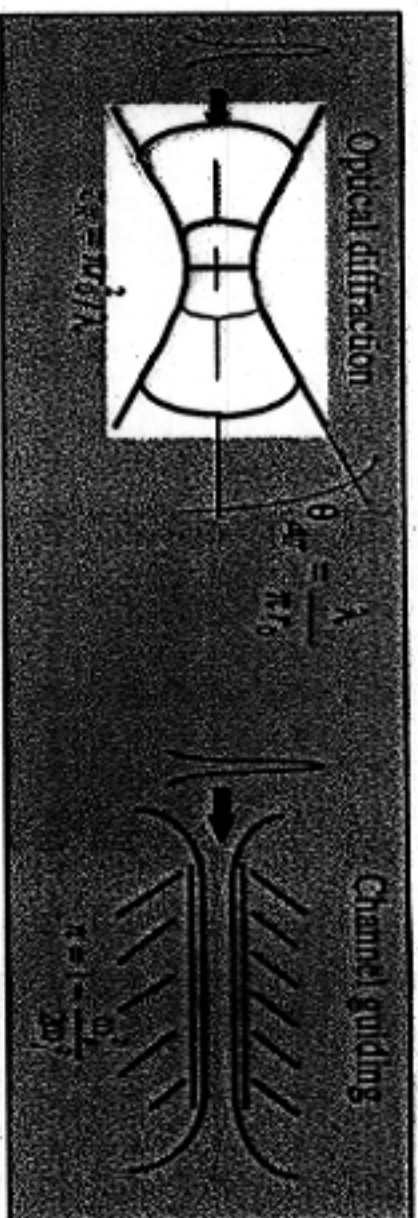
LBL PLASMA WAKEFIELD - FORMATION OF CHANNEL

1000 psi, 0.450 mm from nozzle

+



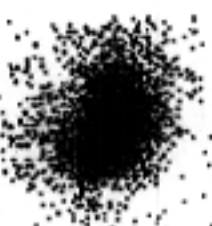
2×10^{17} W/cm² Guided over 5 Rayleigh lengths (1 mm)



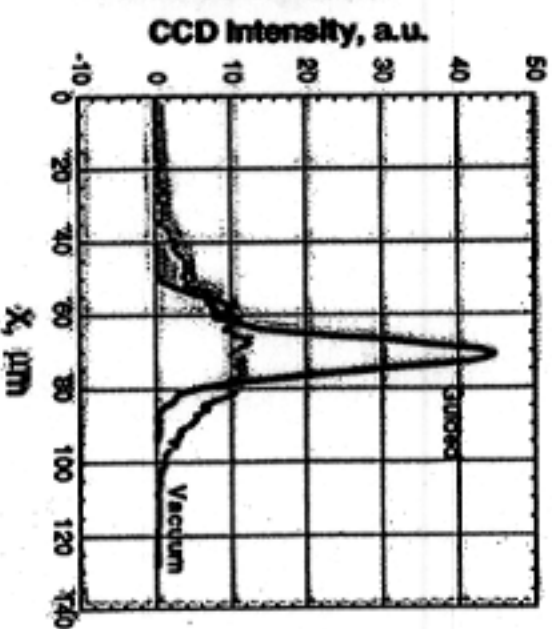
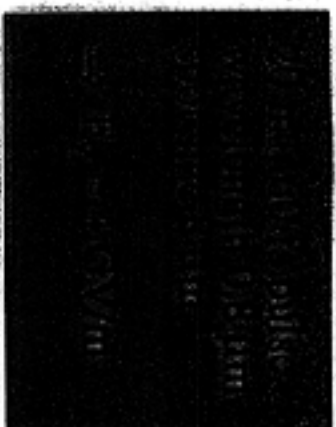
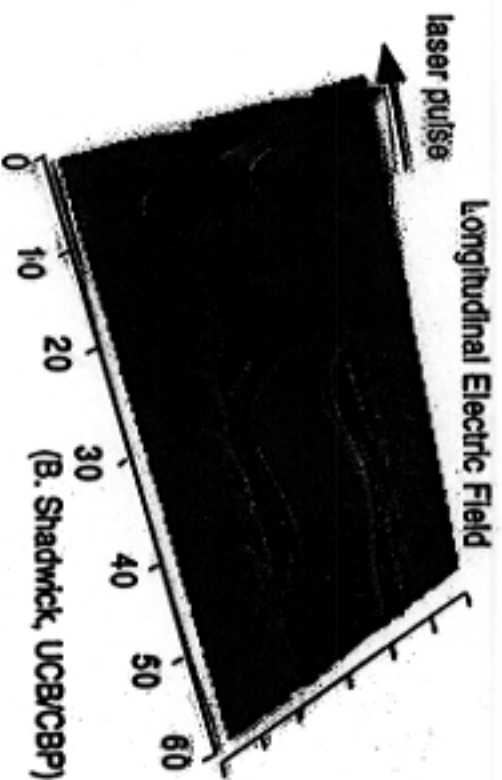
Laser Beam Images:

In Vacuum

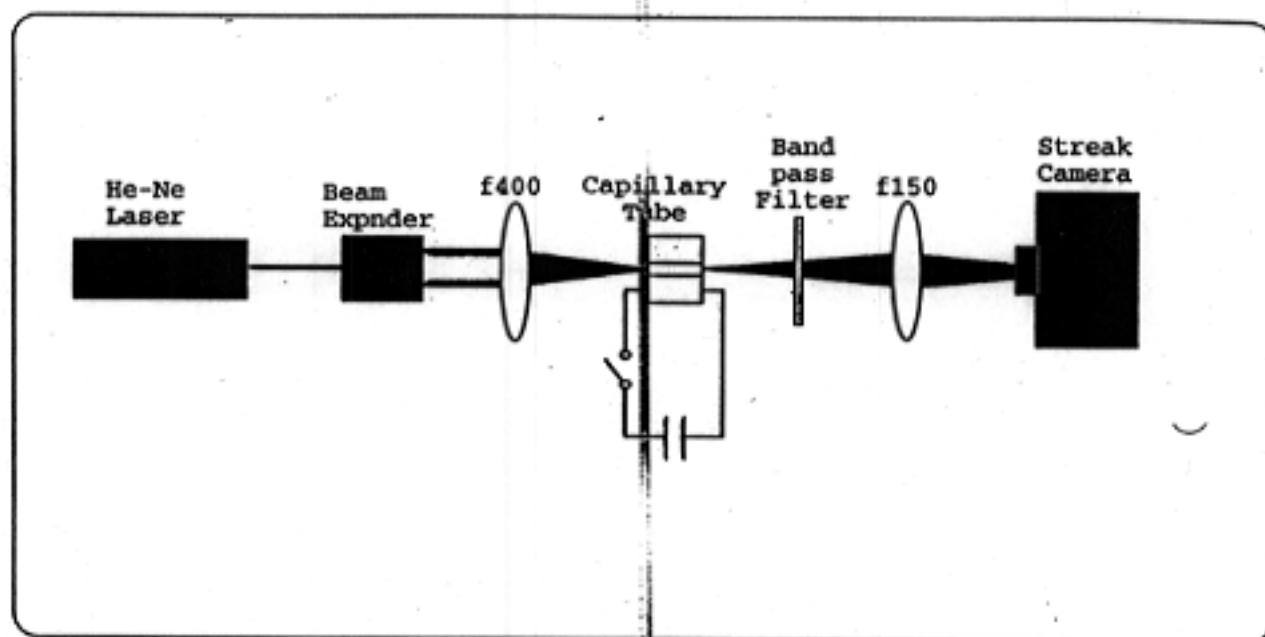
Channel On!



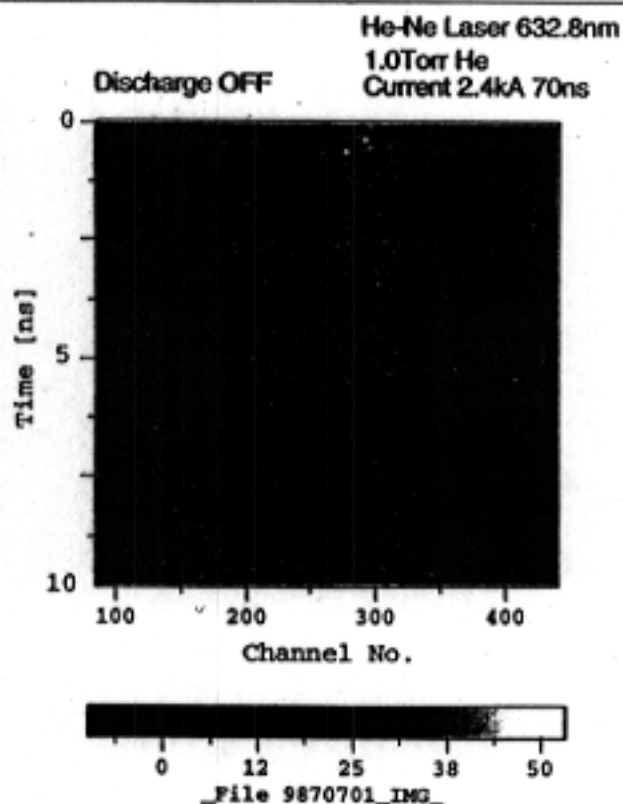
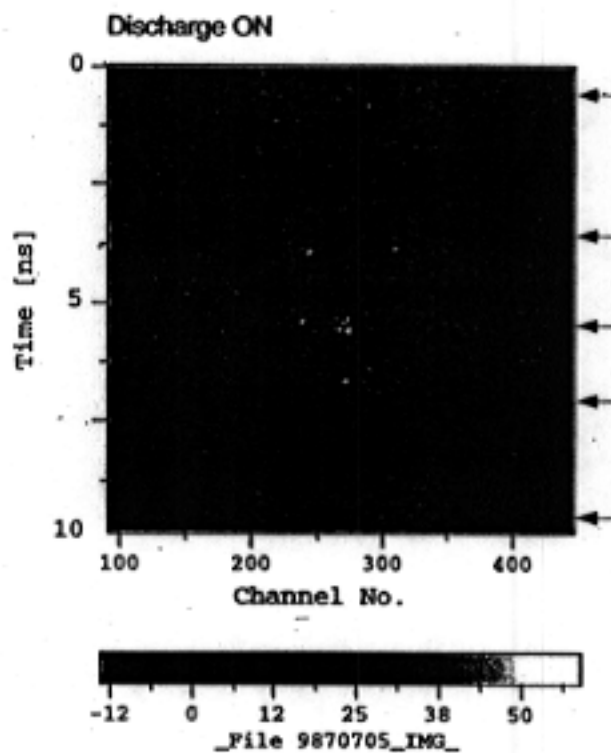
Vacuum Focus Intensity $\sim 2 \times 10^{17}$ W/cm²



Experimental Set up for Plasma wave guide



Laser Beam Profile at Exit of Capillary





Laser Self-Guiding and Preformed-Channel Guiding

DANIEL MÜSTADTER
UMICH

Side Images of Pump Pulse in a Gas Jet

Laser beam images
at the exit of gas jet

1.0 TW



laser power below guiding threshold

100 μm
FWHM

laser power above guiding threshold

2.5 TW



1000 μm

10 μm
FWHM

laser ponderomotive force
creates plasma density
depression

probe pulse



Plasma Density Profile
Obtained Using Probing Interferometry

T+40 ps



6 x 10¹⁹ cm⁻³

Probe beam images

no pump pulse



100 μm
FWHM

with pump pulse



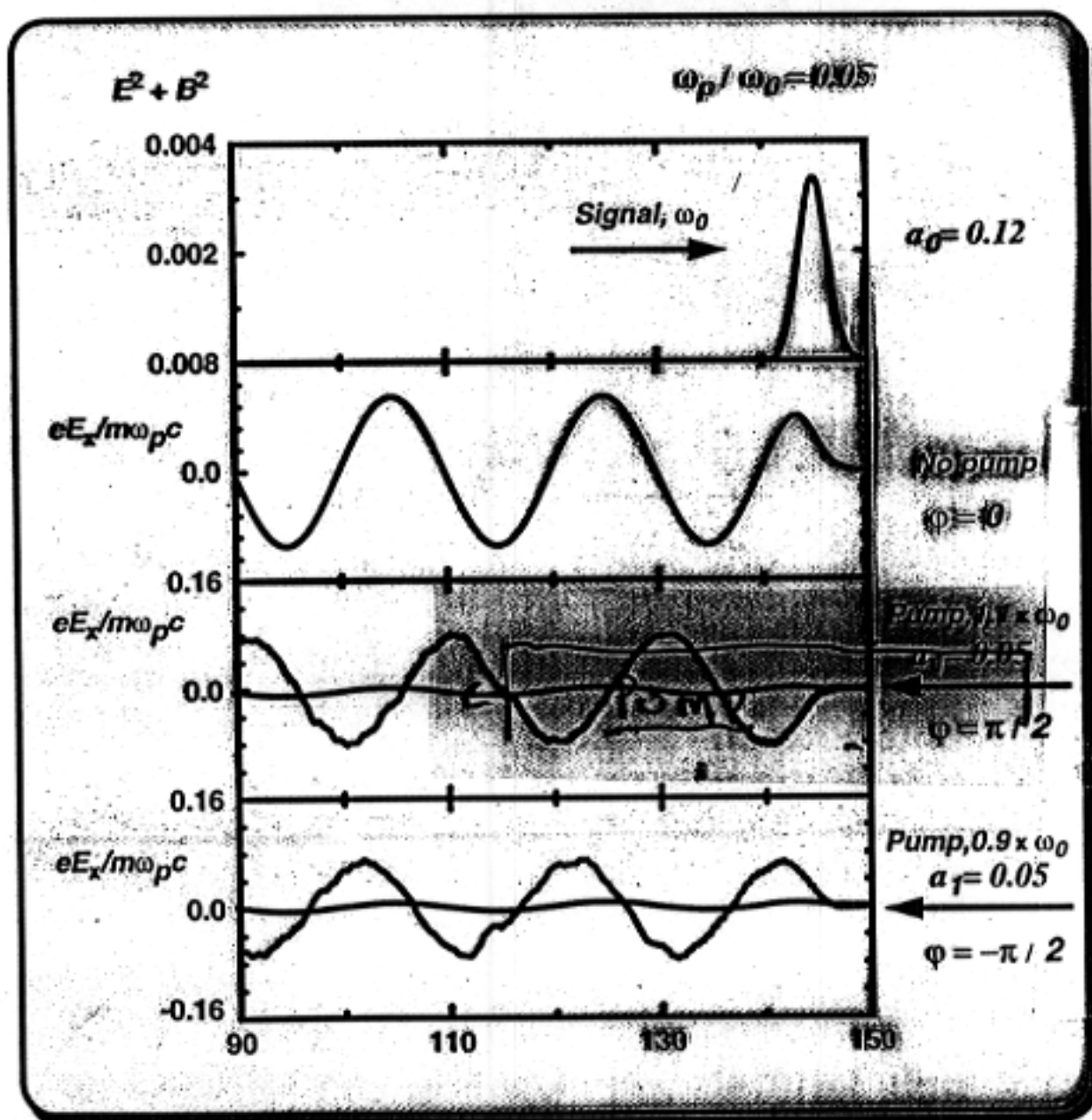
8 μm
FWHM

Enhanced Wake Field Generation by Colliding Laser Beams



Standard wake: $eE_x/m\omega_p c \sim a_0^2$

Enhanced wake: $eE_x/m\omega_p c \sim (a_0^2)^{1/2}$



Shvets, Fisch, PPPL

Pukhov, Meyer-Ter-Vehn, Max Planck



Plasma Wakefield Accelerator

- Linear excitation

- ➔ Nonlinearities in accelerating and focusing fields
- ➔ Higher accelerating gradient

- Nonlinear excitation

- ➔ most of beam focused linearly by ions
- ➔ Lower accelerating gradient

- Physics issues:

- ➔ Beam dynamics--drive bunch and accelerated bunch
- ➔ Efficiency

- Potential advantages in ultra relativistic regime

- ➔ Stiff beam driver--longer interaction lengths,
- ➔ higher coupling efficiency (?).

NO NEED FOR Preformed Plasma Structure



Ultra-Relativistic Beam-Plasma Dynamics



Develop diagnostics, benchmark simulations,

improve design tools, understand plasma as a *dynamic* element of an accelerator.

- Experiments: UCLA, USC, LBNL, BINP, SLAC E150 & E157

Is a multi-GeV beam driver superior to a laser?

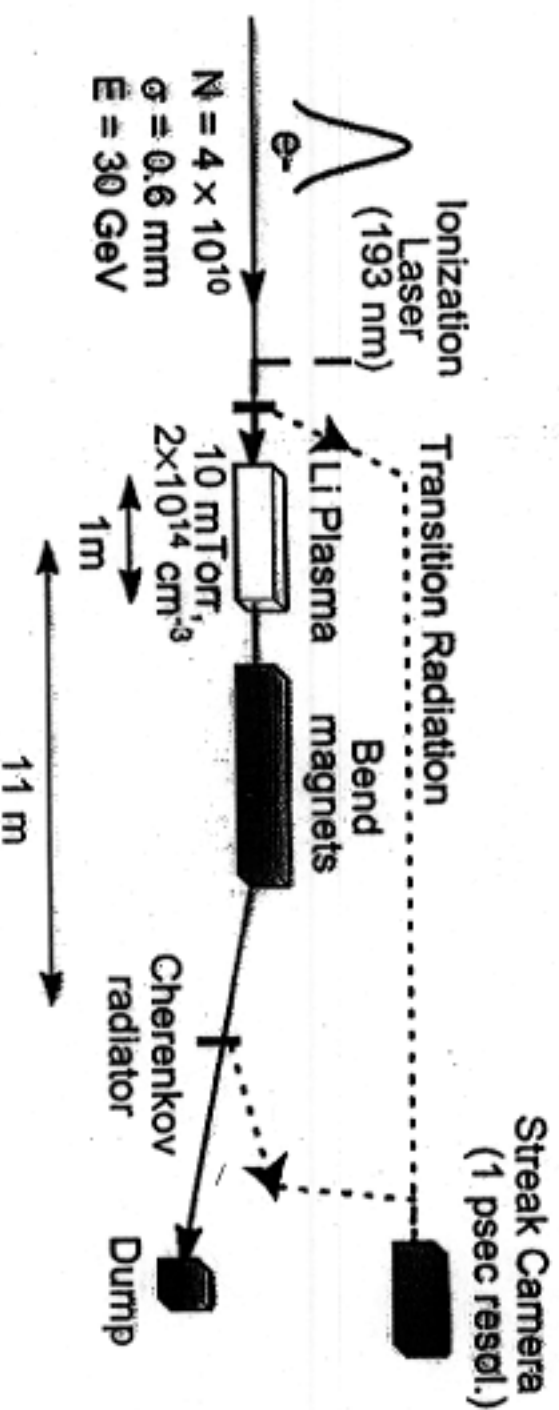
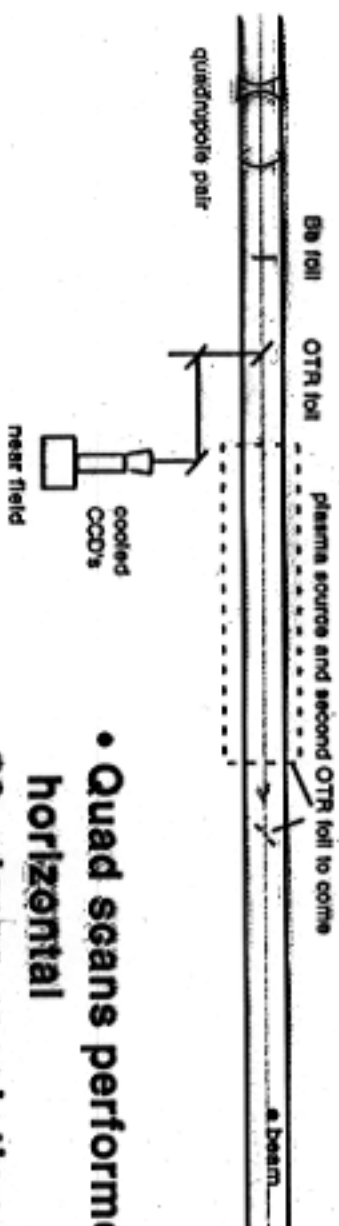


Figure 1: Layout of E-157 (PAC99).

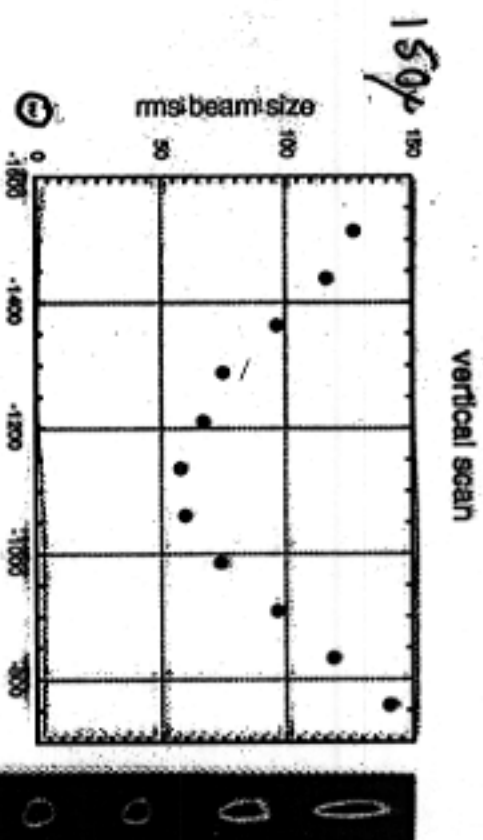
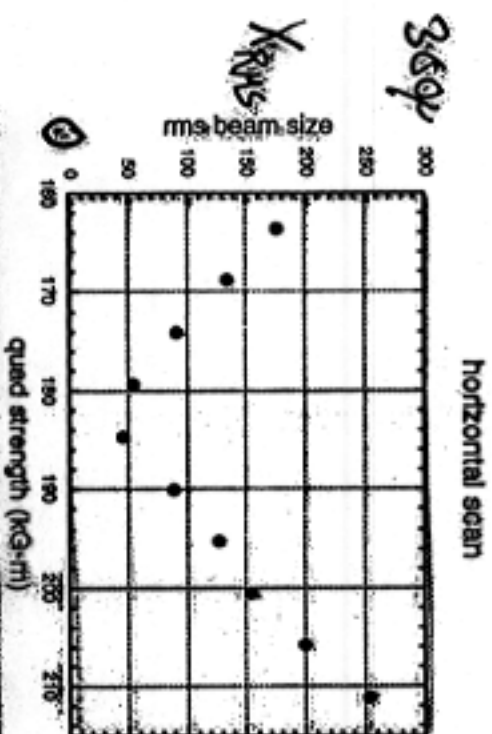
Transition radiation diagnostic of beam size with quad scan



OTR at 30 GeV - Beam spot size resolution



- Quad scans performed in vertical and horizontal
- 50 micron resolution achieved in both axes using visible wavelength OTR.



← OTR near field images →

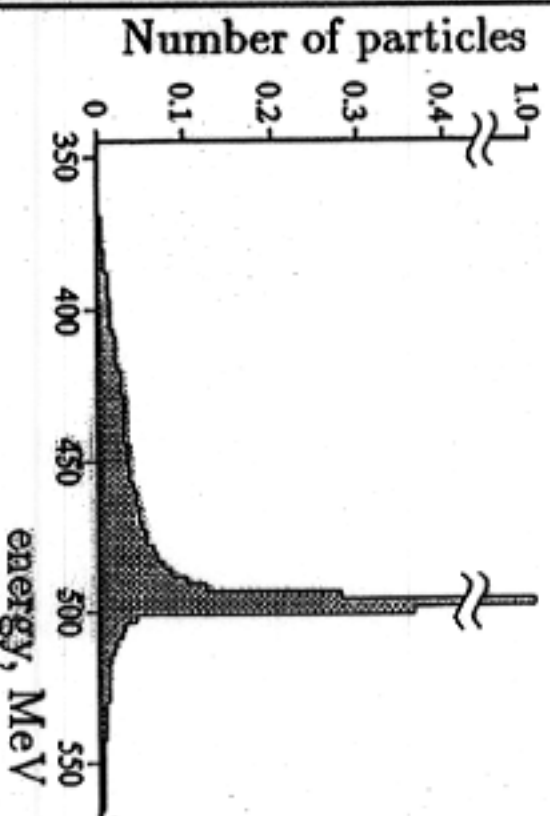


E157: FROM CATRAVNS et al, PAC99



The BINP beam-driven plasma linac

- Density: $10^{15} / \text{cm}^3$
- Gradient: 1 GeV/m , 100 stages @ 10 m each
- Final energy: 1 TeV , energy spread 3% , $N = 5 \times 10^9$,
- Excitation by train of drive pulses
- From Kudryatsev et al. in Proc. ICFA workshop on Advanced Accelerators, NIM 410, 1998. Also work by Chen, Katsouleas, Nakajima, Rosenzweig, others.



Predicted performance of initial experiment at BINP

Drive bunches

High energy bunch



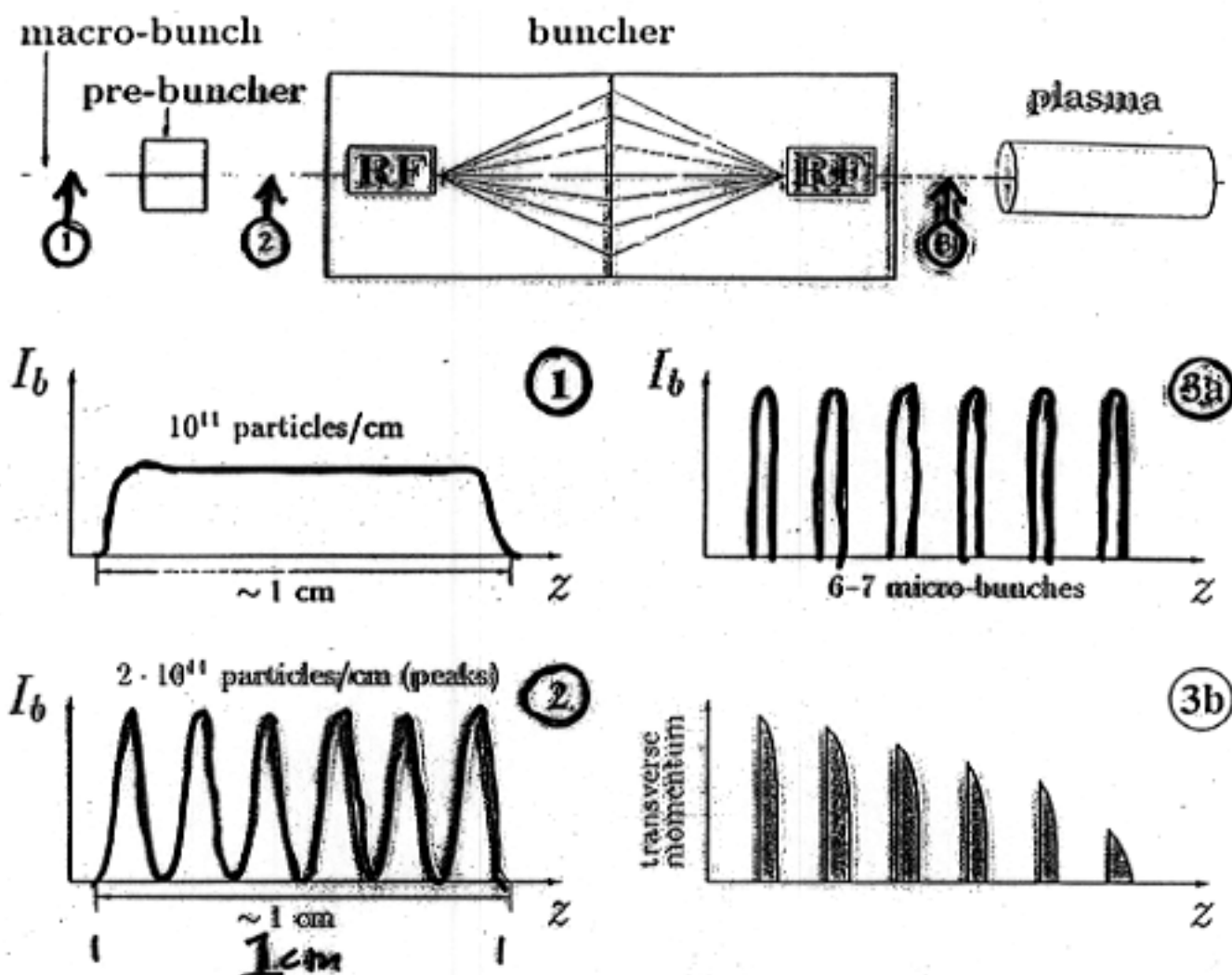


Fig. 5. Driver preparation with the transverse cutting.

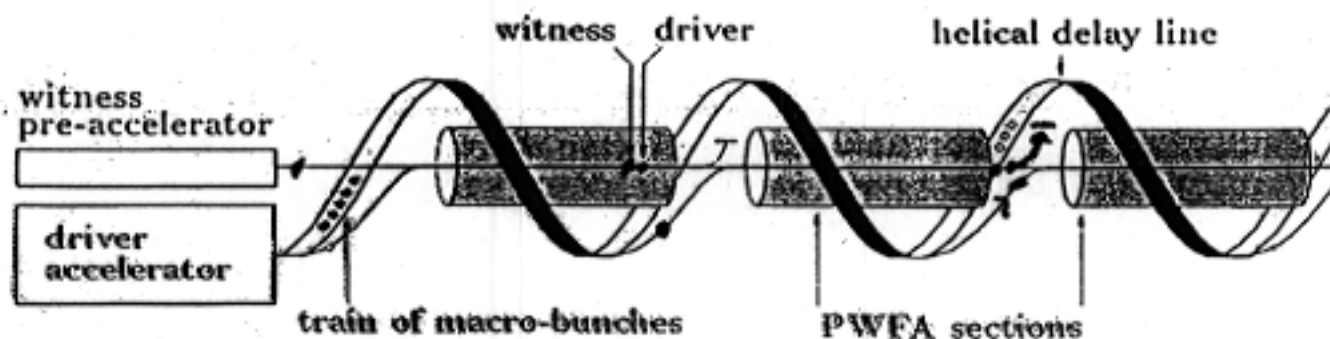
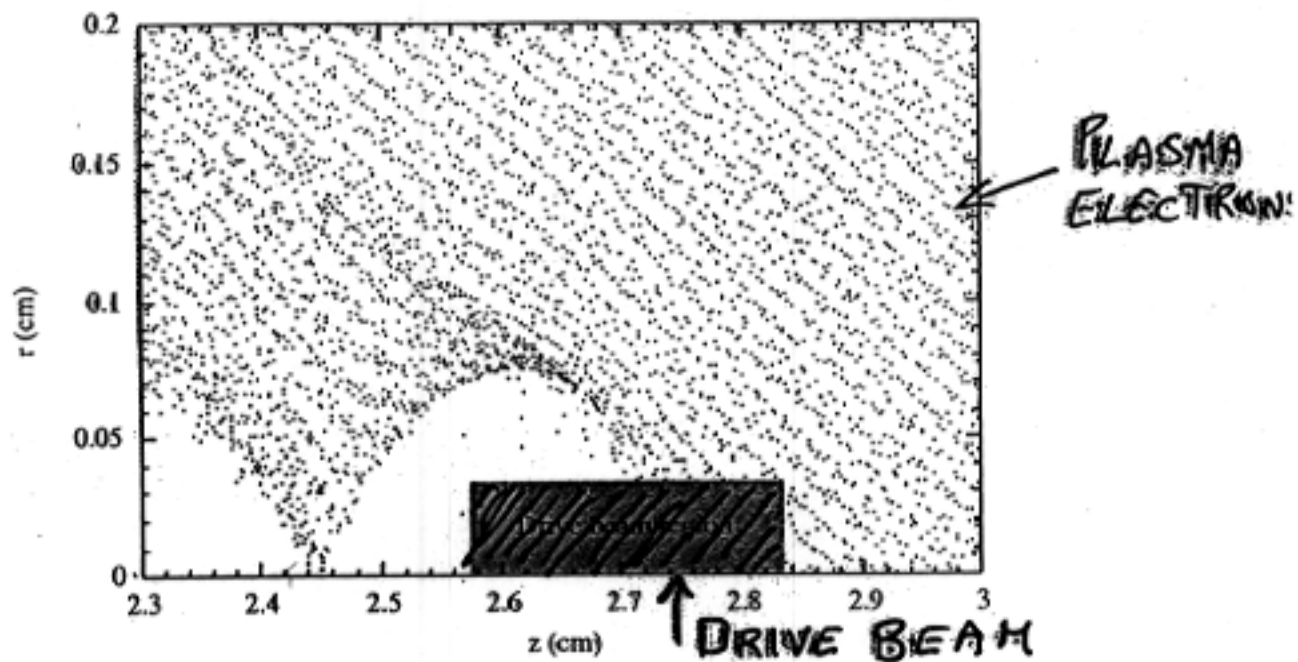


Fig. 6. The staging scheme for the PWFA-based collider.

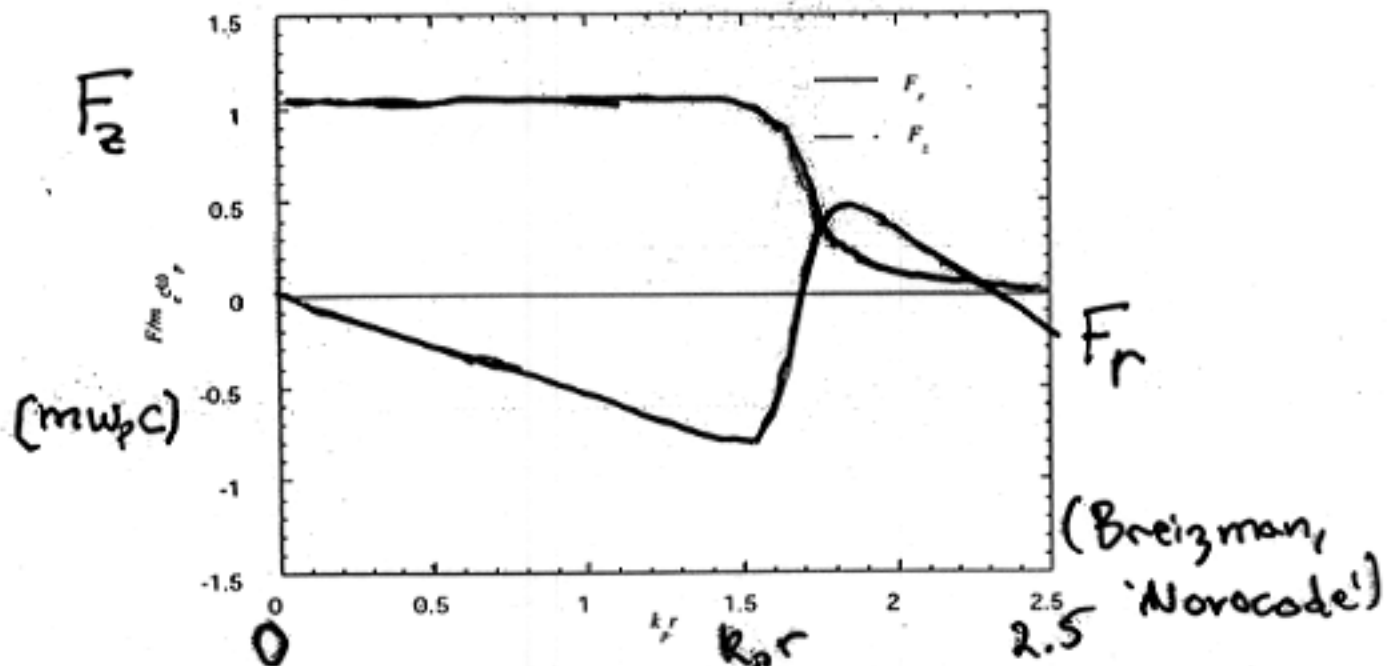
SKRINSKY, et al
BINP

NONLINEAR STRUCTURE FORMATION

PHYSICS OF THE PWEA BLOW-OUT REGIME



PIC simulation results of blow-out regime case.

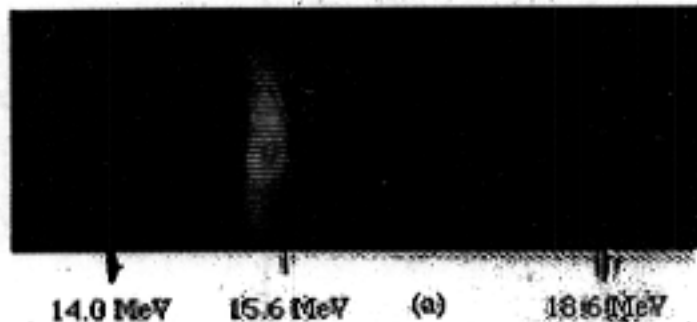


Complete rarefaction of beam channel allows linear electrostatic focusing and electromagnetic acceleration independent of r .

ACCELERATION EXPERIMENTS

- Drive beam deceleration, with acceleration of tail observed in high resolution spectrometer
- Beam parameters: $E=15.6$ MeV, $Q=18$ nC, Pulse length 20-22 psec FWHM, norm. emittance 180 mm-mrad.

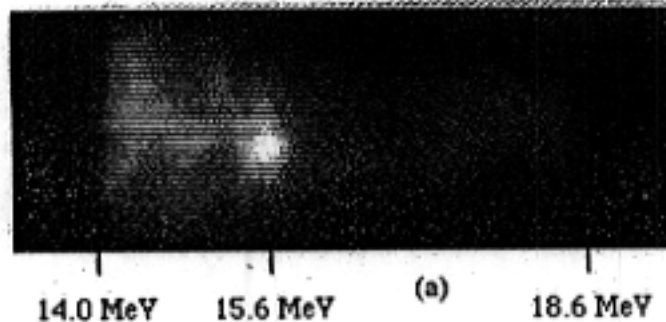
Spectrometer images



14 MeV 15.6 MeV (a) 18.6 MeV
 No plasma case

$$\frac{n_b}{n_0} \sim 3$$

$$n_0 = 1.3 \times 10^{13} / \text{cm}^3$$



14 MeV 15.6 MeV (a) 18.6 MeV
 Plasma case

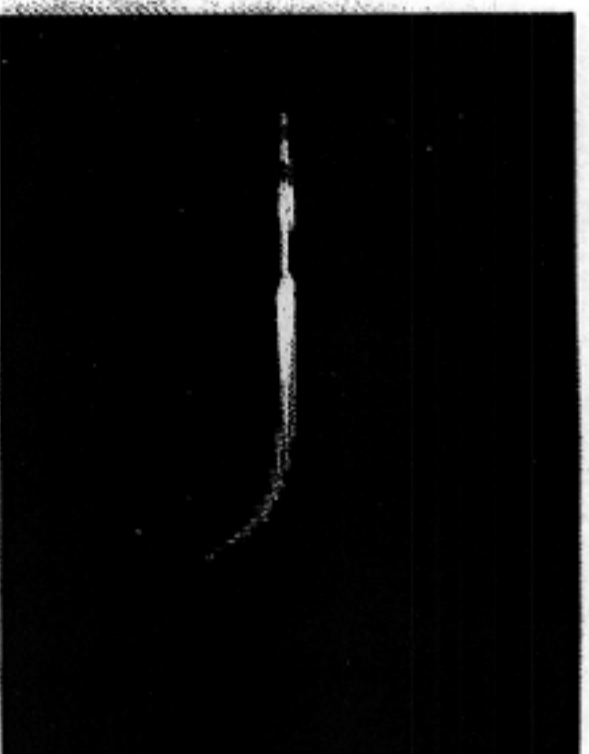
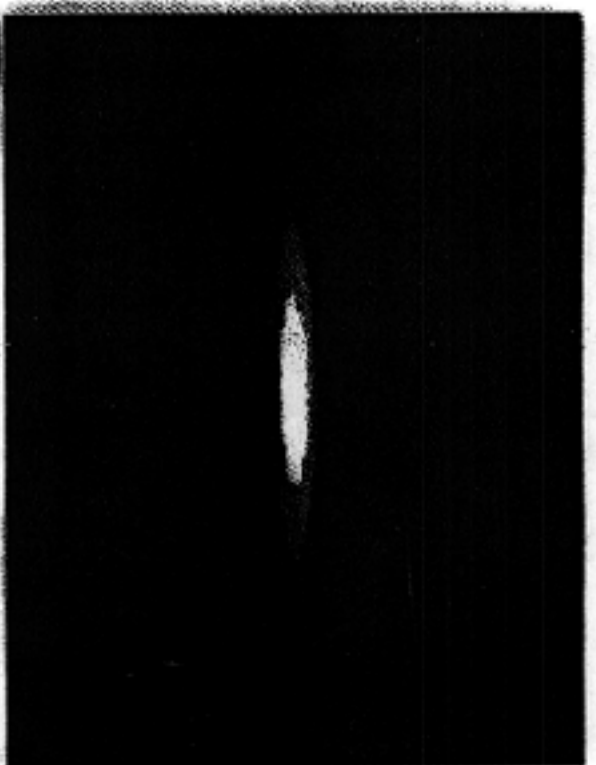
$$R_p G \sim 2$$

(deceleration of 13 MeV/m, acceleration of 25 MeV/m)

FNAL 'PLAN': • Photocathode

- COMPRESS DRIVE BEAM
- GV/m Gradients
- Efficiency studies

Wakefield Amplification

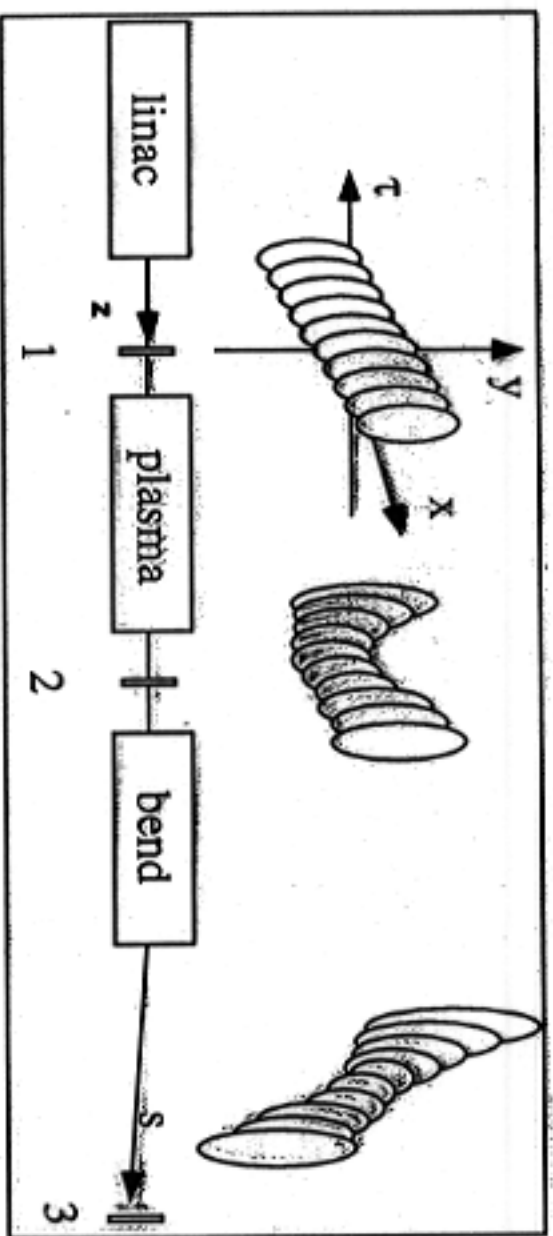


- Incoming head-tail "tilt" is amplified and distorted after 1.5 meter plasma

A. Geraci (SLAC)

Nominal Parameters

Electrons in Bunch	4. E10	Dipole Perturbation
Bunch Length (σ_z)	0.7 mm	
RMS Emittance	6. E-5 m-rad	
RMS Beam size	1. E-5 m-rad	
Beam Energy	30. GeV	Transverse Centroid differs by
Peak Beam Density	1.45 E15 cm ⁻³	0.3 * (beam radius), e.g.
Plasma Density	3 E14 cm ⁻³	between beam head(-3 σ) and
Plasma Channel Radius	1 mm	tail(+3 σ), i.e linear "tilt"
Plasma Length	1.5 m	



Courtesy A. Geraci and D. Whittum, SLAC

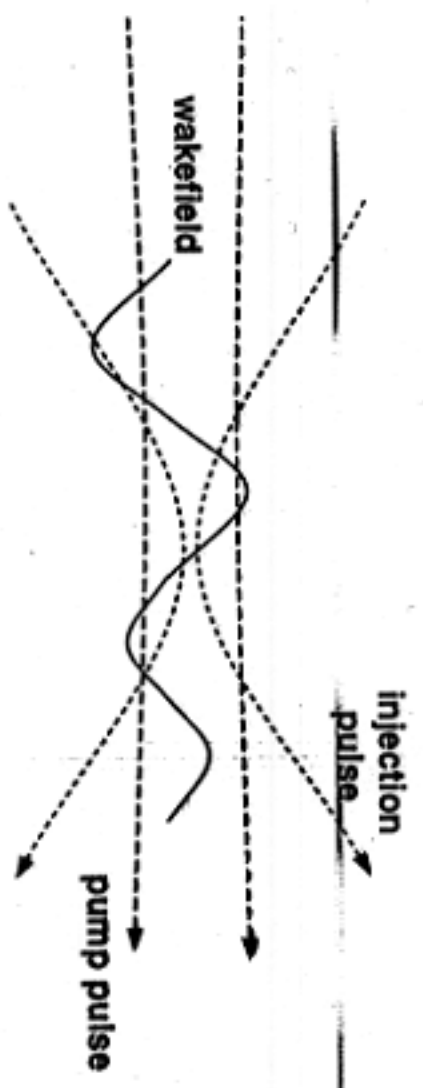


Injectors into high-frequency structures

- fs laser and RF gun+chicane (e.g., UCLA Neptune Lab)
- Optical injection concepts
- Inverse free-electron laser

LILAC: Laser Injected Laser Accelerator

- Laser wakefield accelerators have ultrashort acceleration buckets (10-50 fs).
 - They require an injector that produces femtosecond duration electron bunches.
 - We proposed an all-optical single-stage plasma-cathode injector/accelerator
- D. Umstadter, J. K. Kim, and E. Dodd, Phys. Rev. Lett. 76, 2073 (1996).



Several potential advantages:

- higher first-stage accelerating fields
⇒ lower transverse emittance
- monoenergetic ultrashort electron bunches
- absolute synchronization between wakefield and electrons
- compact and economical

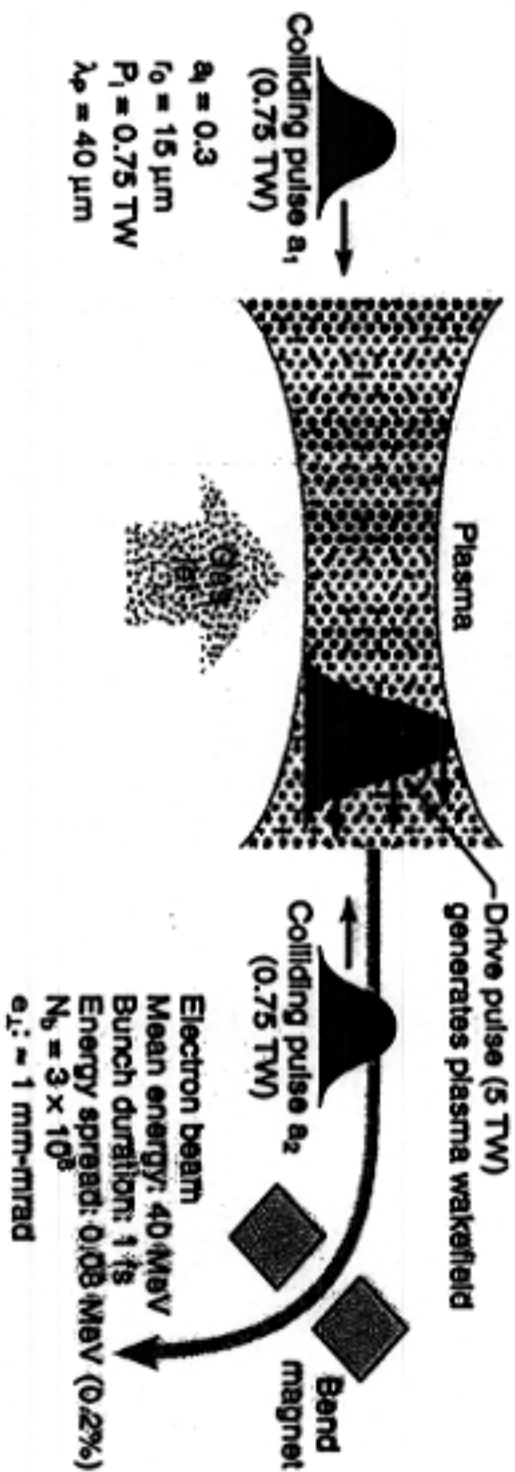
Simulations show LILAC has excellent beam characteristics:

- Single bunch $\tau_b = 9.0$ fs (3. μm) with $n_b = 1 - 4 \times 10^7$ per bunch (2-6 pC)
- Emittances $\epsilon_{\perp n} \approx 2\pi\text{mm} \cdot \text{mrad}$ and $\Delta E/E @ 100 \text{ MeV} = 2.0\%$



Optical injection using colliding pulse method

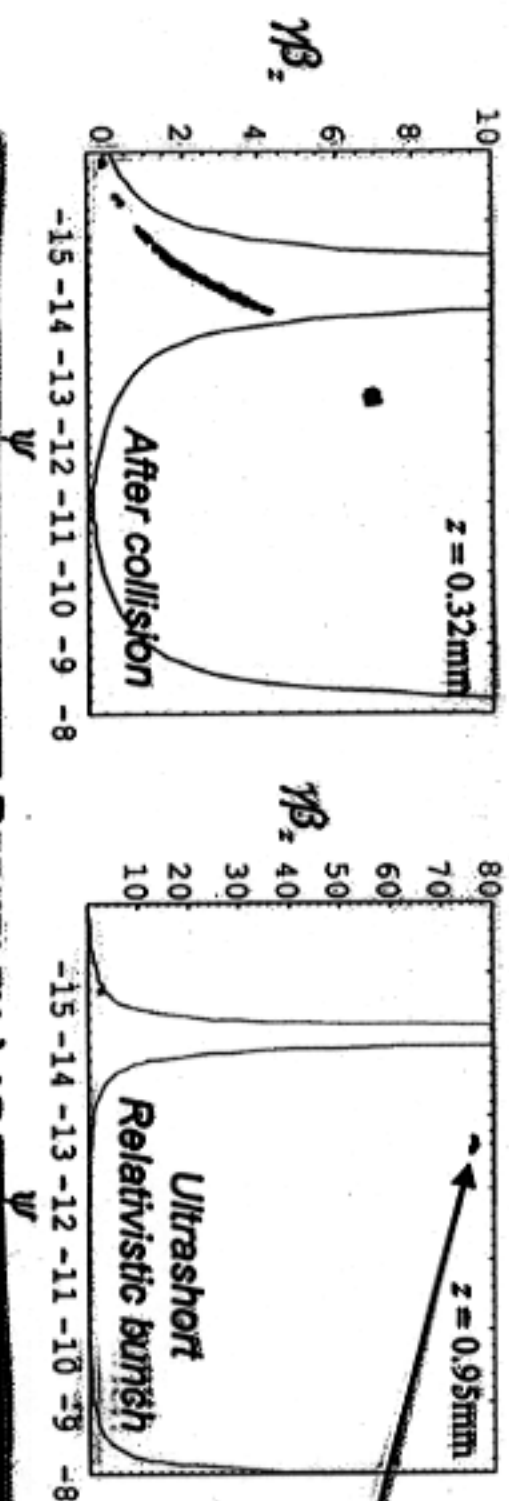
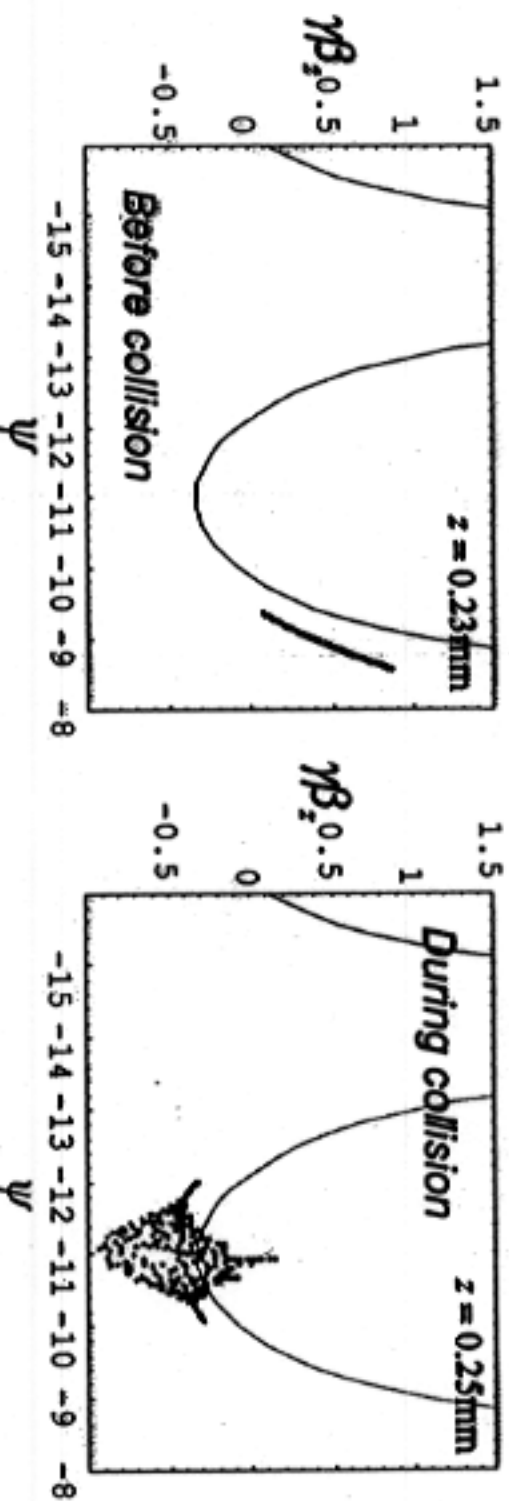
Injection pulses collide, producing a slow-moving beat wave that allows electrons to be trapped.





3-D simulation of colliding pulse: high quality bunch

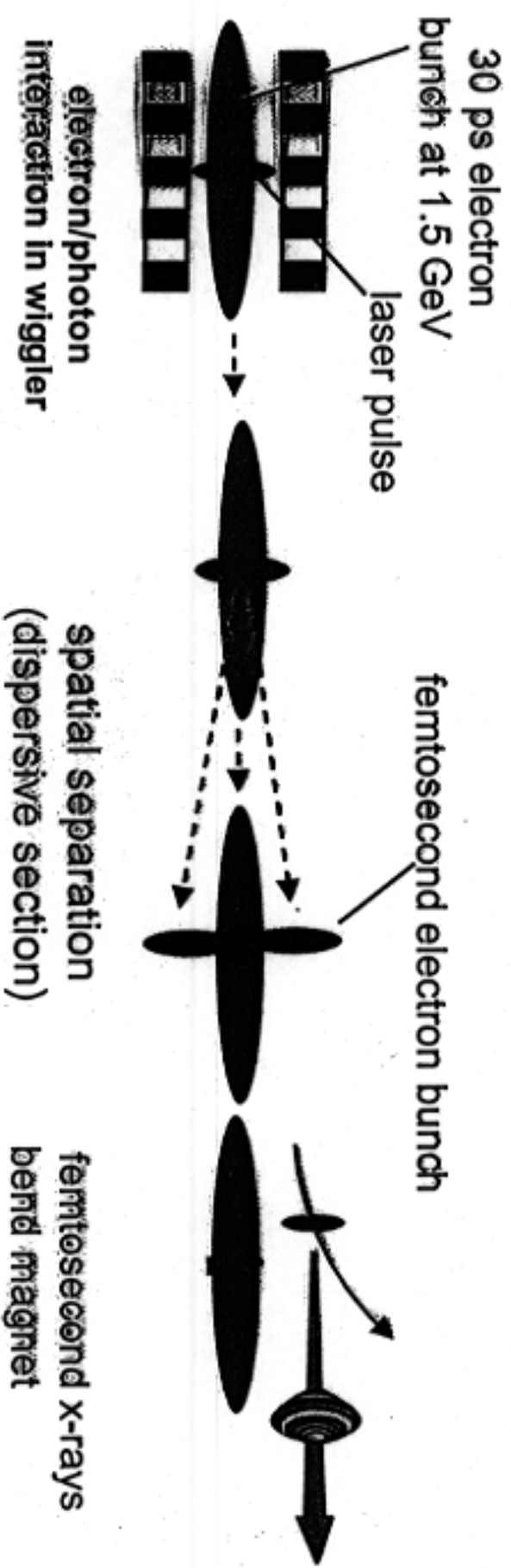
Longitudinal phase space evolution of distribution of plasma electrons



Mean energy = 39 MeV
Bunch duration = 1 fs
Energy spread = 0.08 MeV

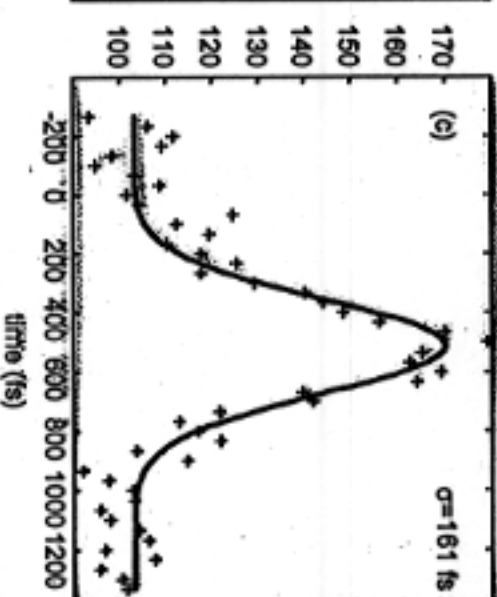
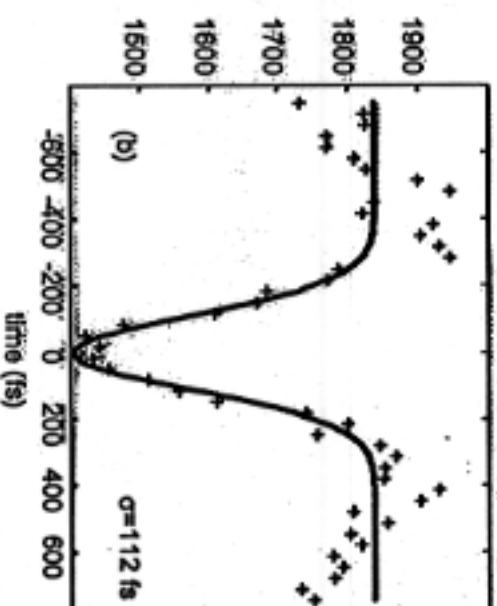
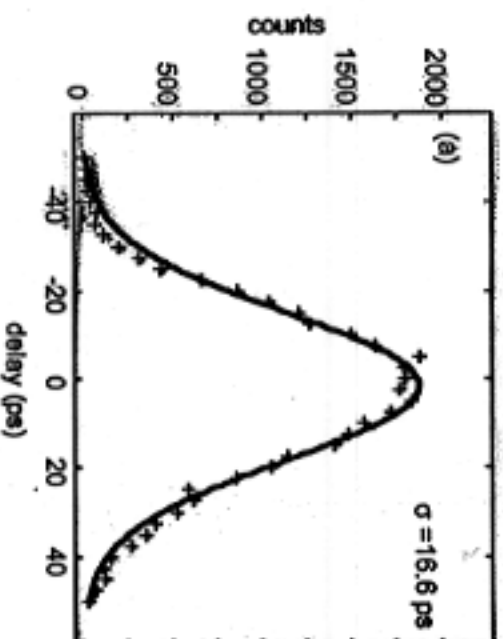
Generation of Femtosecond X-rays

A. Zholents, M. Zolotarev, PRL, 76, 916, 1996



$$t(\mathcal{E}), \quad r(\mathcal{E}), \quad r'(\mathcal{E}), \quad \omega(\mathcal{E})$$

Correlation Measurements of Visible Femtosecond Synchrotron Pulses



Zakharovs & Zolotarev '99

Also, BAL/ATF I FEL Benchmarking, 1999

Ber & Uv, et al



Plasma Lens

- Underdense

$$n_p < n_b$$

NON LINEAR RESPONSE

LINEAR W_L, W_{II}

- Overdense

$$n_b < n_p$$

LINEAR RESPONSE

NON LINEAR W_L, W_{II} (depend on R)

- Current Neutralization

$$R_p R_b > 1 \quad R_p \sigma_z \lesssim 1$$

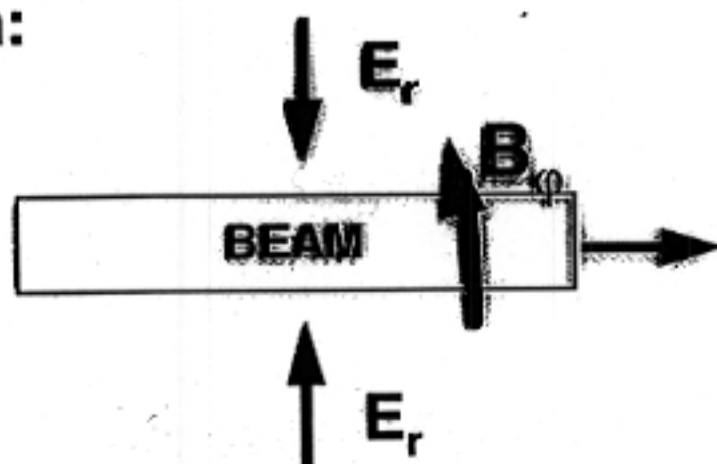
(adiabatic if $\gg 1$)

Groups: ANL, CCLA, LBNL, USC, SLAC

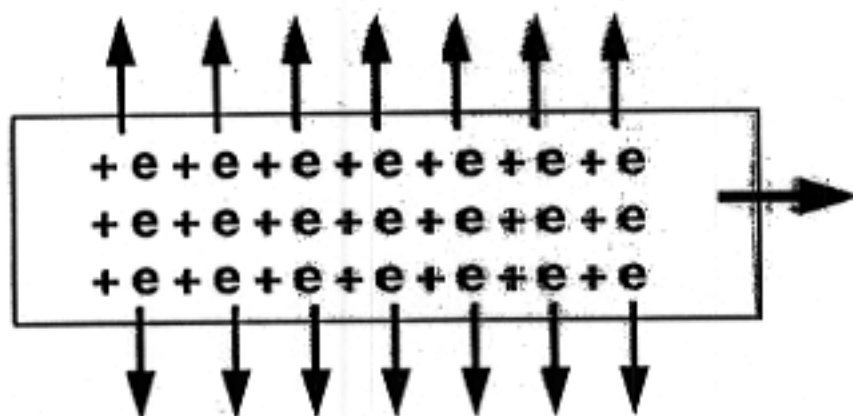
BERKELEY LAB

Ion Focusing

- In vacuum:



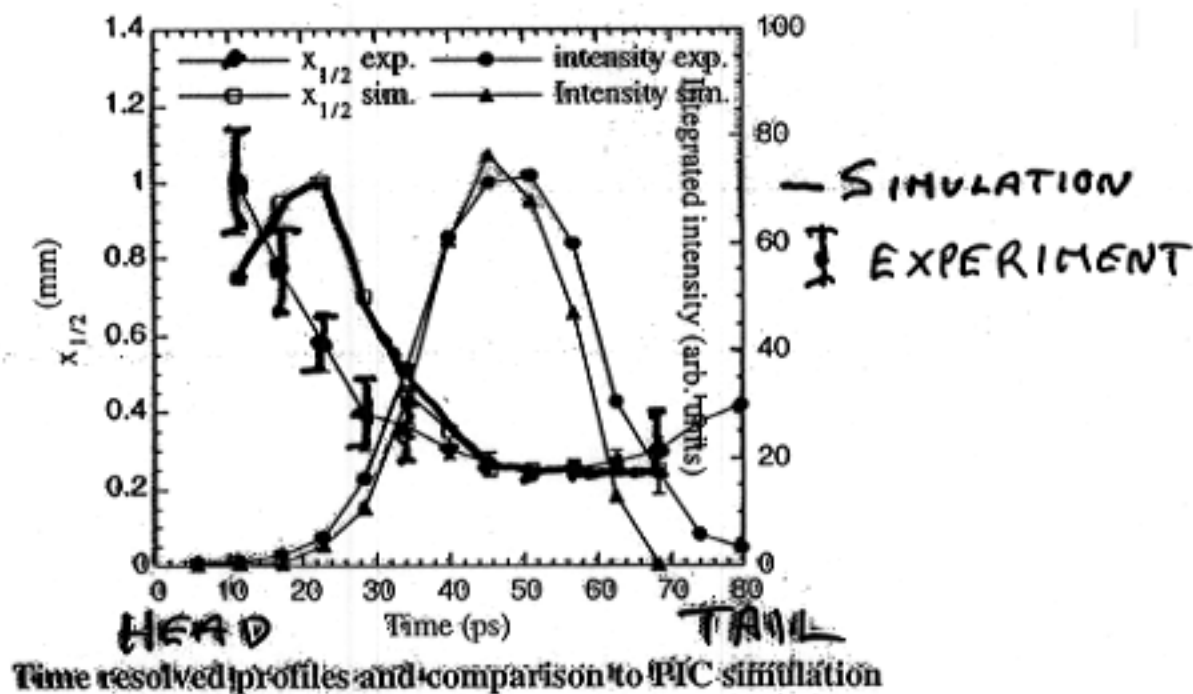
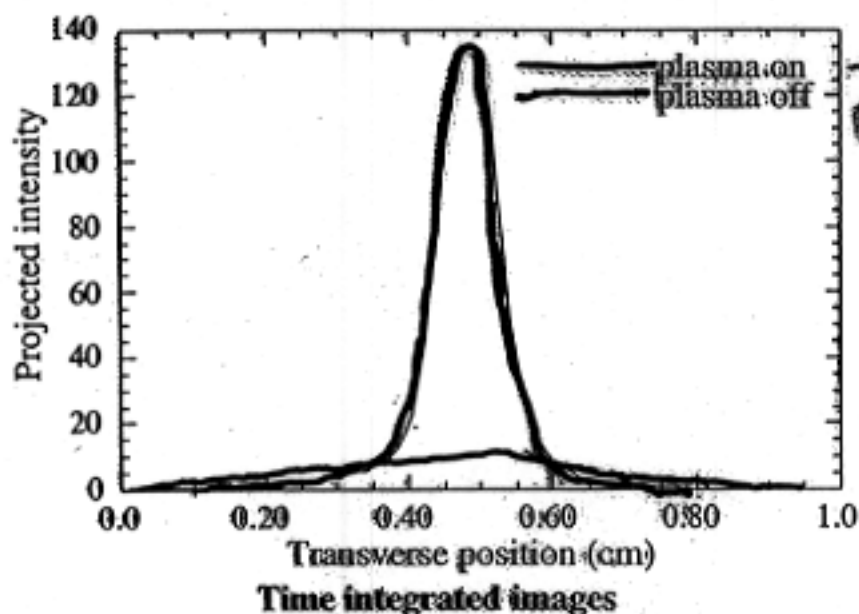
- In underdense case ($n_p < n_b$, ion focused regime):



$$F_r = \frac{4\pi e^2}{r} \int_0^r n_b(r)(1-\beta^2) r dr - 2\pi e^2 n_p r$$

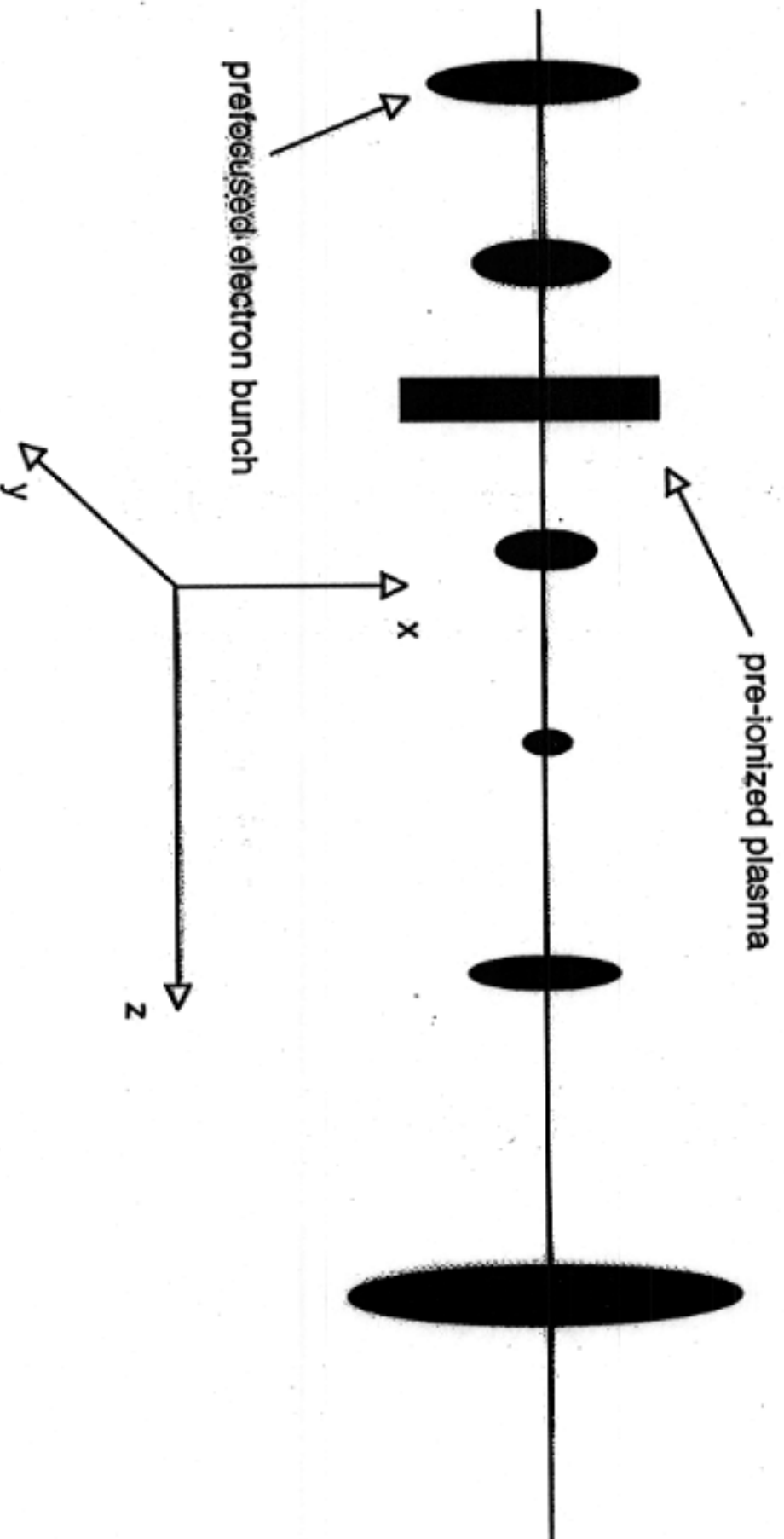
PROPAGATION EXPERIMENTS

- Time-resolved (psec resolution) and time integrated measurement of beam profiles at plasma exit. Near equilibrium confinement over 12 β -functions!



RETURN CURRENT PHYSICS LBNL/BTF (ALS INJECTOR)

k_{prof} variable, $n_b < n_p$



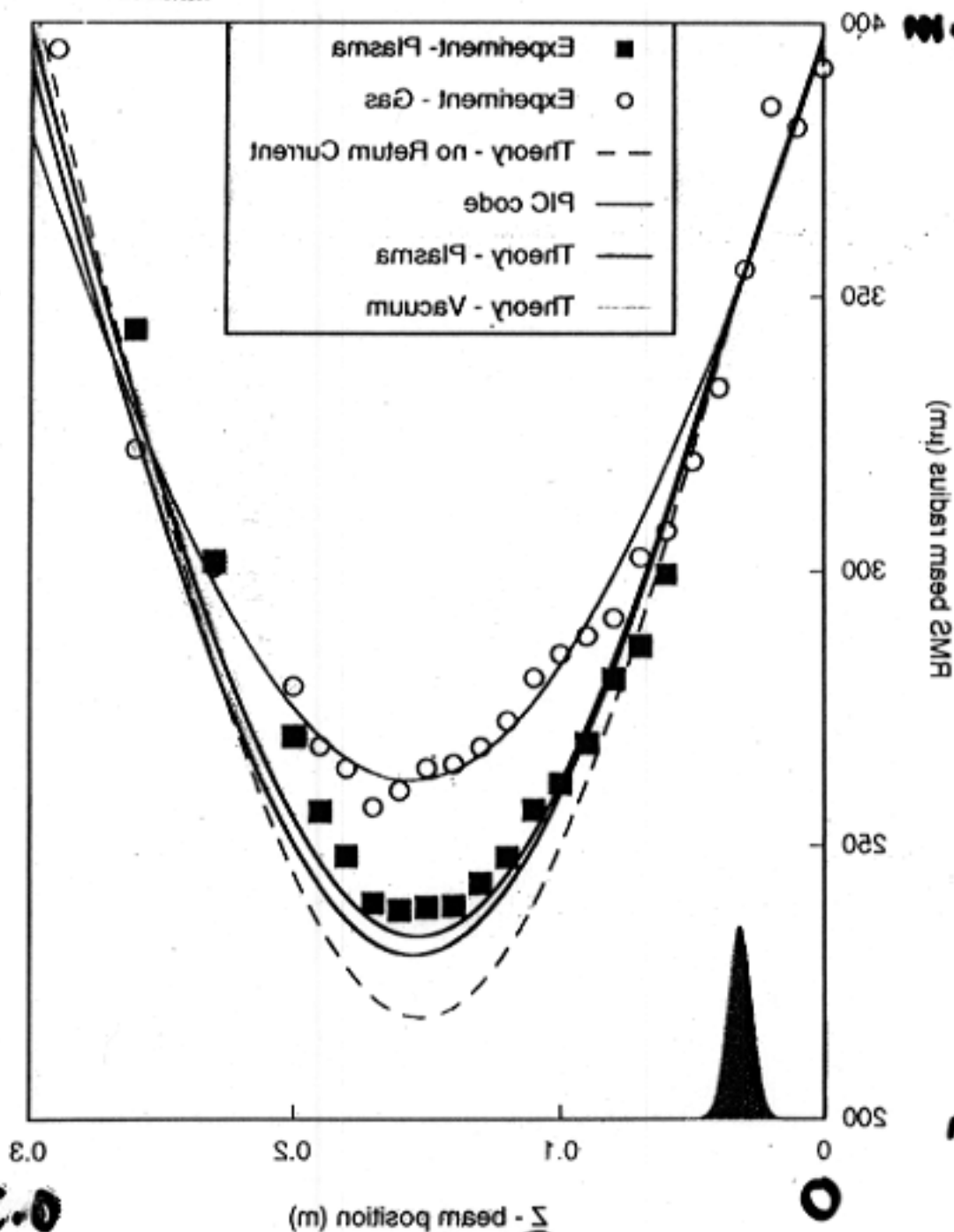
Katy Beckhaus Th.
 (45. w)

R. Govil (25. w)
 (+ W. Leemann)

NOTE: PIC SIMULATION; INVERSE FOR
& EXP. AREA.

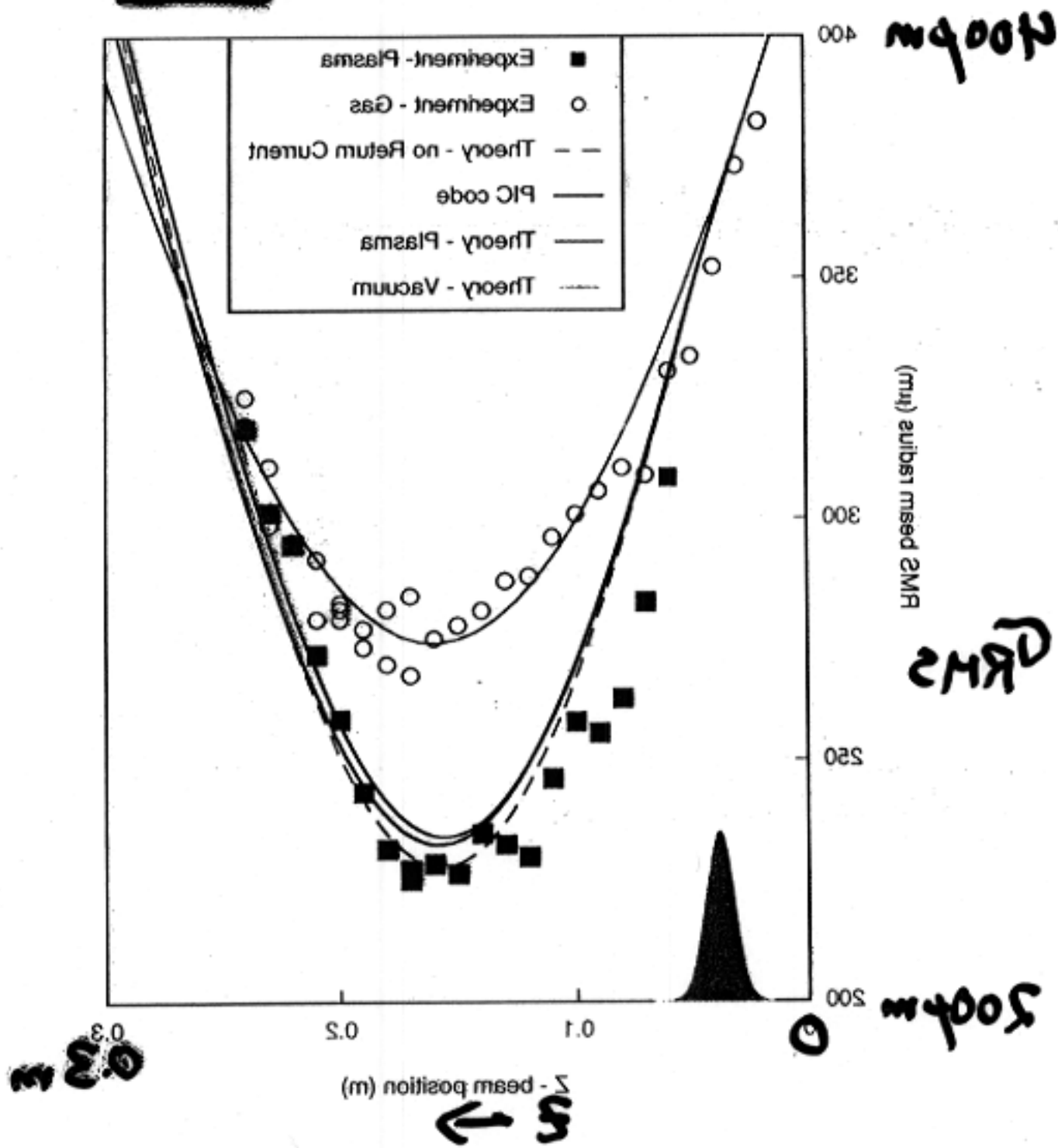
0.3 m

5 →



Plasma Lens with Large Return Current ($k_p = 1.1$)

Plasma Lens with Small Return Current ($k_p = 0.33$)





Plasma Lens

- LONG LENS $k_p \sigma_z^2 > 1$
- $n_p > n_b$ space-charge neutralization

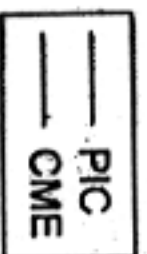
- $V_z \approx c$ and constant

$$X_{rms}'' = \frac{\xi^2}{X_{rms}^3} + \frac{1}{\gamma m c^2} \frac{\langle x F_x \rangle}{X_{rms}}$$

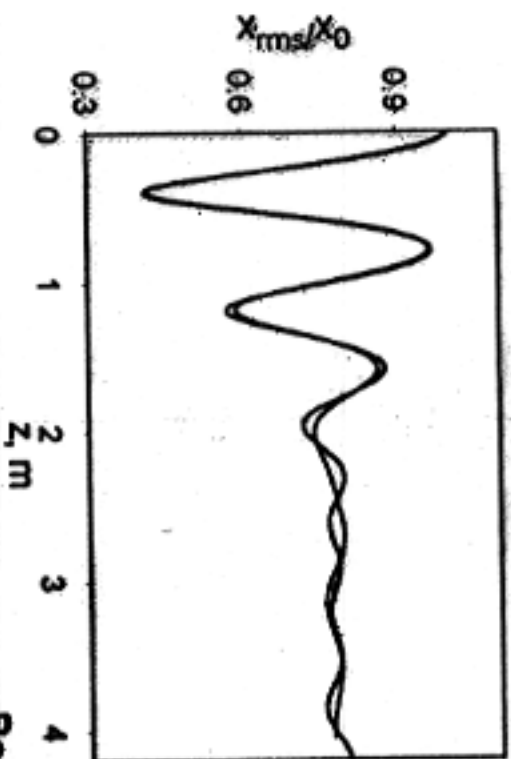
SAME E QNS FOR
SELF-GRAVITATING
SLAB!

$$E_x = -q c B_y \quad B_y = q c \mu_0 \int_0^x dx' n_b(x')$$

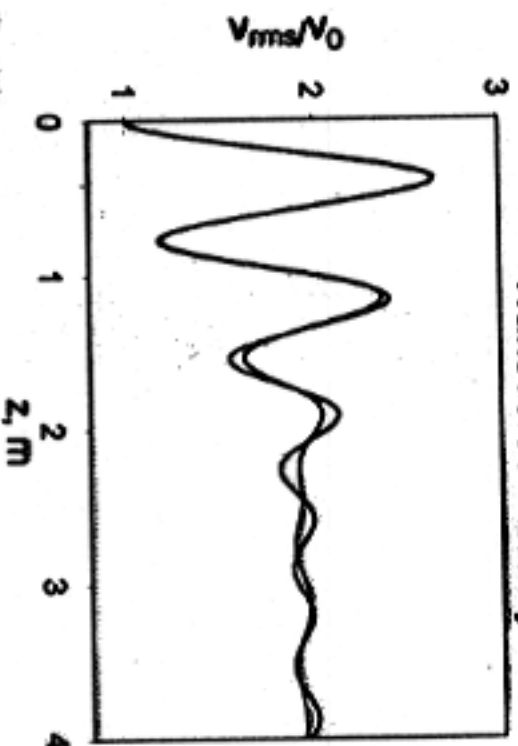
Comparison with PIC Simulation (RMS)



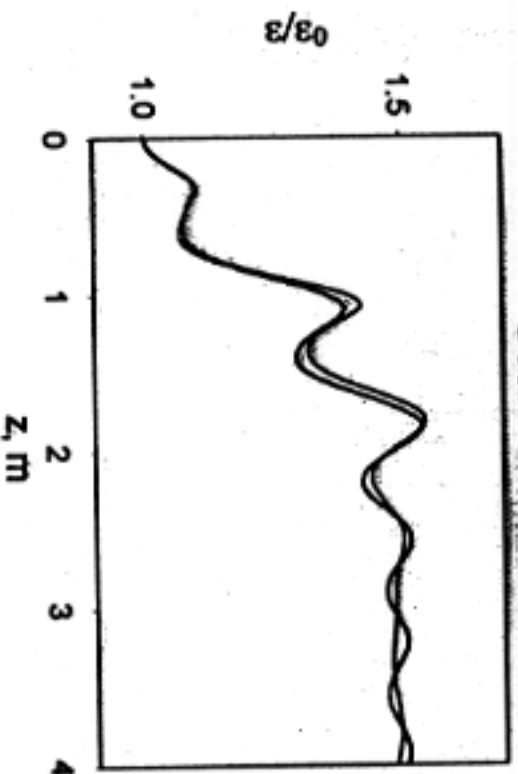
Beam Size



Transverse Velocity

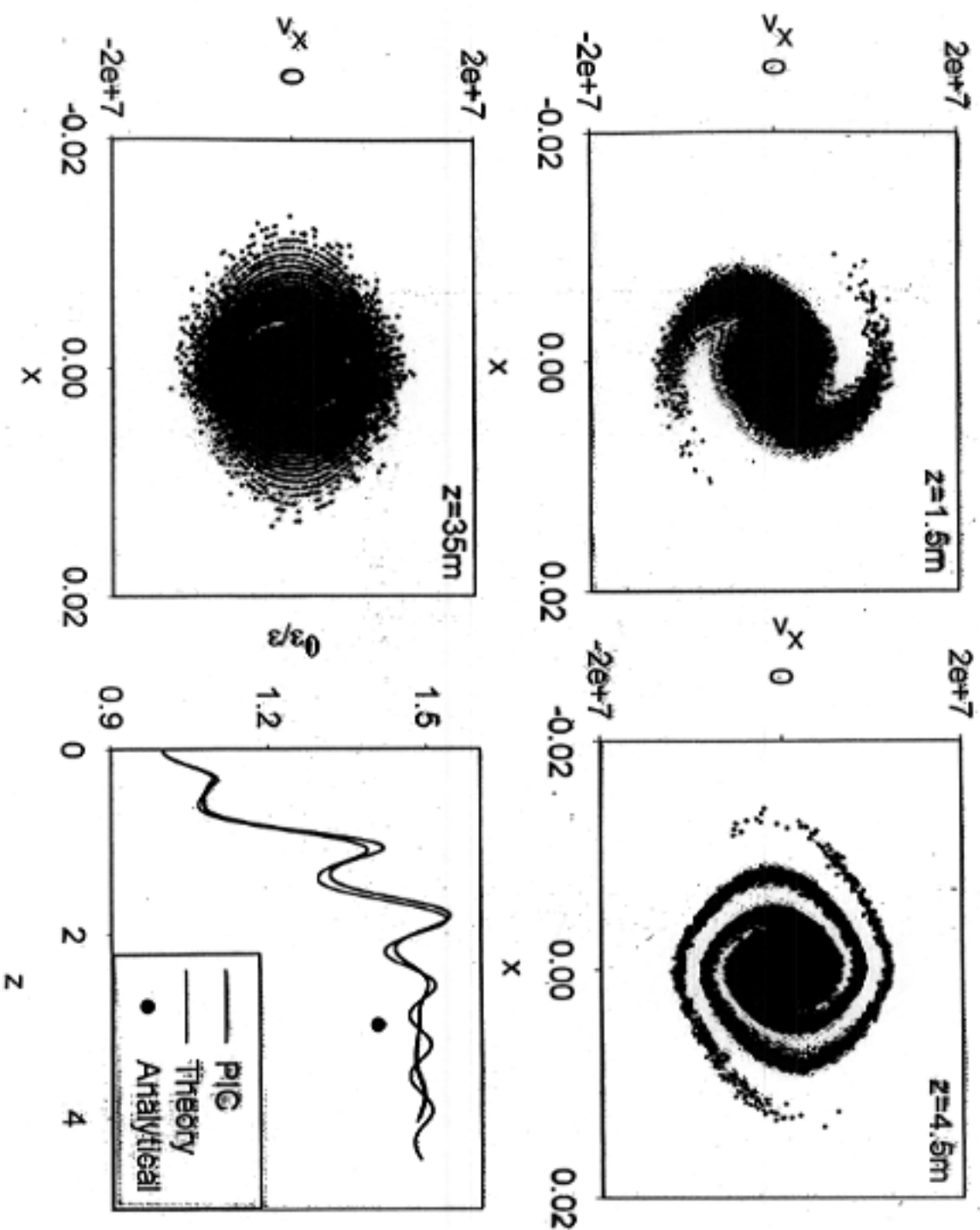


Beam Emittance



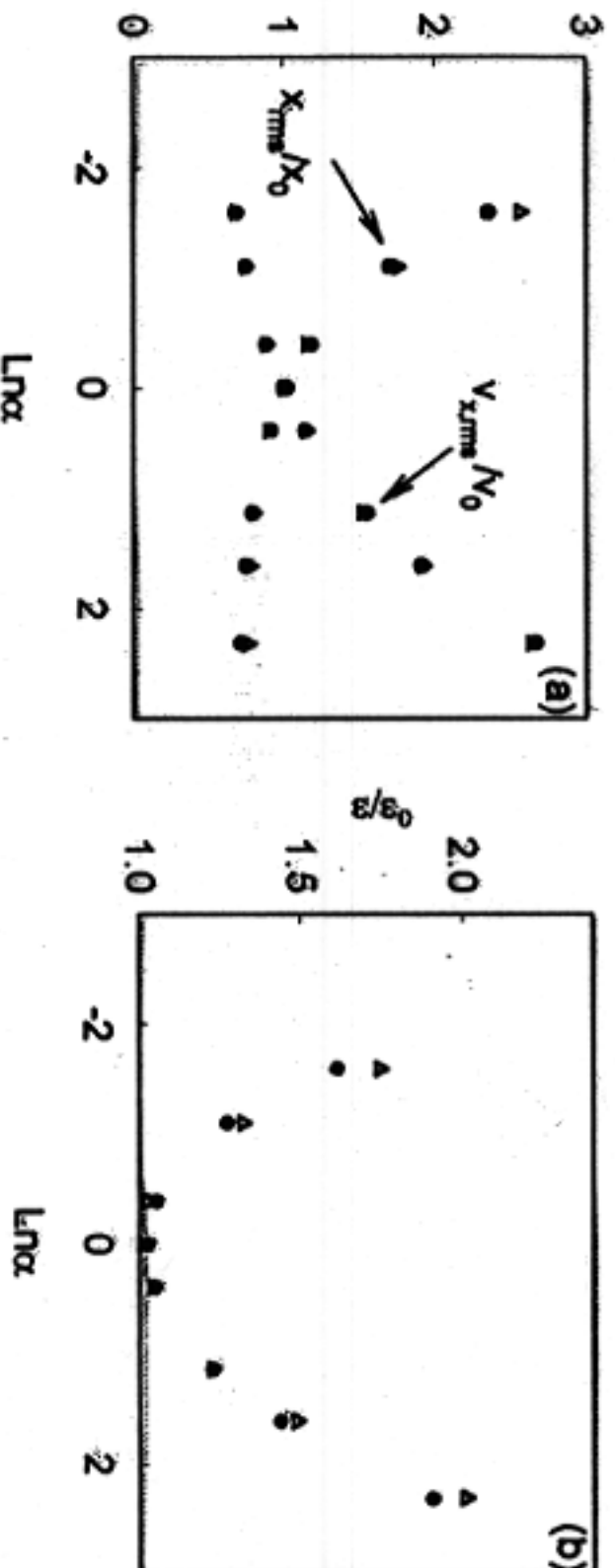
$$k_{max} = l_{max} = 20$$

Beam Transport in Slab Plasma ($k_{\perp} \sigma_z \gg 1$)





RMS Equilibrium in a long overdense plasma lens: Analysis and Simulation





Important considerations for accelerating structures (plasma and others)

- Method of excitation: ebeam or laser.
- Number of oscillations (Q) of wake
 - ➔ limits number of bunches—more charge/bunch at fixed efficiency
- Gradient/stored energy (loss factors k_n)
 - ➔ Determines beam dynamics and stored energy requirements
- Dark Current (*Nonlinear kinetic plasma physics*)
 - ➔ Must be controlled for high-energy applications
 - ➔ High gradient ($\sim 100\text{GeV/m}$) plasma experiments operate in this regime. Earlier results (e.g., Joshi et al) achieved $\sim 1.5\text{GeV}$ without dark current.
 - ➔ Limitations unclear



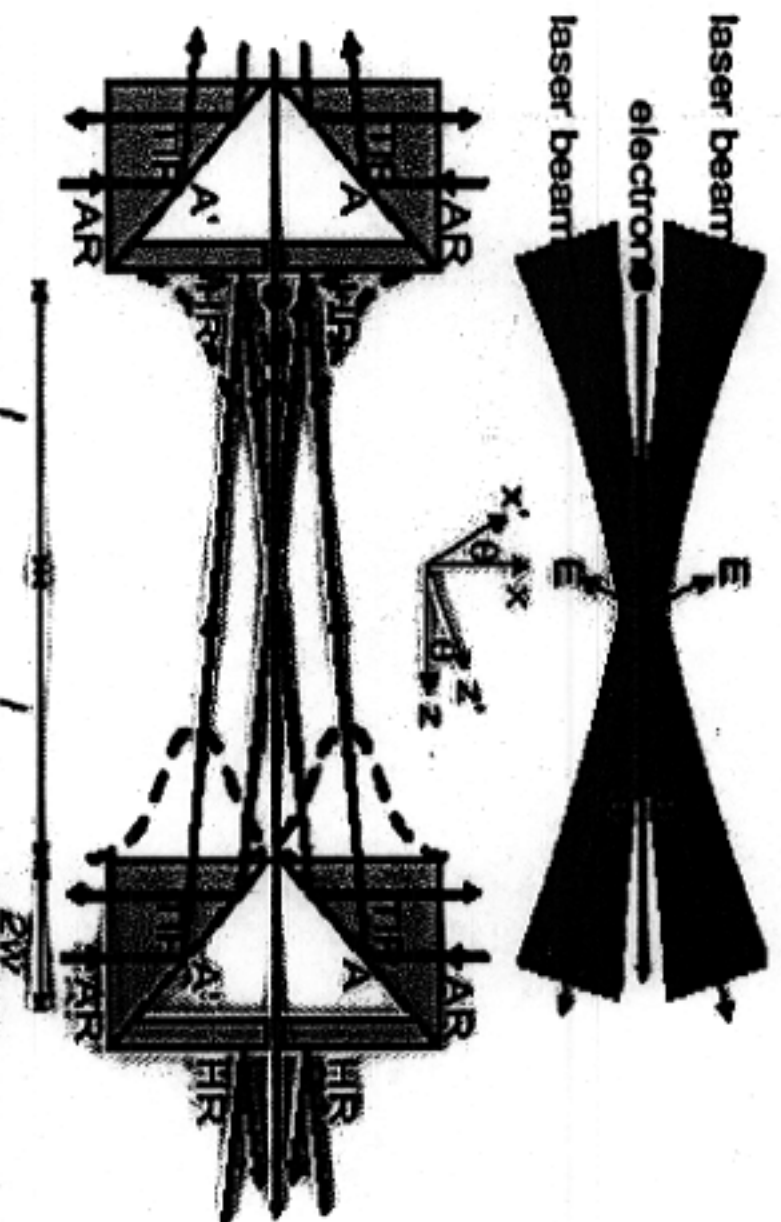
Summary: Plasma-based concepts

- Significant progress in plasma-based concepts: Experiment, theory, simulations, diagnostics.
- Next few years:
 - ➔ LASER driver: High power (10^{18} W/cm²) propagation ($L \gg Z_r$) and wake generation in channels
 - ➔ BEAM driver: Ultra-relativistic dynamics in plasmas (SLAC E150/E157)
 - ➔ Short bunch generation: Optical injection
 - ➔ Dark current limits
 - ➔ *Plasma structure design*--Femtosecond engineering
- For the future: High energy, beam quality and intensity, staging, efficiency



Vacuum Acceleration

- Laser (no plasma): Construct and excite small scale overmoded structures. Directly use the intense electric field of the laser.
- Issues: very short bunch length, structure damage, wakefields.
- A schematic (Laser Acceleration Experiment, Byer et al.). Bunch from Stanford FEL. Accelerate 330keV at $\sim 1\text{GeV/m}$:





Summary

- Many ideas for solving critical problems towards building high-energy colliders. There are no obvious "show stoppers" for any of the schemes--but *strong* opinions on the correct direction to proceed. Even if all fail, a clever student may come up with a better idea!
- Significant research ongoing involving extended collaborations of laboratories and universities
- The timescale for investigating these concepts is longer than most would like.
- The next generation of colliders will not be the last!