

QCD Corrections to $t\bar{t}H$ Production

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Workshop on the Future of Higgs Physics
FNAL, May 2001

- Motivations (discovery, precision studies)
- $p\bar{p} \rightarrow t\bar{t}H$ (Run II)
- $pp \rightarrow t\bar{t}H$ (LHC)
- $e^+e^- \rightarrow t\bar{t}H$ (Tesla/NLC)

Crucial to know first order QCD corrections

- increase/decrease cross section considerably
- reduce μ_r -dependence substantially

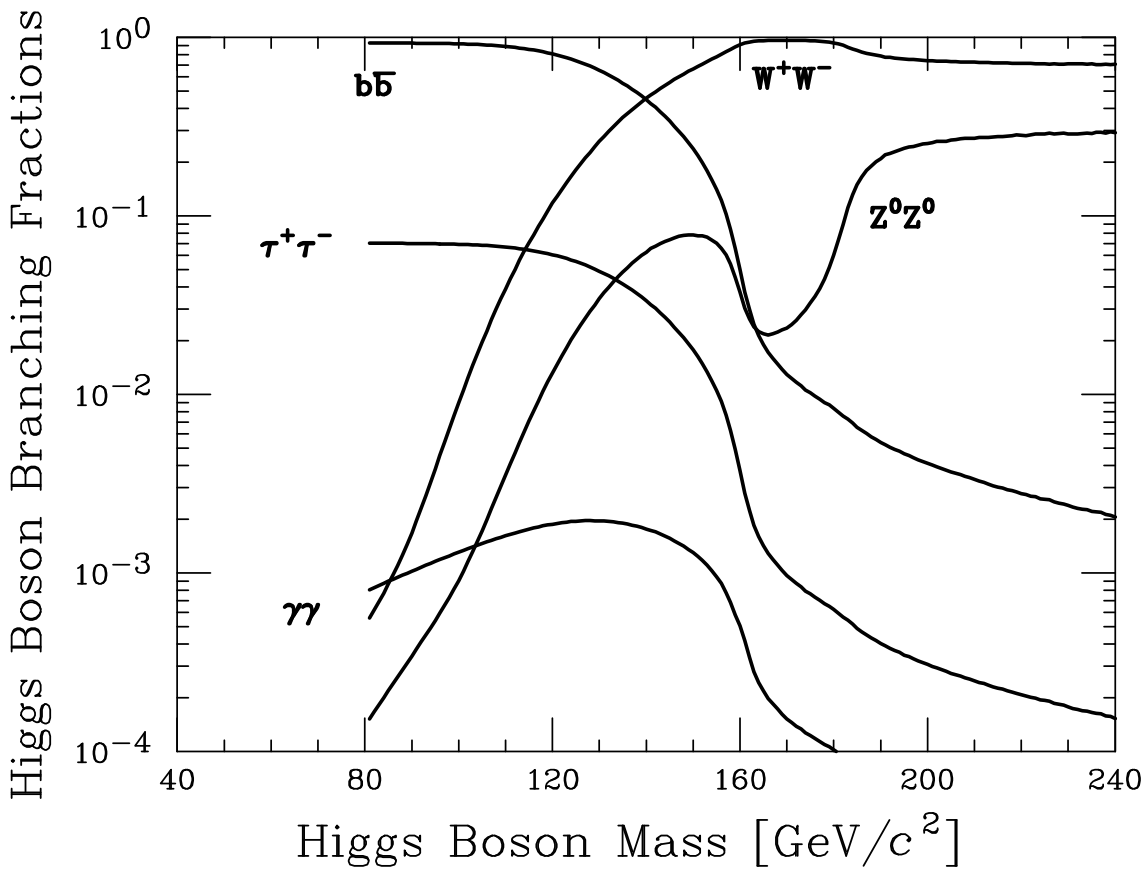
(with S. Dawson)

MOTIVATIONS,1

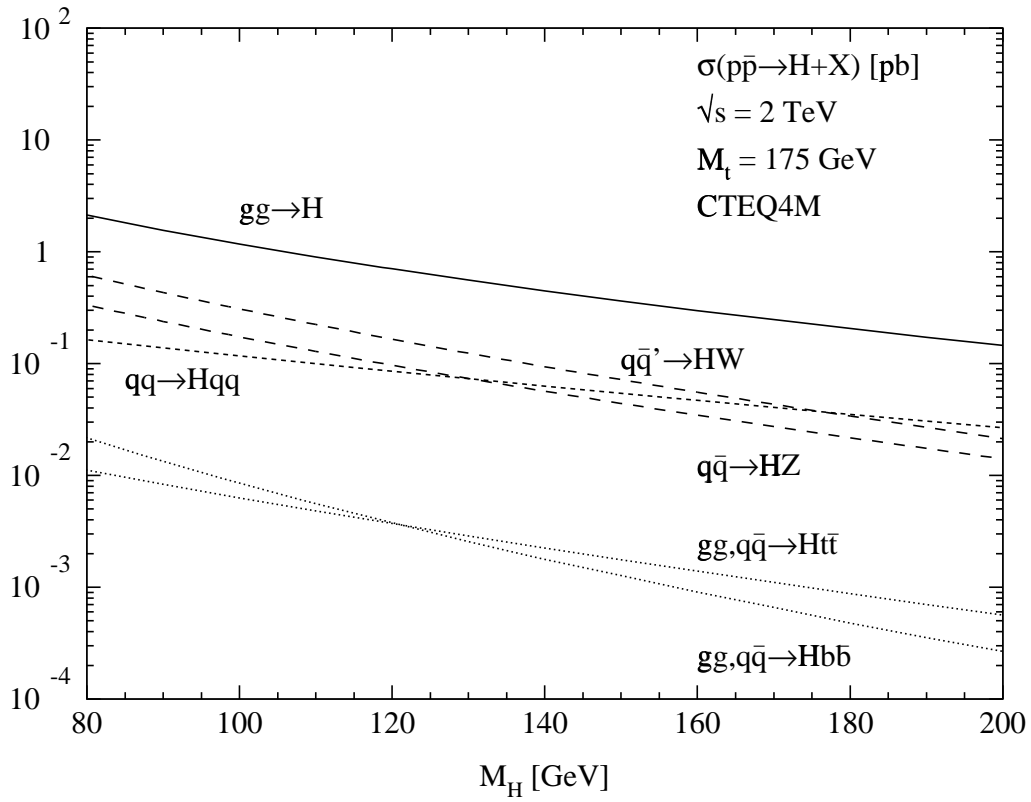
$110 < M_H < 140$ GeV : Crucial Region for future colliders (Tevatron, LHC, NLC)

BUT

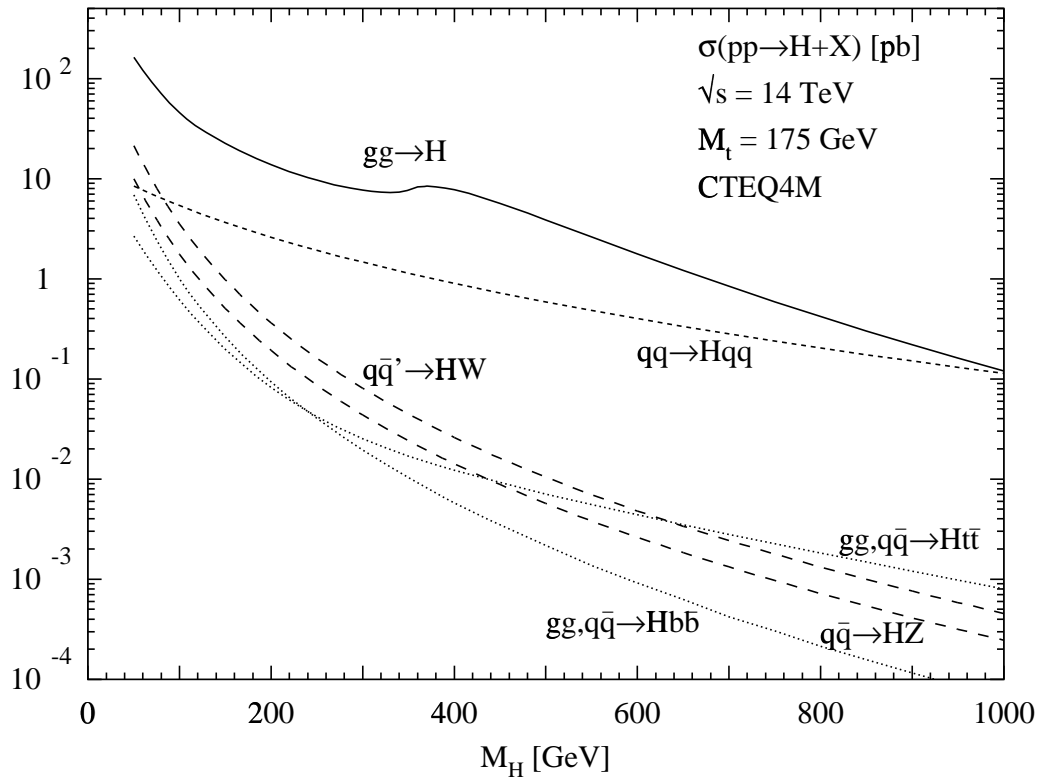
For **Hadron Colliders** this is a **problematic region**: non leading production modes will also have to be used to confirm Higgs signal



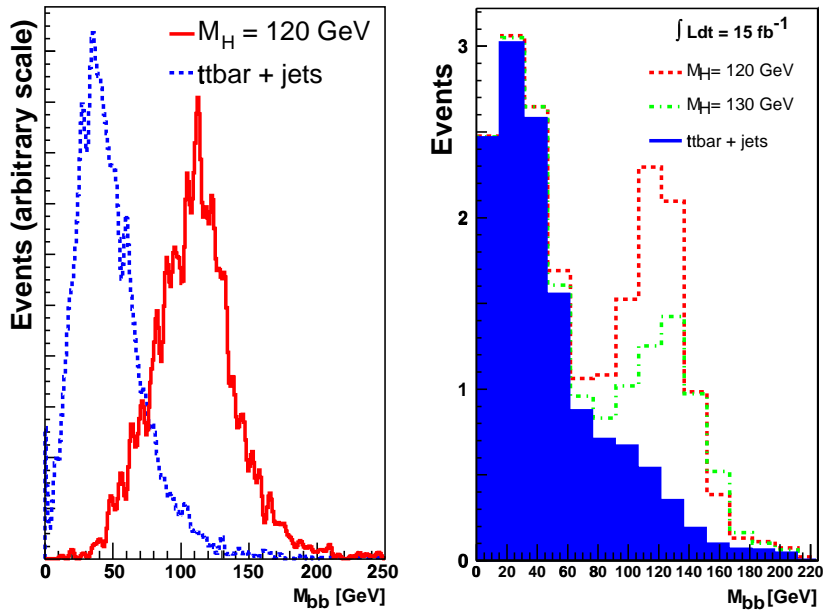
Tevatron, SM Higgs



LHC, SM Higgs



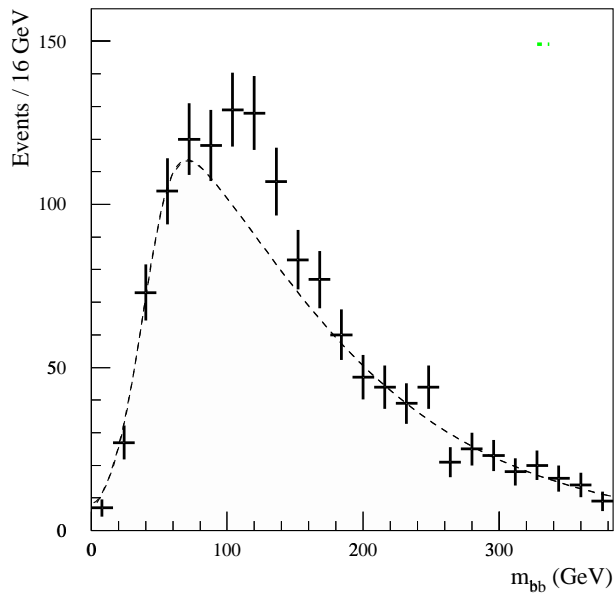
Tevatron, $p\bar{p} \rightarrow t\bar{t}h_{SM}$



Goldstein et al.
hep-ph/0006311

LHC, $pp \rightarrow t\bar{t}h_{SM}$

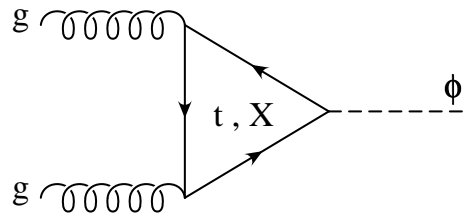
LHC top working group



Precision studies : measure the Top Yukawa coupling in $pp, e^+e^- \rightarrow t\bar{t}\phi$

- **direct** measurement

($\phi \rightarrow gg$ or $gg \rightarrow \phi$: possible spurious contributions from heavy new particles)



- **high sensitivity** to the $t\bar{t}\phi$ coupling: probe **new physics**
- **spectacular signature** : $W^+W^-b\bar{b}b\bar{b}$

NLC :

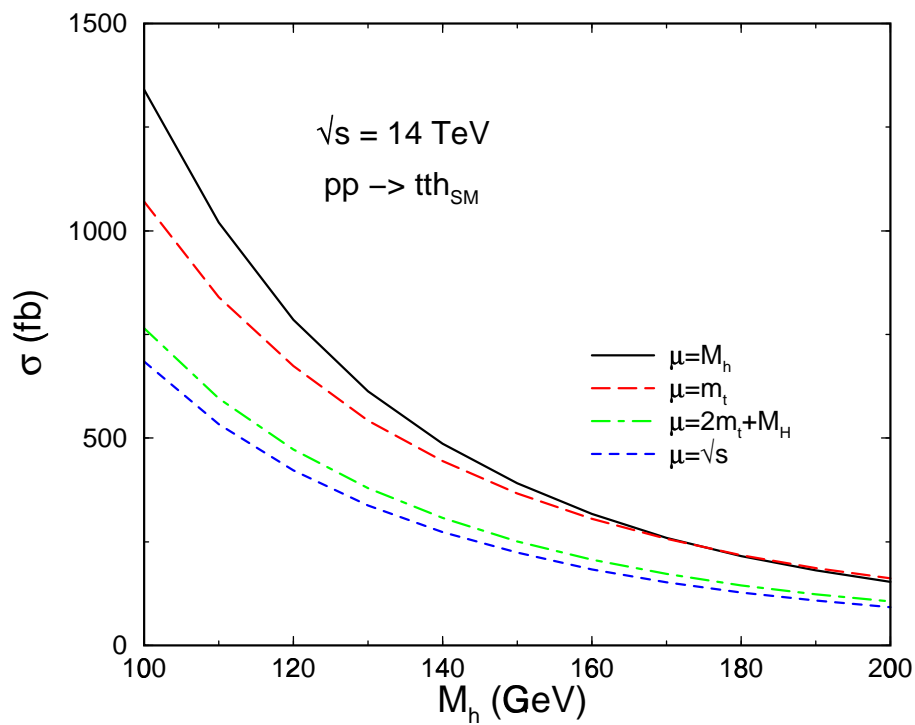
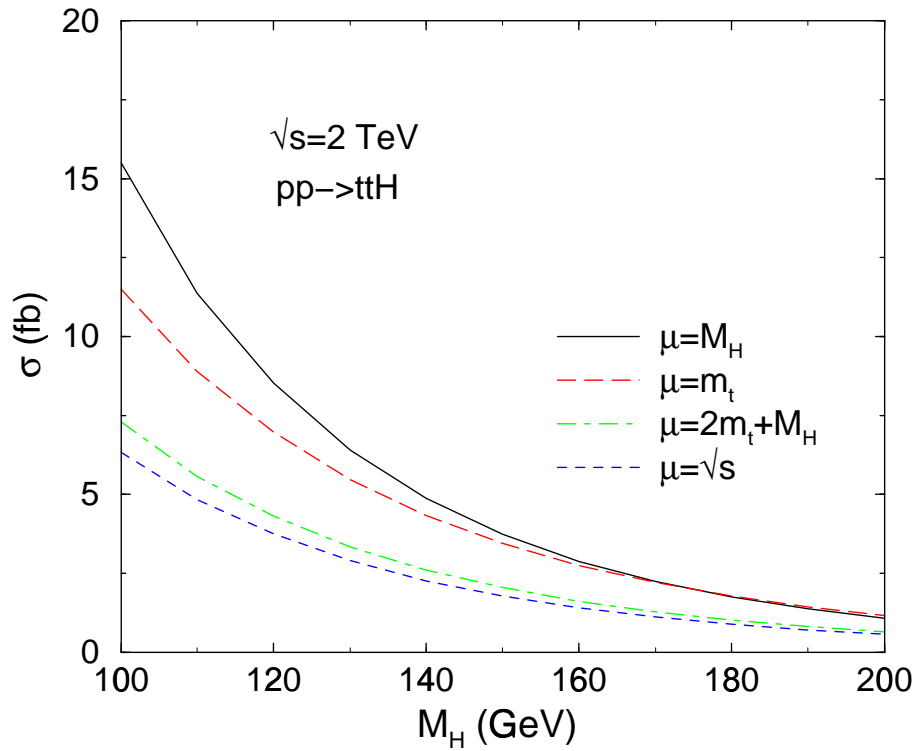
$$\frac{\delta g_{tth}}{g_{tth}} \simeq 5.5\%$$

LHC :

$$\frac{\delta g_{tth}}{g_{tth}} \simeq 16\%$$

Need for NLO calculation

LO calculation have very strong μ -dependence



Status of QCD corrections

- $\sigma(e^+e^- \rightarrow t\bar{t}h)$ calculated at $O(\alpha_s)$:
 - $\sqrt{s} = 500$ GeV : K=1.4-2.4
 - $\sqrt{s} = 1$ TeV : K=0.8-0.9

Dittmaier, Krämer, Liao, Spira, Zerwas
L.R., S.Dawson

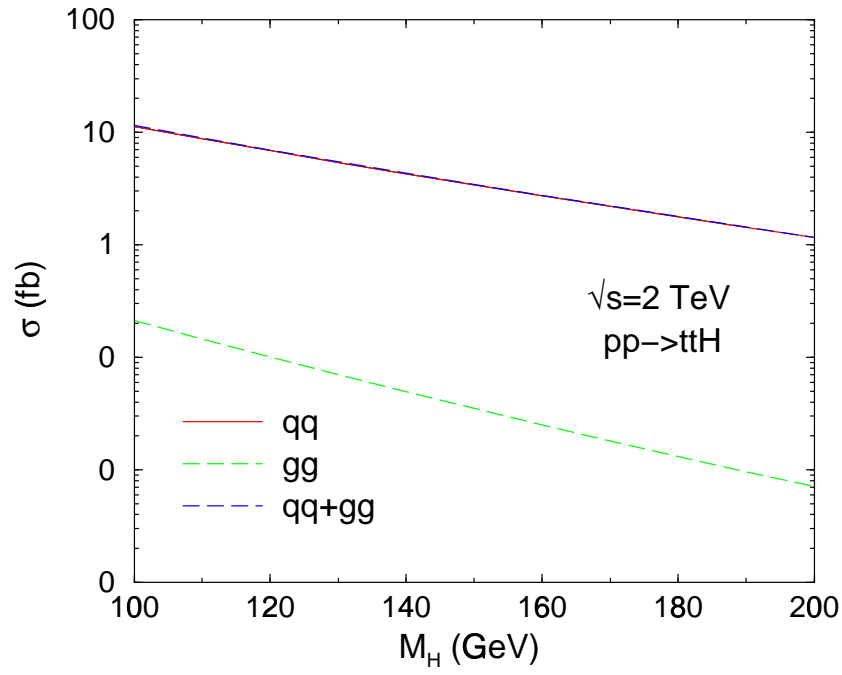
- $\sigma(p\bar{p} \rightarrow t\bar{t}h)$ calculated at $O(\alpha_s^3)$

Beenakker, Dittmaier, Krämer, Pluemper, Spira,
Zerwas
S.Dawson, L.R.

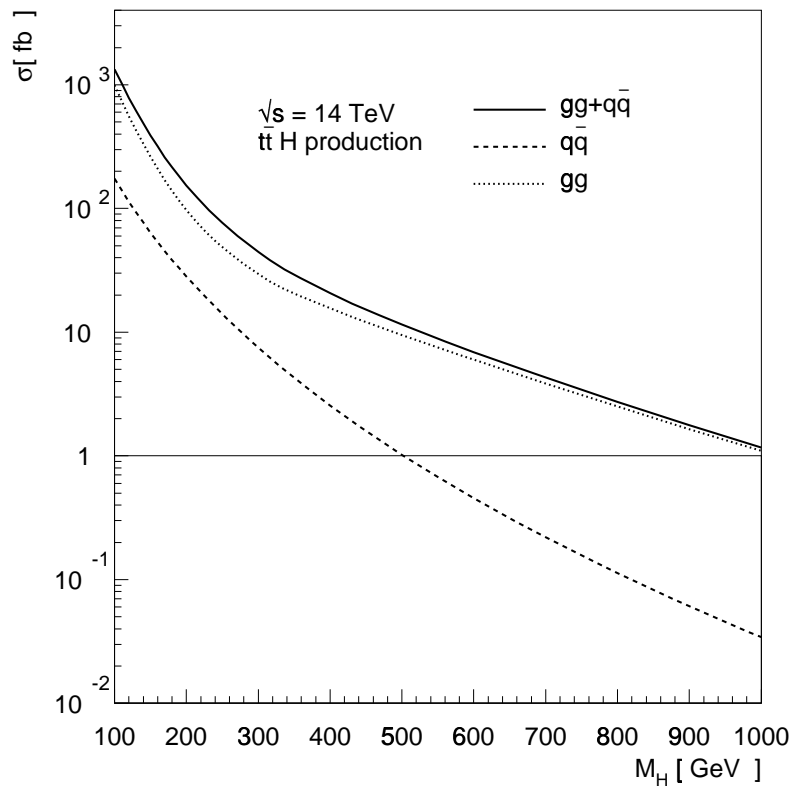
- $\sigma(pp \rightarrow t\bar{t}h)$ at $O(\alpha_s^3)$: work in progress

Beenakker, Dittmaier, Krämer, Pluemper, Spira,
Zerwas
S.Dawson, L. Orr, L.R., D. Wackerroth

$pp \rightarrow ttH$: dominant channel $q\bar{q} \rightarrow ttH$



$pp \rightarrow t\bar{t}H$: dominant channel $g\bar{g} \rightarrow t\bar{t}H$



General Approach

NLO cross section:

$$\sigma_{\mathcal{H}_1\mathcal{H}_2} = \sum_{a,b} \int dx_1 dx_2 \mathcal{F}_a^{\mathcal{H}_1}(x_1) \mathcal{F}_b^{\mathcal{H}_2}(x_2) \sigma_{ab}^{NLO}(x_1, x_2)$$

where

$$\sigma_{ab}^{NLO} = \sigma_{ab}^{LO} + \alpha_s \delta\sigma_{ab}^{NLO} + O(\alpha_s^2)$$

NLO correction $\delta\sigma_{ab}^{NLO}$ made off:

- one loop **virtual** corrections to $q\bar{q} \rightarrow t\bar{t}H$
- **real** one gluon emission: $q\bar{q} \rightarrow t\bar{t}H + g$

We use **Phase Space Slicing** as proposed by:

Giele and Glover, PRD 46 (1992) 1980

Giele, Glover, and Kosower NPB 403 (1993) 633

extended to **massive quarks** by

Keller and Laenen, PRD 59 (1999) 114004

- **Cross** all partons to final state:

$$q\bar{q} \rightarrow t\bar{t}H \text{ becomes } H \rightarrow q\bar{q}t\bar{t}$$

- Compute **virtual corrections** $\rightarrow \sigma_{virtual}$
 - renormalize **UV divergences**
 - calculate **IR divergences**
- Compute **real gluon emission** using **PSS**
 - introduce a **cut-off** parameter s_{min} (GeV²)
 - partons are considered **soft/collinear** if

$$s_{ij} = 2p_i \cdot p_j < s_{min}$$

$\sigma_{hard} \rightarrow$ all $s_{ij} > s_{min}$, numerical

$\sigma_{ir} = \sigma_{soft} + \sigma_{collinear} \rightarrow$ analytic

- **Cross** $q\bar{q}$ to initial state, H to final state:
 - flip signs of s_{ij} invariants accordingly
 - add $\sigma_{crossing} \rightarrow$ use **crossing functions**

\Downarrow

$$\sigma_{real} = \sigma_{hard} + \sigma_{ir} + \sigma_{crossing}$$

$$\delta\sigma^{NLO} = \sigma_{virtual} + \sigma_{real}$$

GGK+KL \rightarrow systematic approach to factorization of soft+collinear divergences.

- use color ordered amplitudes

$$\begin{aligned} \mathcal{A}(H \rightarrow q\bar{q}t\bar{t} + g) &\simeq \delta_{f_t f_{\bar{t}}} \delta_{f_q f_{\bar{q}}} \cdot \\ &\left(\delta_{c_t c_{\bar{q}}} T_{c_q c_{\bar{t}}}^a B_1(t, \bar{t}, q, \bar{q}, k) + T_{c_t c_{\bar{q}}}^a \delta_{c_q c_{\bar{t}}} B_2(t, \bar{t}, q, \bar{q}, k) \right) \\ - \frac{1}{N} &\left(\delta_{c_t c_{\bar{t}}} T_{c_q c_{\bar{q}}}^a B_3(t, \bar{t}, q, \bar{q}, k) + T_{c_t c_{\bar{q}}}^a \delta_{c_q c_{\bar{q}}} B_4(t, \bar{t}, q, \bar{q}, k) \right) \end{aligned}$$

- calculate the soft and collinear approximation of

$$\int dPS_g |\mathcal{A}(H \rightarrow q\bar{q}t\bar{t} + g)|^2$$

keeping both leading and sub-leading terms in N , and find

$$\sigma_{ir} = \sigma_{soft+collinear} = \int dPS_3 (S_F + C_F) |A_{LO}|^2$$

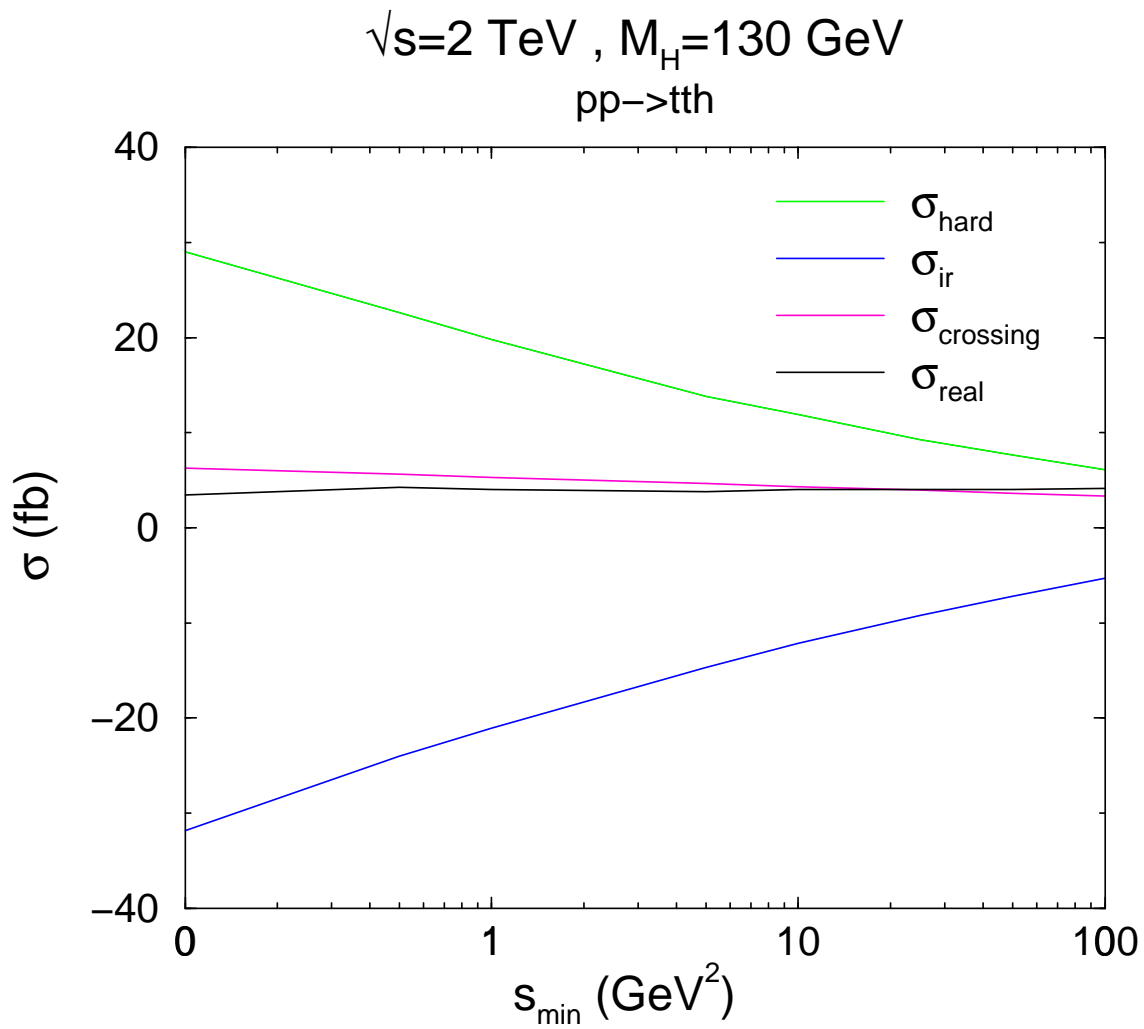
where both S_F and C_F are s_{min} -dependent.

- add $\sigma_{crossing}$ to correct for mismatch between initial state and final state collinear divergences \rightarrow extra s_{min} -dependence

$$\mathcal{F}_a^{\mathcal{H}}(x) = f_a^{\mathcal{H}}(x, \mu_F) + \alpha_s C_a^{\mathcal{H}}(x, \mu_F) + O(\alpha_s^2)$$

Crucial: check s_{min} -independence of σ_{real}

$$\sigma_{real} = \sigma_{hard} + \sigma_{ir} + \sigma_{crossing}$$



- Reduce each diagram in terms of scalar integrals of the form:

$$\int \frac{d^n k}{(2\pi)^n} \frac{1}{\left[(k^2 - m_1^2) \left((k + p_1)^2 - m_2^2 \right) \dots \right]}$$

→ three external massive particle

→ several massive internal propagators

- finite integrals: compute them numerically
- UV-divergent integrals: easy
- IR-divergent integrals : difficult

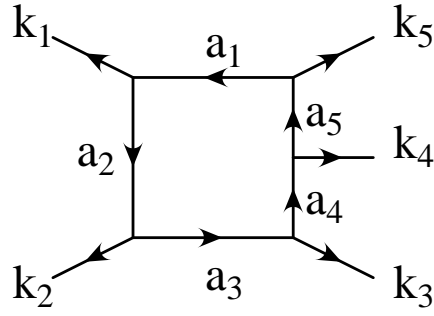
Example: pentagon scalar integral

$$E0 = \int \frac{d^n k}{(2\pi)^n} \frac{1}{N_1 N_2 N_3 N_4 N_5}$$

use method proposed by:

Bern, Dixon, and Kosower,

PLB 302 (1993) 299, NPB 412 (1994) 751



$$E0 = \frac{1}{16\pi^2} (4\pi\mu^2)^\epsilon \Gamma(3+\epsilon) \int d^5 a_i \frac{\delta(1 - \sum_{i=1}^5 a_i)}{[\mathcal{D}(a_i)]^{3+\epsilon}}$$

where

$$\mathcal{D}(a_i) = \sum_{i,j=1}^5 S_{ij} a_i a_j$$

$$S_{ij} = \frac{1}{2} (M_i^2 + M_j^2 - p_{ij}) \quad , \quad p_{ij} = k_i + k_{i+1} + \dots + k_{j-1}$$

scalar pentagon \rightarrow \sum scalar boxes

$$I_n = \frac{(-1)^n}{2} \left[\sum_{i=1}^n c_i I_{n-1}^{(i)} + (n - 5 + 2\epsilon) c_0 I_n^{6-2\epsilon} \right]$$

\Downarrow

$$E0 = -\frac{1}{2} \left[\sum_{i=1}^5 c_i D0^{(i)} + 2\epsilon c_0 D0^{6-2\epsilon} \right]$$

Conclusions

- $t\bar{t}H$ very interesting production mode at present and future colliders:
 - discover/confirm the Higgs
 - SM? New Physics?
- Crucial to know the impact of QCD corrections
- $\sigma(p\bar{p} \rightarrow t\bar{t}H)$ calculated at NLO: comparing results
- $\sigma(pp \rightarrow t\bar{t}H)$ coming soon.