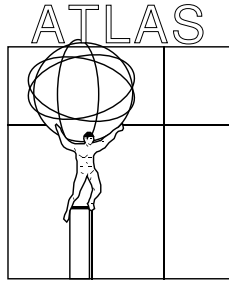


# Extraction of Higgs parameters at the LHC



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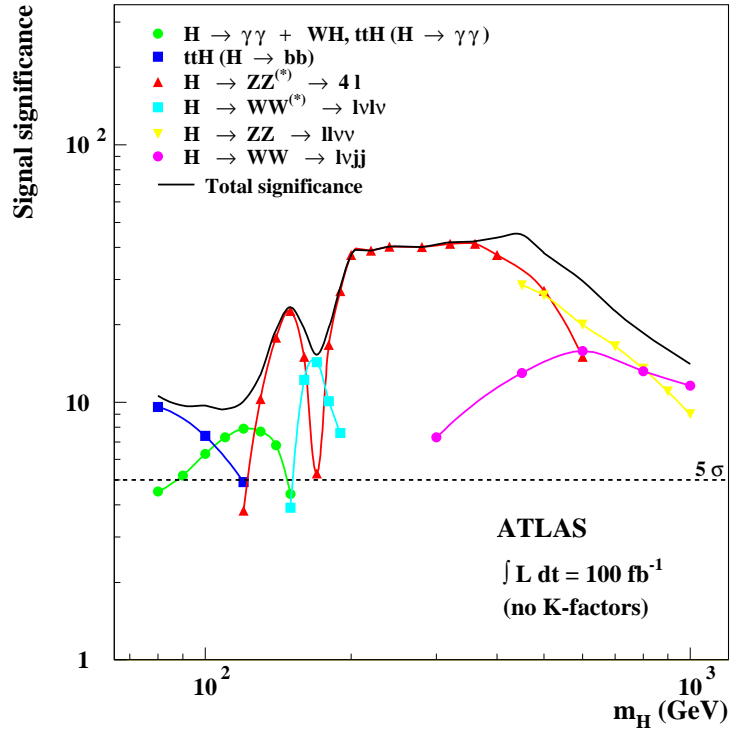
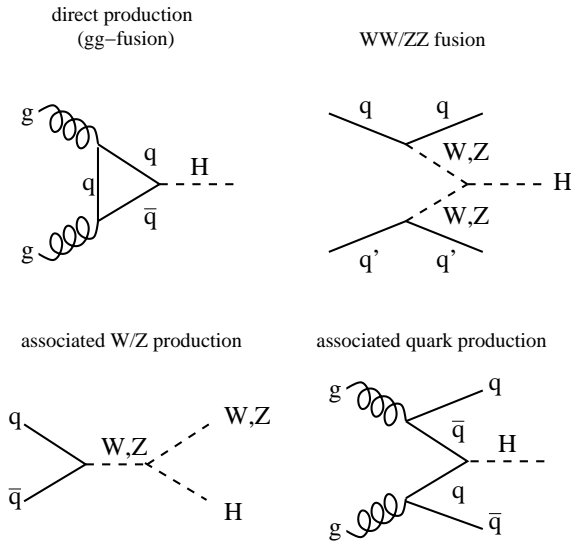
Workshop on the Future of Higgs Physics

FNAL, May 2001

- Introduction
- Measurement of the Higgs boson mass
- Couplings to bosons and fermions
- Measurement of the width
- First approach of a spin determination
- Summary

# Higgs discovery in ATLAS

## Higgs production at LHC



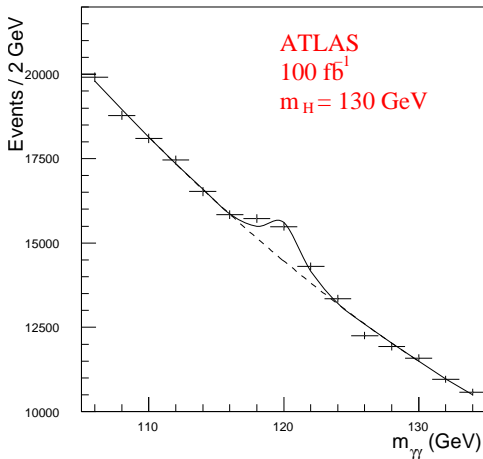
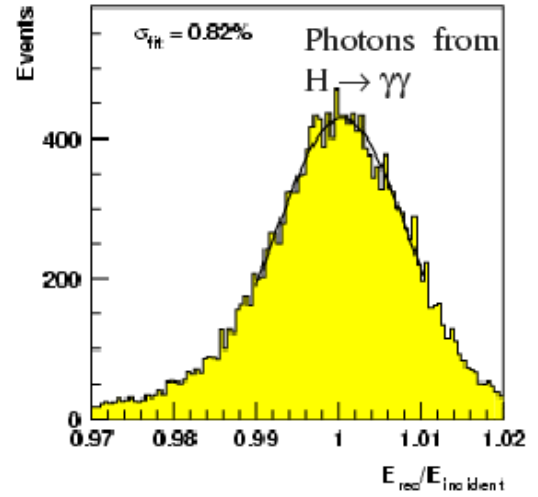
Production process	Decay Channel	Mass Range
direct prod.	$H \rightarrow \gamma\gamma$	80 GeV–150 GeV
	$H \rightarrow ZZ^* \rightarrow 4l$	120 GeV–2 $m_Z$
	$H \rightarrow ZZ \rightarrow 4l$	2 $m_Z$ –600 GeV
	$H \rightarrow WW \rightarrow l\nu l\nu$	150 GeV–190 GeV
	$H \rightarrow WW \rightarrow l\nu jj$	300 GeV–1000 GeV
	$H \rightarrow ZZ \rightarrow ll\nu\bar{\nu}$	500 GeV–1000 GeV
WW fusion	$H \rightarrow WW \rightarrow e\nu e\nu$	120 GeV–200 GeV
	$H \rightarrow \tau\tau \rightarrow ll p_T^{\text{miss}}$	110 GeV–150 GeV
ass. W prod.	$WH \rightarrow WWW \rightarrow l\nu l\nu l\nu$	150 GeV–190 GeV
	$WH \rightarrow WWW \rightarrow l\nu l\nu jj$	150 GeV–190 GeV
ass. quark prod.	$H \rightarrow b\bar{b}$	80 GeV–100 GeV
	$H \rightarrow \gamma\gamma$	80 GeV–140 GeV

# H → γγ, H → ZZ(\*) → 4l

H → γγ results for ATLAS and 100 fb<sup>-1</sup>

m <sub>H</sub> (GeV)	110	130	150
S	1110	1110	617
B	47300	33700	23350
S/√B	5.6	6.5	4.3
σ <sub>m</sub>	1.37	1.55	1.74

CMS, full simulation high L

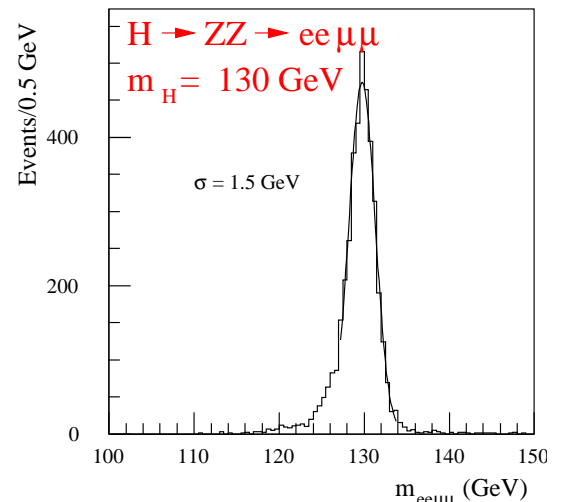


- CMS mass resolution better due to better performance of crystal calorimeter in comparison with ATLAS LAr calorimeter
- Acceptance loss bigger for CMS due to rejection of converted photons
- Assuming same K-factors and kinematic cuts, signal significances for ATLAS and CMS are within 10%

H → ZZ → 4l results for ATLAS and 100 fb<sup>-1</sup>

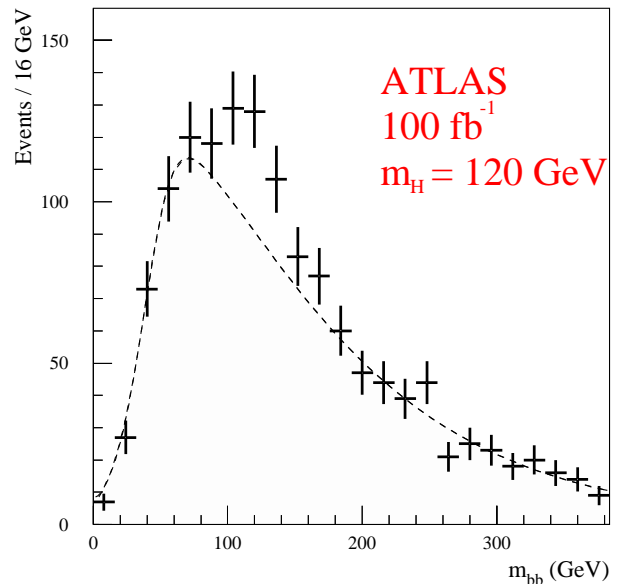
m <sub>H</sub> (GeV)	120	130	150	180
S	10.3	28.7	67.6	49.7
B	4.39	7.76	8.92	8.81
S/√B	4.9	10.3	15.5	11.2
σ <sub>m</sub> (GeV)	1.60	1.65	1.82	2.28
m <sub>H</sub> (GeV)	200	300	400	600
S	445	352	285	76
B	246	129	96	49
S/√B	28.4	31.0	29.1	10.8

1 experiment  
(full detector simulation)

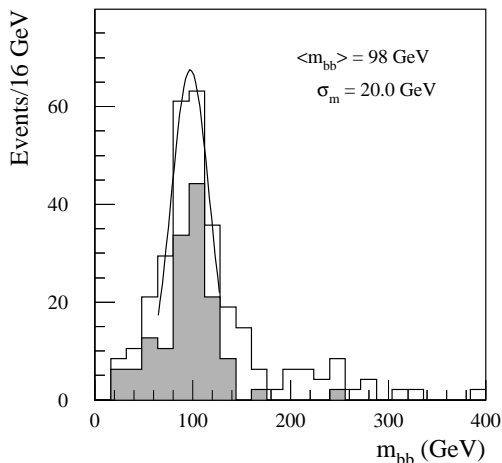


# $t\bar{t}H \rightarrow t\bar{t}b\bar{b}, t \rightarrow bj\bar{j}, t \rightarrow be/\mu\nu$

- Major backgrounds
  - combinatorial from signal
  - $Wjjjjjj, WWbbjj, \text{etc.}$
  - $ttjj$  (dominant, non-resonant)



- Background not flat under the peak
  - ⇒ **Uncertainty of 10%** assumed for background subtraction
  - ⇒ This error is included in determination of rate and coupling



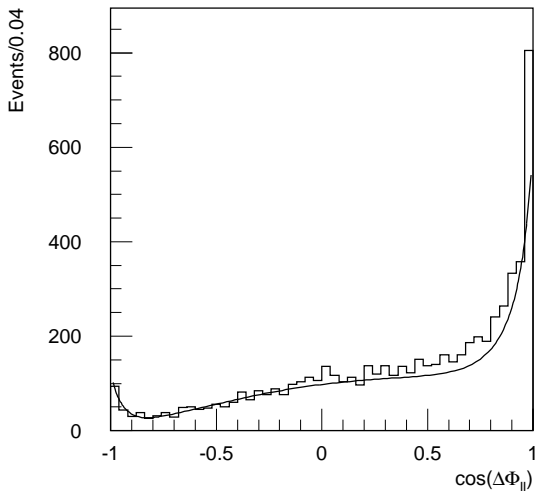
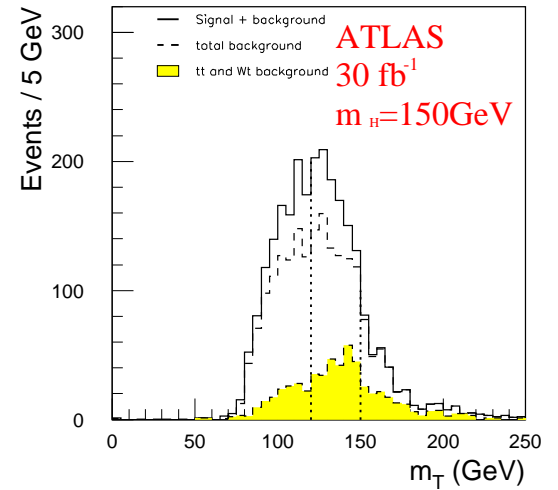
fully simulated events

$m_H$ (GeV)	100	120
S	107	62
B	278	257
$S/\sqrt{B}$	6.4	3.9

# H → WW → |ν|ν

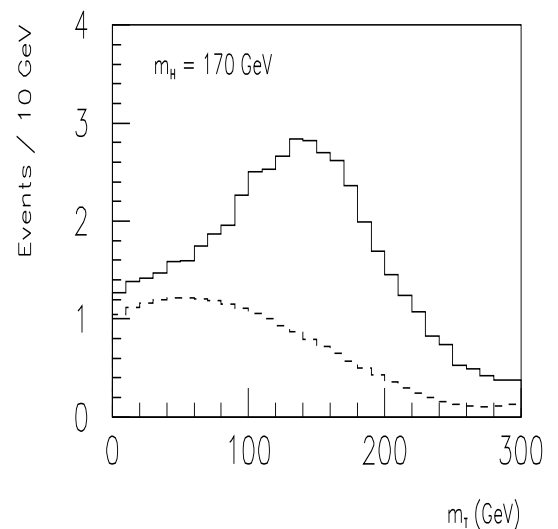
- **No mass peak** due to neutrinos in final state
- **Transverse mass:**  

$$m_H = \sqrt{2 \cdot p_T \cdot E_T^{\text{miss}} (1 - \cos \Delta\phi)}$$
- Signal appears as excess above background (WW and tt)
- Shape similar, sidebands cannot be used to determine background below signal



- Better signal to background separation in  $\cos(\Delta\Phi)$  due to spin correlations
- Normalization of background in low signal region ( $\cos(\Delta\Phi) < 0.2$ )  
 $\Rightarrow$  Allows background estimate with uncertainty of  $\pm 3.5\%$  (result from MC study)  
 $\Rightarrow$  Included in determination of rates and couplings

- Associated channel:  
WH → WW → |ν|ν|ν
- Similar uncertainty on background subtraction assumed as for WW channel
- **Important** for determination of couplings

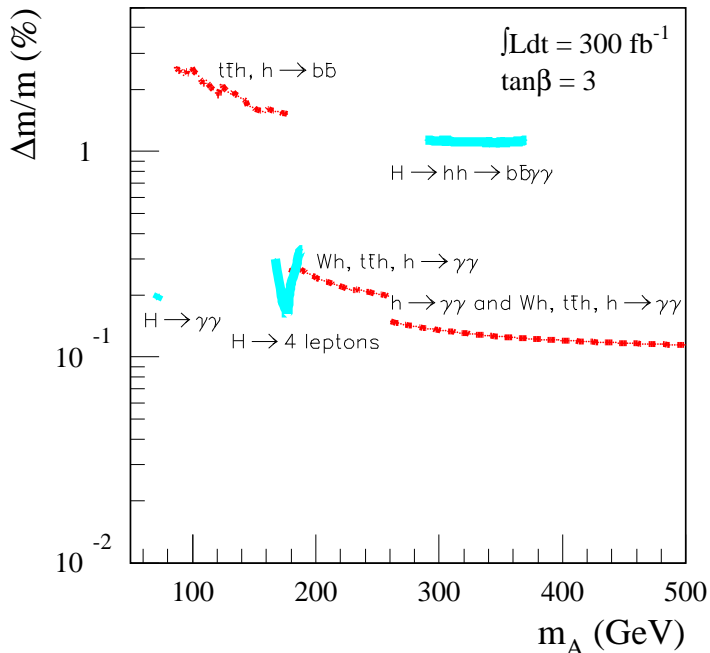
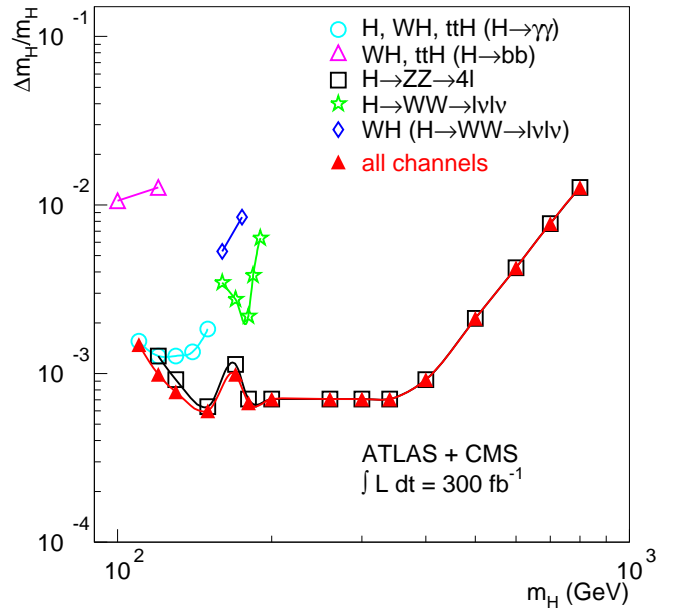


# Results for the mass reconstruction

- Precision includes statistical error due to limited number of signal events and background subtraction
- **Systematic error:**
  - Leptonic E scale (assumed to be 0.1%, goal 0.02%)
  - 1% error on energy scale for hadronic channels
- **No theoretical error** included:
  - **Interference** of signal and background **not included**  
 $\Rightarrow$  shift of  $m_H$  for large  $m_H$  and  $\Gamma_H$
  - Uncertainty from structure functions expected to be small
- Expected mass resolution combining rates of **ATLAS** and **CMS** and assuming  $300 \text{ fb}^{-1}$ /experiment:

$$\Delta m_H / m_H \approx 0.06\% (m_H < 400 \text{ GeV})$$

$$\Delta m_H / m_H \approx 1\% (m_H = 700 \text{ GeV})$$



## MSSM

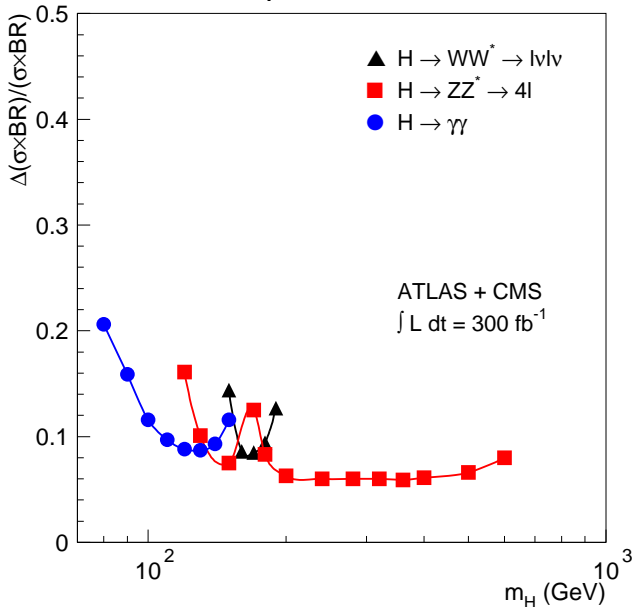
Channel	$\Delta m/m$ (%)
$h, A, H \rightarrow \gamma\gamma$	0.1–0.4
$H \rightarrow ZZ^* \rightarrow 4l$	0.1–0.4
$H/A \rightarrow \mu\mu$	0.1–1.5
$h \rightarrow b\bar{b}$	1–2
$H \rightarrow hh \rightarrow b\bar{b}\gamma\gamma$	1–2
$A \rightarrow Zh \rightarrow b\bar{b}$	1–2
$H/A \rightarrow \tau\tau$	1–10

# Measurement of the rate

- Measurement of production cross-section times decay branching ratio

$$\sigma \times \text{BR} = \frac{N_{S+B} - \langle N_B \rangle}{\epsilon \cdot \eta \cdot \int \mathcal{L} dt}$$

direct production



- Included errors:

Luminosity:  $\Delta L/L = 5\%$

Efficiency and acceptance:  $\Delta(\epsilon \cdot \eta)/(\epsilon \cdot \eta) = 2\%$

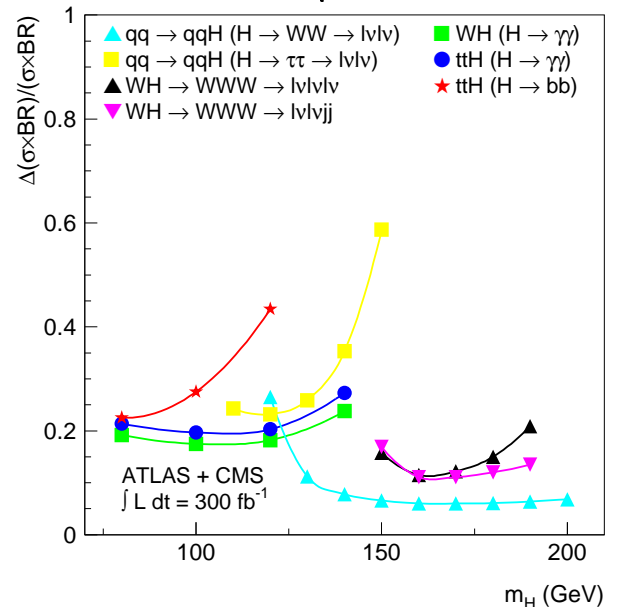
Uncertainty of background:

- 5% for  $H \rightarrow WW$
- 10% for  $H \rightarrow b\bar{b}$

Typical precision for ATLAS+CMS and  $300 \text{ fb}^{-1}$ /experiment:

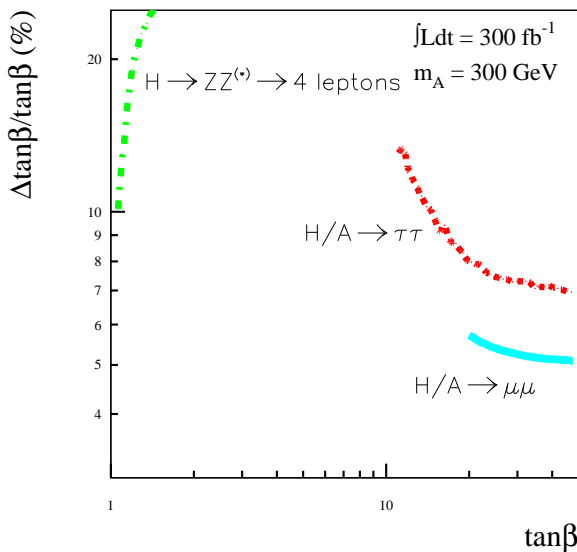
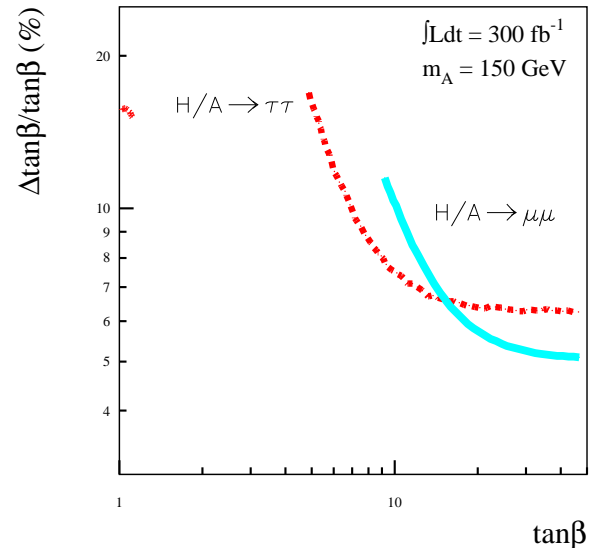
- Direct production:** 6%–20% in mass range  $120 < m_H < 600 \text{ GeV}$
- Associated production:** 6%–30% in mass range  $100 < m_H < 200 \text{ GeV}$

associated production



# Measurement of $\tan\beta$

- Measurement of the signal rate of heavy Higgs H provides good **sensitivity to  $\tan\beta$**
- Accuracy for  $\tan\beta$  assuming  $300 \text{ fb}^{-1}$  and  $m_A=150\text{GeV}$
- $H/A \rightarrow \tau\tau$  channel:  
 $\Delta(\tan\beta)/\tan\beta = 6\text{--}15\%$
- $H/A \rightarrow \mu\mu$  channel:  
 $\Delta(\tan\beta)/\tan\beta = 5\text{--}12\%$
- $H \rightarrow ZZ^{(*)} \rightarrow 4l$  channel:  
 $\Delta(\tan\beta)/\tan\beta = 10\text{--}25\%$



- Systematic error is dominated by luminosity (assumed here to be 10%)
- $H \rightarrow ZZ^* \rightarrow 4l$  channel is rate limited

- Compared to the SM the rate in the  $H \rightarrow ZZ^{(*)} \rightarrow 4l$  channel is **suppressed by an order of magnitude** in the MSSM  $\Rightarrow$  Distinction between SM and MSSM may be feasible



# Coupling to bosons

- Without theoretical input only measurement of **coupling ratios** possible

- Direct measurement**

$$- \frac{\sigma \times \text{BR}(H \rightarrow WW^*)}{\sigma \times \text{BR}(H \rightarrow ZZ^*)} = \frac{\Gamma_g \Gamma_W}{\Gamma_g \Gamma_Z} = \frac{\Gamma_W}{\Gamma_Z}$$

- QCD corrections cancel (in first order)

- Indirect measurement**

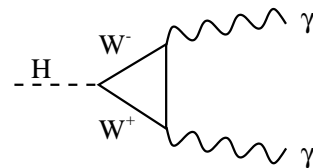
$$- \frac{\sigma \times \text{BR}(H \rightarrow \gamma\gamma)}{\sigma \times \text{BR}(H \rightarrow ZZ^*)} = \frac{\Gamma_g \Gamma_\gamma}{\Gamma_g \Gamma_Z} \sim \frac{\Gamma_W}{\Gamma_Z}$$

- Indirect measurements  $\Rightarrow$  Proportionalities are used

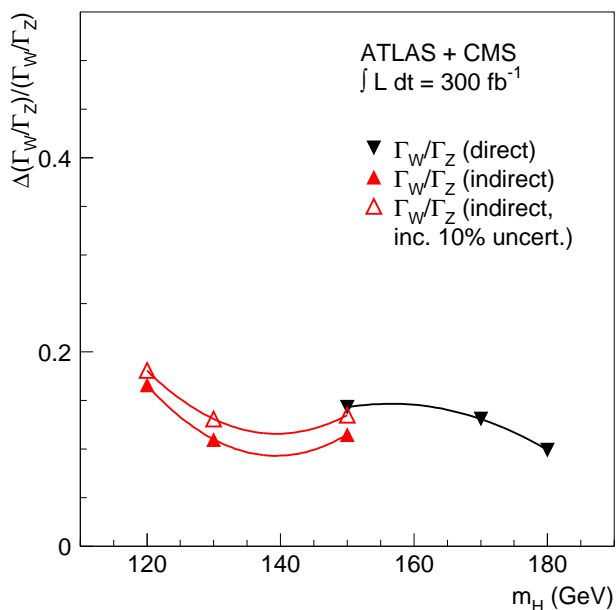
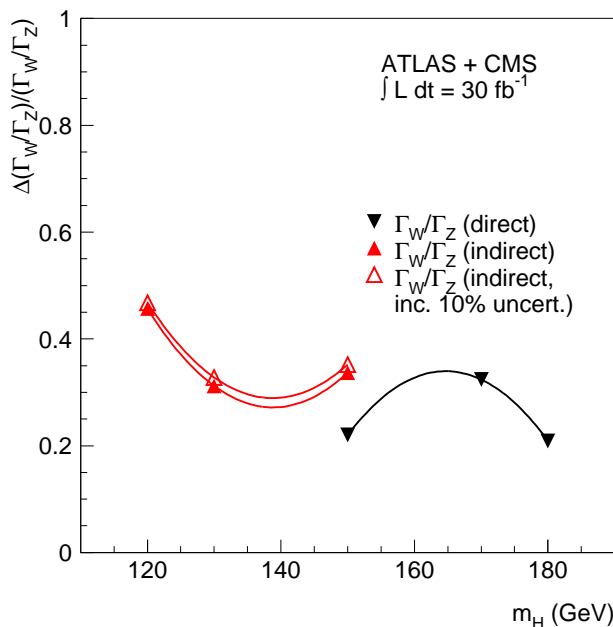
$$\Gamma_W = \alpha \Gamma_\gamma$$

$$\Gamma_t = \beta \Gamma_g$$

Error of **10%** assumed for  $\alpha$  resp.  $\beta$



Results for  $30 \text{ fb}^{-1}$  and  $300 \text{ fb}^{-1}$  per experiment



- No theoretical errors included:

**Different K-factors** for gg-Fusion ( $K_g=1.7-1.8$ ) and WW-Fusion ( $K_W=1.06-1.08$ )

$\Rightarrow$  Additional uncertainty, if channels with different production processes are compared

# Ratio of boson/fermion coupling

- Direct measurement

$$- \frac{\sigma \times \text{BR}(qq \rightarrow qqH(H \rightarrow WW))}{\sigma \times \text{BR}(qq \rightarrow qqH(H \rightarrow \tau\tau))} = \frac{\Gamma_W \Gamma_W}{\Gamma_W \Gamma_\tau} = \frac{\Gamma_W}{\Gamma_\tau}$$

- Indirect measurement

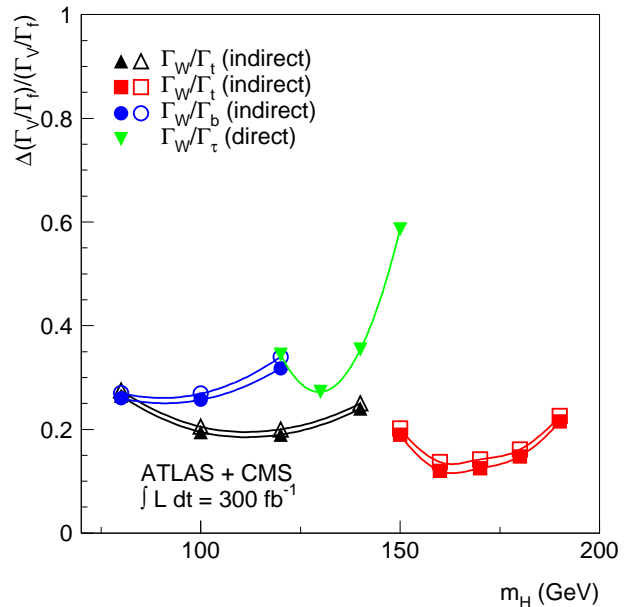
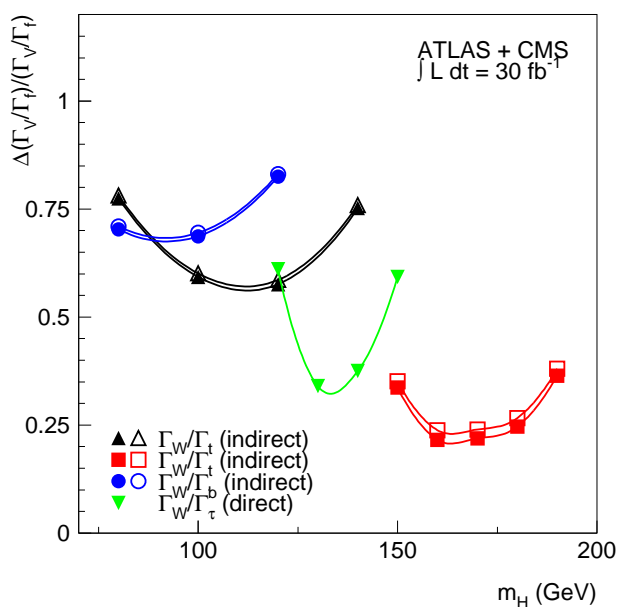
$$- \frac{\sigma \times \text{BR}(WH(H \rightarrow \gamma\gamma))}{\sigma \times \text{BR}(H \rightarrow \gamma\gamma)} = \frac{\Gamma_W \Gamma_\gamma}{\Gamma_g \Gamma_\gamma} \sim \frac{\Gamma_W}{\Gamma_t} * C_{QCD}$$

$$- \frac{\sigma \times \text{BR}(WH(H \rightarrow WW))}{\sigma \times \text{BR}(H \rightarrow WW^*)} = \frac{\Gamma_W \Gamma_W}{\Gamma_g \Gamma_W} \sim \frac{\Gamma_W}{\Gamma_t} * C_{QCD}$$

$$- \frac{\sigma \times \text{BR}(ttH(H \rightarrow bb))}{\sigma \times \text{BR}(ttH(H \rightarrow \gamma\gamma))} = \frac{\Gamma_t \Gamma_b}{\Gamma_t \Gamma_\gamma} \sim \frac{\Gamma_b}{\Gamma_W}$$

\* Uncertainties on the ratio arising through different production processes are not included

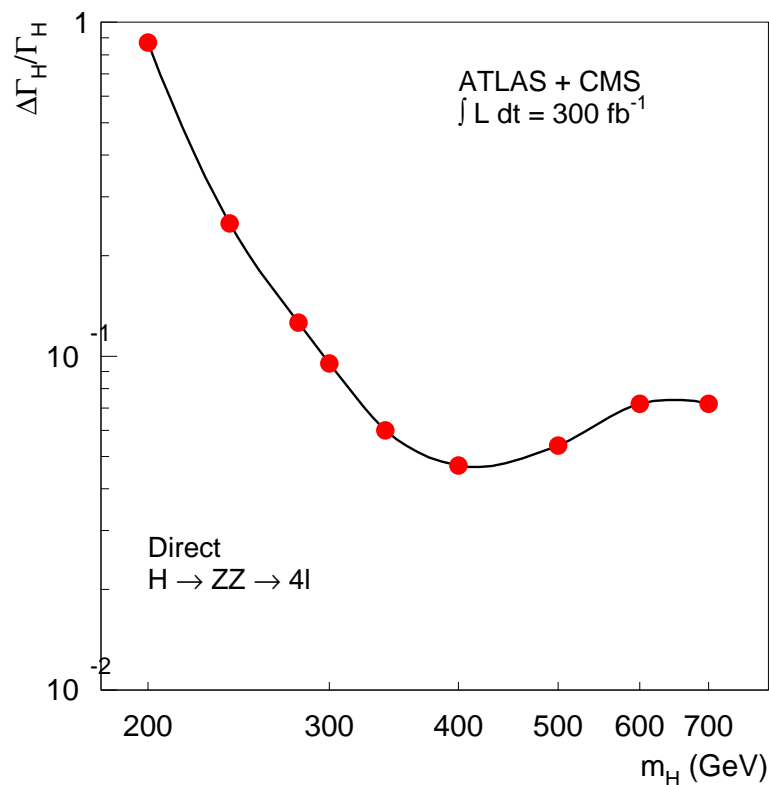
Results for  $30 \text{ fb}^{-1}$  and  $300 \text{ fb}^{-1}$  per experiment



# Measurement of the total width

- Determined from reconstructed width of reconstructed mass peak
- Direct measurement only possible for  $m_H > 200$  GeV
  - $\Gamma_H^{\text{tot}} < 0.1$  GeV for  $m_H < 2M_Z$
  - expected detector resolution 1.2–1.4 GeV
- mostly covered by  $H \rightarrow ZZ \rightarrow 4l$  decays

Result for ATLAS and CMS and  $300 \text{ fb}^{-1}$  per experiment

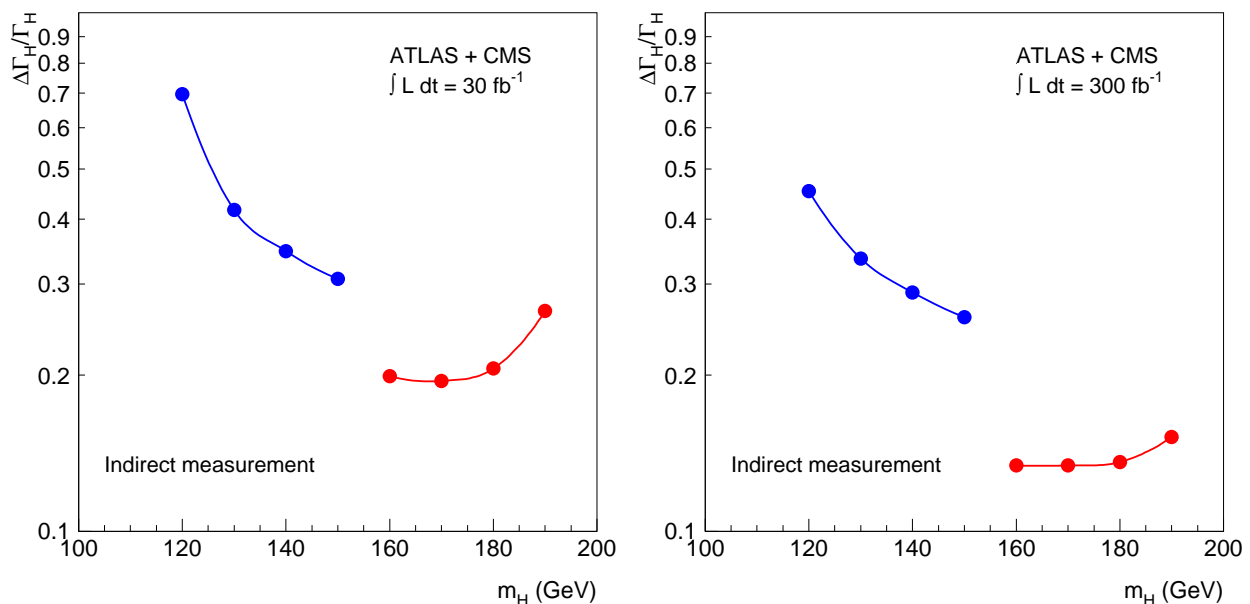


- systematic error dominated by uncertainty of radiative decays, assumed to be 1.5%

# Indirect measurement for $m_H < 200$ GeV

- Proposed by D. Zeppenfeld et al., Phys. Rev. **D 62** (2000) 013009/1–10
  - Computation of partial width  $\tilde{\Gamma}_W$   
 $[\Gamma_b + \Gamma_\tau + \Gamma_W + \Gamma_Z + \Gamma_g] \frac{\Gamma_W}{\tilde{\Gamma}} = \Gamma_W(1 - \epsilon) = \tilde{\Gamma}_W$
  - Using SU(2)-symmetry  
 $\Gamma_Z = z \cdot \Gamma_W \quad \Gamma_b = y \cdot \Gamma_\tau$
  - Computation of total width  $\tilde{\Gamma}$   
 $\tilde{\Gamma} = \Gamma(1 - \epsilon)^2 = \frac{\Gamma_W^2}{X_W}$
- Important channels: **Weak Boson Fusion**  
 $qq \rightarrow qqH$  ( $H \rightarrow WW \rightarrow \nu\nu$ ):  $120 \text{ GeV} \leq m_H \leq 190 \text{ GeV}$   
 (Detector simulation)
   
 $qq \rightarrow qqH$  ( $H \rightarrow \tau\tau \rightarrow \ell p_T^{\text{miss}}$ ):  $110 \text{ GeV} \leq m_H \leq 150 \text{ GeV}$   
 (Detector simulation, statistics corresponding to  $30\text{fb}^{-1}$ /experiment)
- $\epsilon < 5\%$  over whole mass range

Results for  $30\text{fb}^{-1}$  and  $300\text{fb}^{-1}$  per experiment



# Decay correlation function for $H \rightarrow ZZ^{(*)} \rightarrow 4l$

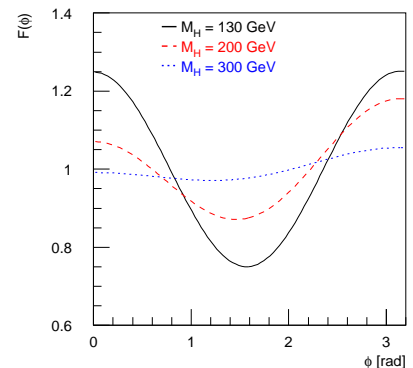
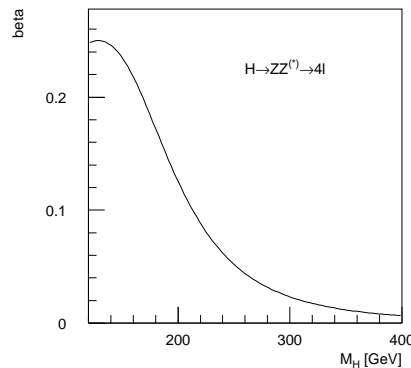
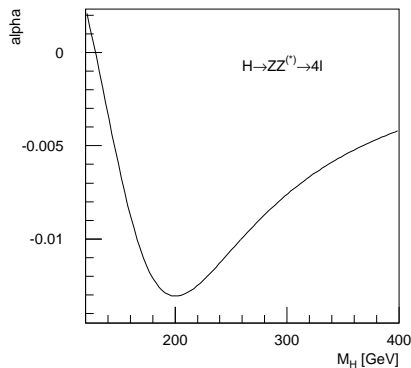
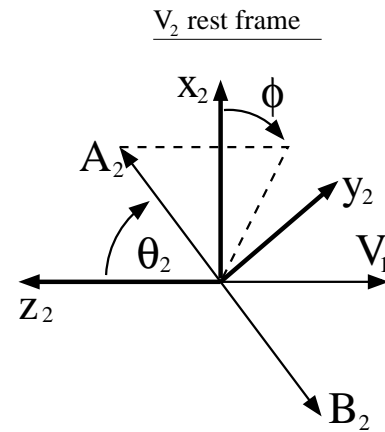
- From C.A. Nelson, Phys. Rev. **D 37** (1988) 1220

$$F(\Phi) = 1 + \alpha \cos\Phi + \beta \cos 2\Phi$$

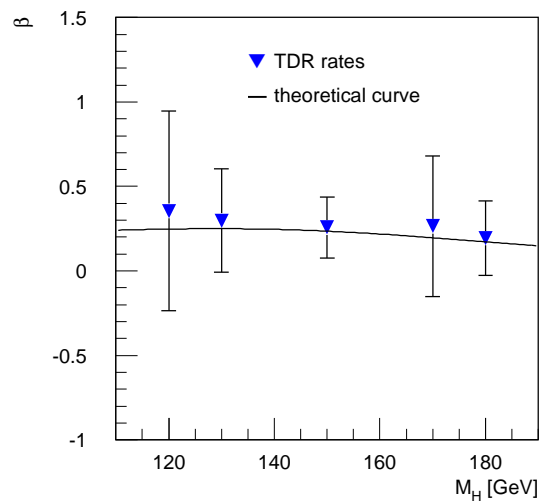
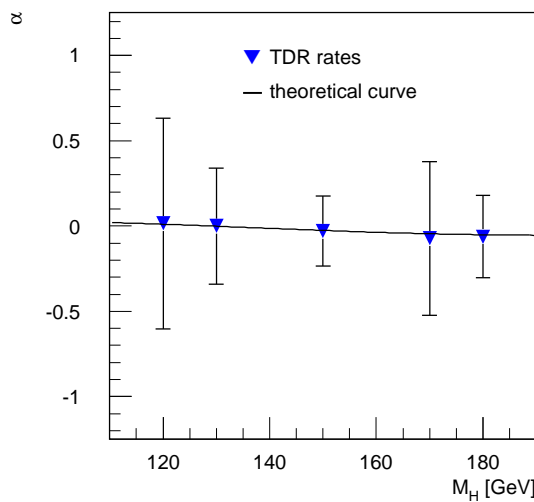
$$\alpha = \left(\frac{3\pi}{8}\right)^2 \frac{\gamma}{(1 + \gamma^2/2)}$$

$$\beta = \frac{1}{4(1 + \gamma^2/2)}$$

$$\gamma = \frac{(M_H^2 - 2m_V^2)}{2m_V^2}$$



Sensitivity for  $\alpha$  and  $\beta$  for  $100 \text{ fb}^{-1}$  and  $m_H < 2m_Z$



- Similar results for mass range  $200 \text{ GeV} < m_H < 400 \text{ GeV}$

# Summary

- **Mass:**
  - Mass resolution around **0.06 %** for  $m_H < 400$  GeV
  - **Interference** of signal and background must be included in further studies
- **Rates:**
  - In various channels the rate can be determined with a relative precision of **6%**
  - Good rate measurement desirable for determination of couplings and width
- **Couplings:**
  - The ratio of the coupling to W and Z bosons can be measured with a relative precision of **10–20%**
  - The ratio of boson and fermion coupling can be determined with an accuracy of **10–35%**
- **Width:**
  - In the mass range above 200 GeV a direct measurement of the width is possible with a relative precision of **5–20%** for  $m_H > 220$  GeV
  - In the intermediate mass range the width can be determined **indirectly** with an accuracy of **13–45%**
- **Spin:**
  - **Consistency with spin 0** over the whole mass range 120 to 400 GeV can be demonstrated
  - Comparison with particle with spin  **$S \neq 0$**  necessary