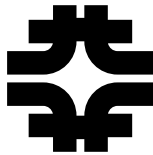


Backgrounds to Tevatron Higgs production

Keith Ellis

Fermilab



In collaboration with:

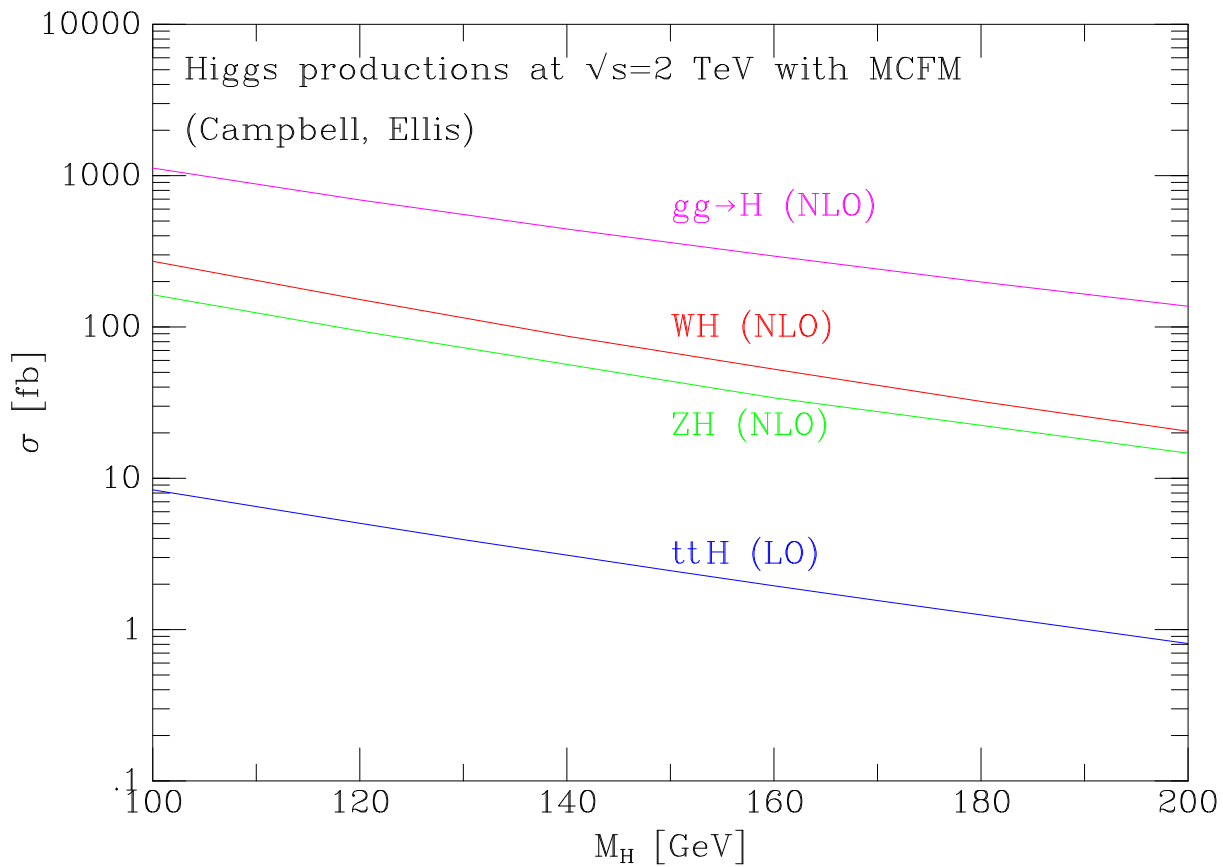
John Campbell and Sinisa Veseli

hep-ph/9810489

hep-ph/9905386

hep-ph/0006304

Higgs production at 2 TeV



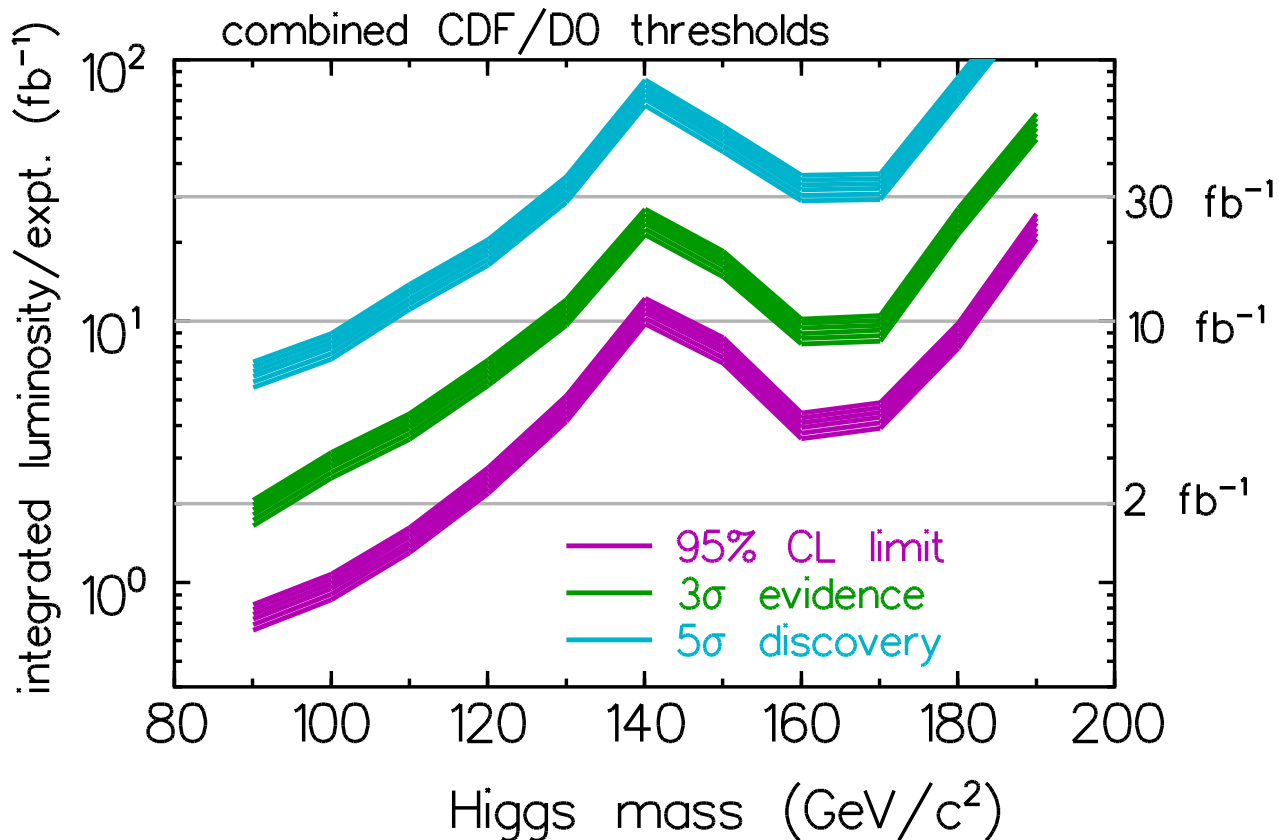
- Higgs cross sections are not so small at the Tevatron.
- Issue of background is crucial.
- This talk will consider SM physics backgrounds.

Scope of the talk

- Aim is to provide the most accurate NLO information on signals and backgrounds, so that relative rates and shapes of cross sections are predicted.
- NLO calculations only contain partial effects of parton showering. For the heavy objects considered here this may be a good approximation.
- No detector effects on backgrounds. So a full assessment of the Higgs reach will not be possible.
- These are preparatory studies which will inform the more complete analysis to be performed when the run II data becomes available.



Higgs working group results



- Does this contain the most accurate theoretical information?
- $t\bar{t}H$ not included (4.1σ effect for $M_H = 120\text{GeV}$ and $15\text{fb}^{-1}/\text{expt}$).



Synopsis of Talk

- Description of the NLO program MCFM.
- Backgrounds to WH production (low mass Higgs).
- Backgrounds to ZH production.(low mass Higgs)
- Backgrounds to l^+l^- +missing energy signature (high mass Higgs)
- No MCFM studies yet of $t\bar{t}H$ production, (largest background $t\bar{t} + 2$ jets)



MCFM Philosophy

- The Tevatron Run II will be sensitive to processes at the femtobarn level.
- Particularly interesting are final states involving heavy quarks, leptons and missing energy.
- Standard model processes which result in these final states often include W , Z or Higgs. Data samples may not be large enough to determine SM backgrounds experimentally.
- **MCFM** aims to provide a unified description of such processes at **NLO** accuracy.
- The extension to **NLO** is made possible in many cases by the recent calculations of virtual matrix elements involving a vector boson and four partons.



MCFM Process List

Included at NLO

$$p\bar{p} \rightarrow W^\pm/Z$$

$$p\bar{p} \rightarrow W^\pm + Z$$

$$p\bar{p} \rightarrow W^\pm/Z + H$$

$$p\bar{p} \rightarrow W^\pm + g^* (\rightarrow b\bar{b})$$

$$p\bar{p} \rightarrow W^+ + W^-$$

$$p\bar{p} \rightarrow Z + Z$$

$$p\bar{p} \rightarrow W^\pm/Z + 1 \text{ jet}$$

$$p\bar{p} \rightarrow Z + b\bar{b}$$

- Various leptonic and/or hadronic decays of the bosons are included as further sub-processes.
- Where applicable Z includes the Z/γ^* interference, (important e.g. for trilepton signatures).
- MCFM v1.0 may be downloaded from <http://www-theory.fnal.gov/people/campbell/mcfm.html>



Processes included in MCFM at LO (including vector boson and top decays).

- $p\bar{p} \rightarrow t\bar{t}$
- $p\bar{p} \rightarrow t\bar{b}$ (single top production)
- $p\bar{p} \rightarrow qt\bar{b}$ (single top production, W gluon fusion)
- $p\bar{p} \rightarrow t\bar{t}H$



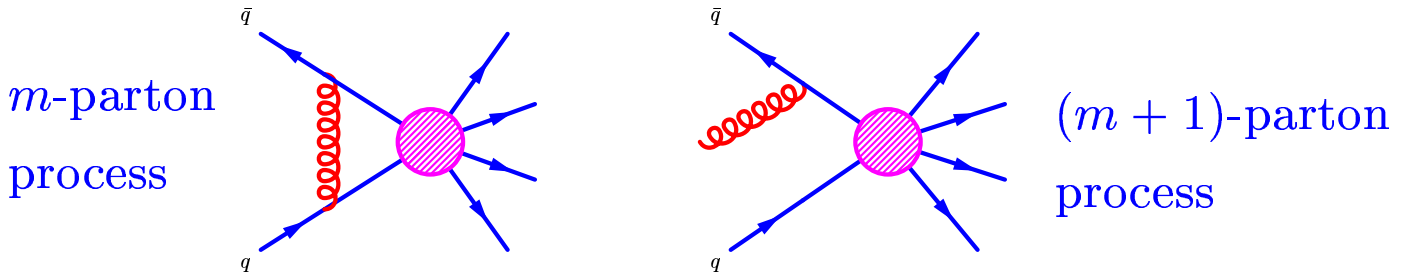
Processes which are under construction at NLO

- $p\bar{p} \rightarrow W^\pm + 2 \text{ jet}$
- $p\bar{p} \rightarrow Z/\gamma^* + 2 \text{ jet}$



Monte Carlo Ingredients - 1

- Helicity amplitudes for the virtual and real ME's

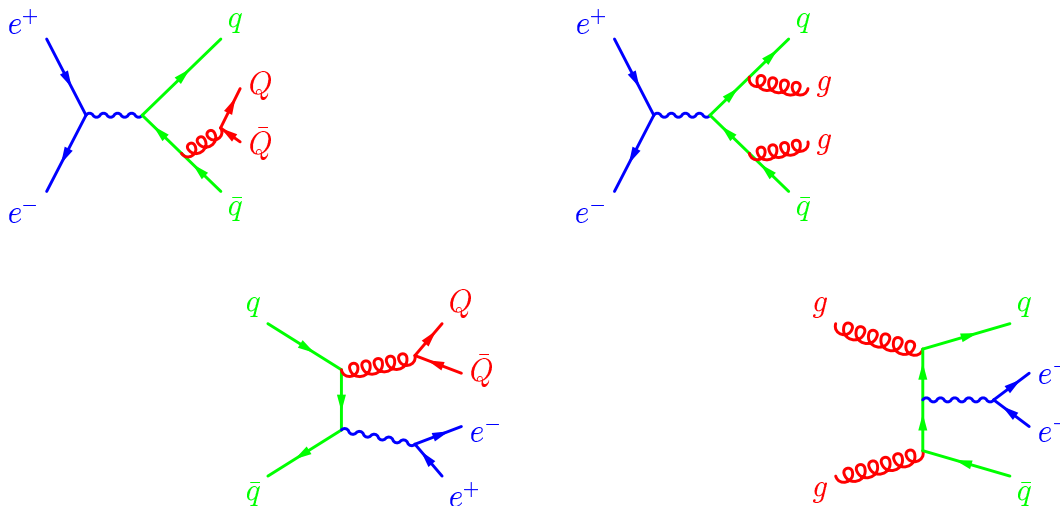


- Many of the NLO matrix elements are obtained by crossing the ones calculated for $e^+e^- \rightarrow 4$ jets.

Bern, Dixon, Kosower and Weinzierl, Nucl. Phys. **B489** (1997) 3

Glover and Miller, Phys. Lett. **B396** (1997) 257

Campbell, Glover and Miller, Phys. Lett. **B409** (1997) 503



Monte Carlo Ingredients - 2

- Singular pieces of the real matrix elements must be identified and cancelled by an appropriate set of counter-terms. [ERT](#), Nucl.Phys.**B178** (1981) 421
- MCFM uses the **dipole** method to cancel the infrared divergences between real and virtual contributions.

[Catani and Seymour](#), Nucl. Phys. **B485** (1997) 291

$$\begin{aligned}\sigma_{real}^{m+1} &= \int_{(m+1)} (d\sigma_{real} - d\sigma_{counter}) + \int_{(m+1)} d\sigma_{counter} \\ &= (\text{integrable terms}) + \sum_{dipoles} \int_m d\sigma \otimes \int_1 dV_{dipole}\end{aligned}$$

where the 1-dimensional integral over the dipoles leads to soft and collinear divergences (poles in ϵ).

- These poles manifestly multiply m -parton ME's and may be cancelled against poles from the loop diagrams.



Higgs search using MCFM

- Studies using LO Monte Carlos and other event generators show that for a Higgs in the mass range of 100-130 GeV, the most promising channels for discovery at Run II are **associated Higgs production**.

Stange, Marciano, Willenbrock, Phys. Rev. **D49** (1994) 1354, **D50** (1994) 4491

$$p\bar{p} \longrightarrow W(\rightarrow e\nu)H(\rightarrow b\bar{b})$$

$$p\bar{p} \longrightarrow Z(\rightarrow \nu\bar{\nu}, \ell\bar{\ell})H(\rightarrow b\bar{b})$$

- Mass range interesting in the light of hints from SUSY.
- Backgrounds for the WH signal:

$$p\bar{p} \longrightarrow W g^*(\rightarrow b\bar{b})$$

$$p\bar{p} \longrightarrow W Z/\gamma^*(\rightarrow b\bar{b})$$

$$p\bar{p} \longrightarrow t(\rightarrow bW^+)\bar{t}(\rightarrow \bar{b}W^-)$$

$$p\bar{p} \longrightarrow W^{\pm*}(t(\rightarrow bW^+)\bar{b})$$

$$qg \longrightarrow q't(\rightarrow bW^+)\bar{b}$$



Cuts for a WH study

- Use a set of “standard” cuts from the literature

$$|p_b^T|, |p_{\bar{b}}^T| > 15 \text{ GeV}$$

$$|y_b|, |y_{\bar{b}}| < 2 ,$$

$$|p_e^T|, |p_{\nu}^T| > 20 \text{ GeV}$$

$$|y_e| < 2.5 ,$$

$$|p_{jet}^T| > 20 \text{ GeV}$$

$$|y_{jet}| < 2 ,$$

$$R_{b\bar{b}}, R_{eb}, R_{e\bar{b}} > 0.7$$

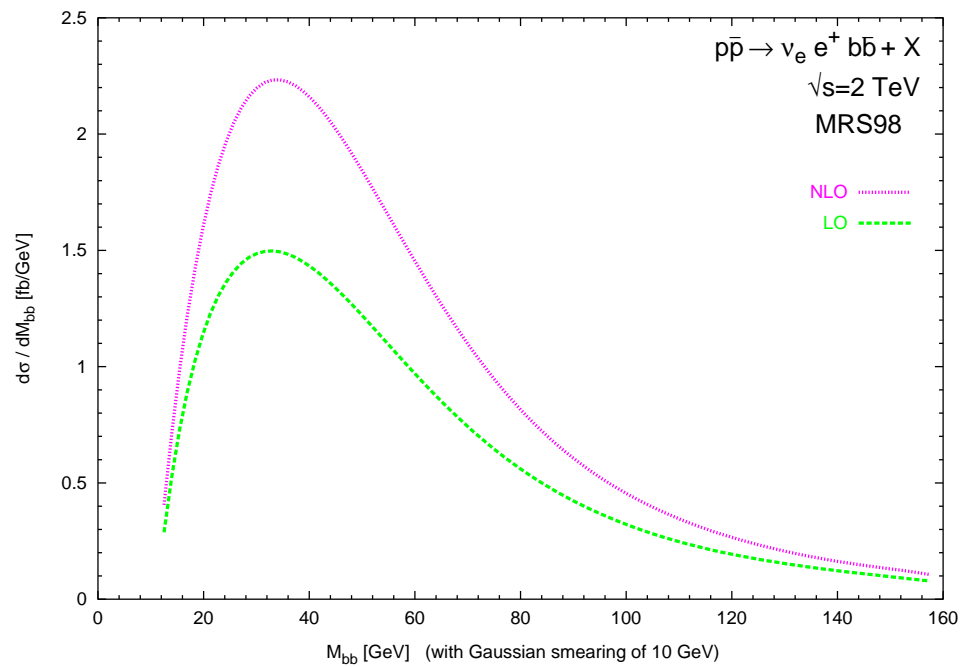
$$|\cos \theta_{b\bar{b}}| < 0.8$$

- $\theta_{b\bar{b}}$ is the scattering angle of the $b\bar{b}$ system in the Collins-Soper frame
- MRS98 parton distribution functions



$m_{b\bar{b}}$ mass distribution

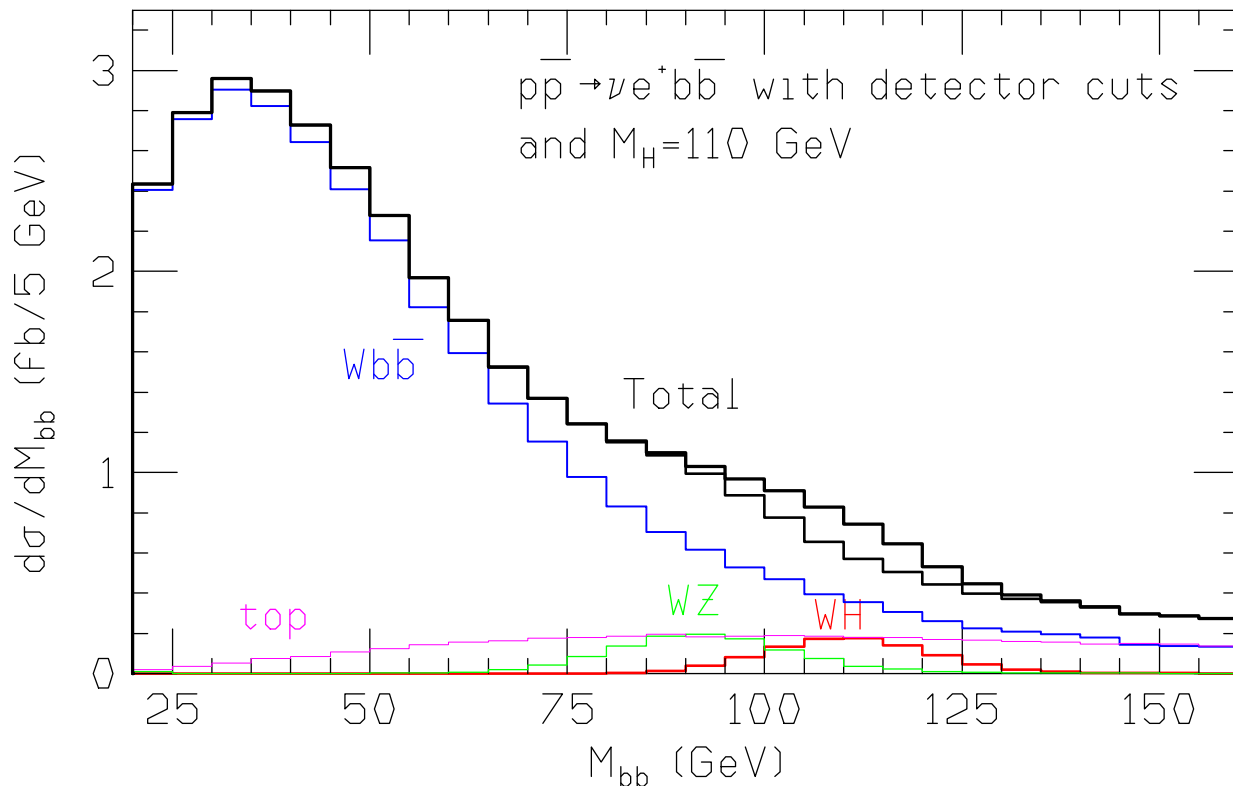
- Examine the $m_{b\bar{b}}$ mass distribution of the largest background at LO and NLO and a scale of 100 GeV



- We see that the actual shape changes very little and the K -factor is around 1.5



Signal and Backgrounds for $m_H = 110 \text{ GeV}$



- Double b -tagging efficiency of $\epsilon_{b\bar{b}} = 0.45$
- We have to understand a cocktail of backgrounds, both in normalization and shape. ($WZ \equiv WH$).



Signal and backgrounds at 2 TeV

- We take a resolution of $\Delta = 0.1 \times m_H$ and integrate signal and background over the range

$$m_H - \sqrt{2}\Delta < m_{b\bar{b}} < m_H + \sqrt{2}\Delta$$

M_H [GeV]	100	110	120	130
$W^\pm H(\rightarrow b\bar{b})$	8.8	6.4	4.2	2.5
$W^\pm g^*(\rightarrow b\bar{b})$	22.9	20.3	16.8	14.2
$W^\pm Z(\rightarrow b\bar{b})$	6.7	4.3	2.0	1.0
$t(\rightarrow bW^+)\bar{t}(\rightarrow \bar{b}W^-)$	4.2	4.7	4.9	5.0
$W^{\pm*}(t(\rightarrow bW^+)\bar{b})$	5.3	6.0	6.2	6.2
$q't(\rightarrow bW^+)$	3.1	3.2	2.9	2.6
Total B	42.1	38.5	32.9	28.8
S/B	0.21	0.17	0.13	0.09
S/\sqrt{B}	1.36	1.03	0.73	0.47



Conclusions on WH channel

- Tau-decay in the $t\bar{t}$ channel is not properly implemented, and may result in an under-estimate of this background
- Our analysis confirms the rates used in the SUSY Higgs study for the WH channel.



ZH Study

- Similar cuts to WH . Cut also on the missing transverse energy and the azimuthal angles between the direction of the missing energy and the b jets:

$$E_T^{miss} > 35 \text{ GeV}$$
$$\Delta\phi(E_T^{miss}, jets) > 0.5$$

Caveat

- A further background to the ZH study could be from ordinary ($p\bar{p} \rightarrow b\bar{b}$) events which also contain sufficient missing energy - for example, from a jet that is outside the acceptance region or hits a crack in the detector
- Since the total cross-section for this process is ~ 9 orders of magnitude above the femtobarn signal, even the very



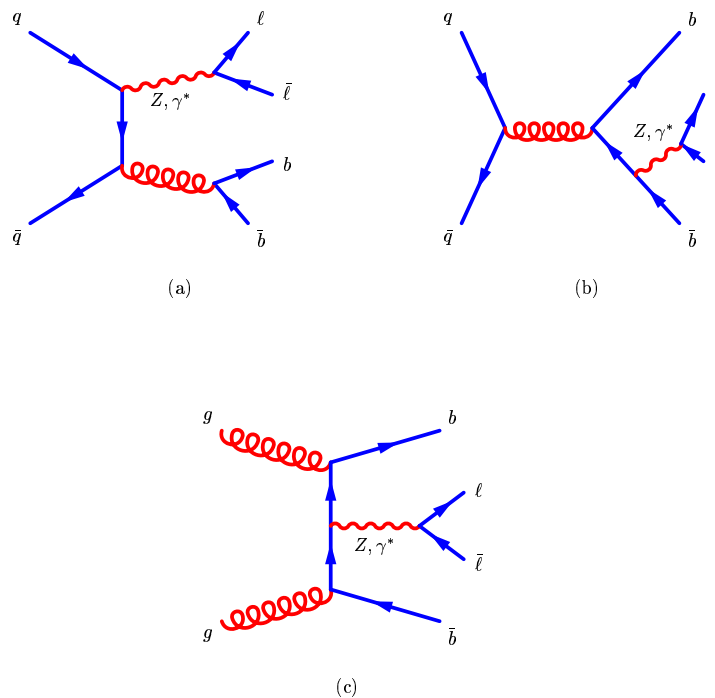
small number of events that might fall into this category could be significant

- It is unlikely that even a full shower Monte Carlo would provide a reliable estimate of such a background, much less our NLO Monte Carlo
- Ultimately, we shall have to rely on data to assess the importance of such effects



Results for $Zb\bar{b}$

- New results include radiative corrections, relevant for a further Higgs search in the channel ZH .
- The required matrix elements are very similar to the $Wb\bar{b}$ case,

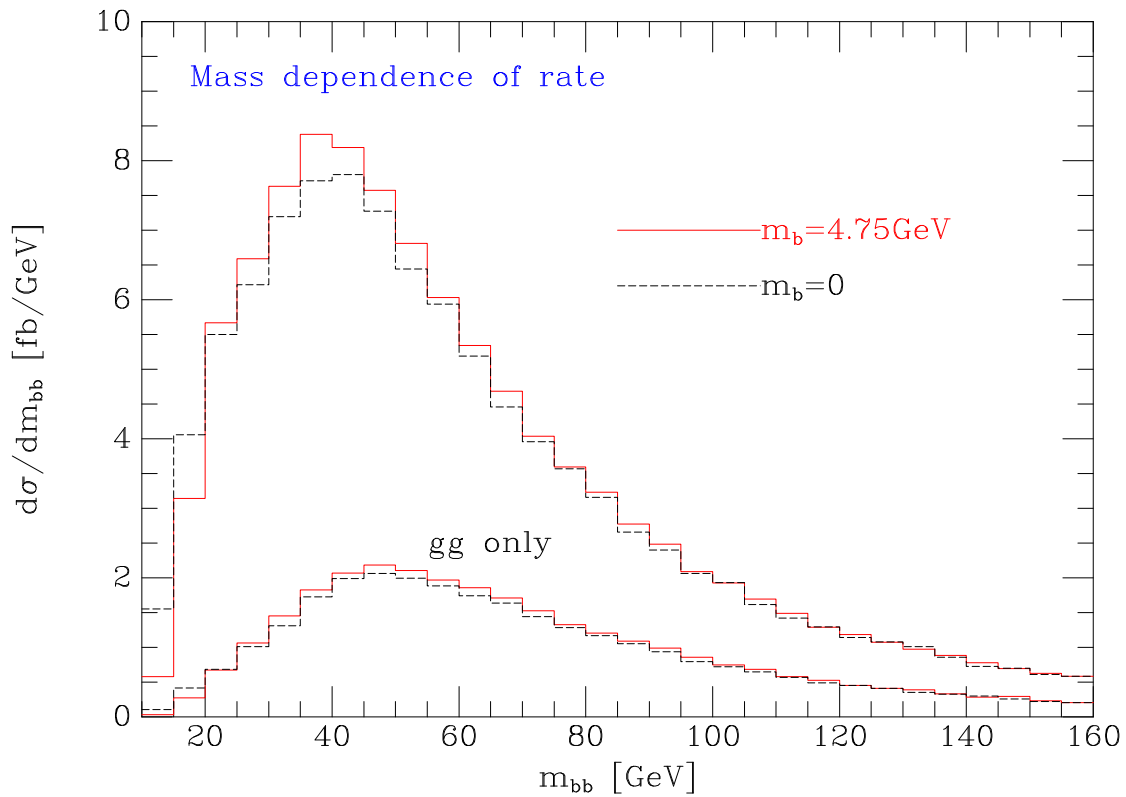


with additional contributions from gg initial states.



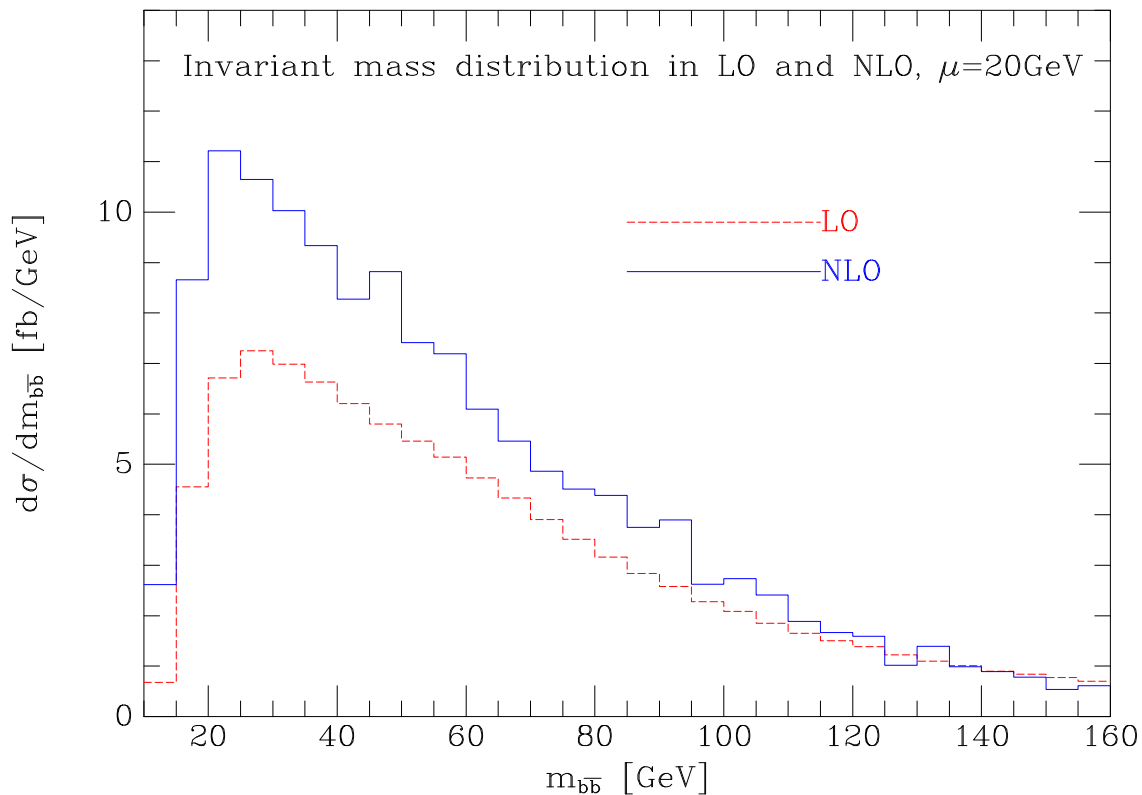
The gg sub-process in LO

- A $b\bar{b}$ pair with a large invariant mass can be produced by the gg initial state process, without off-shell propagators.
- Test of $m_b = 0$ approximation



$m_{b\bar{b}}$ mass distribution for $Zb\bar{b}$

- For a 'conventional' scale of 100 GeV, there is a large K -factor in the region of interest, around 1.8.



- The entire distribution is changed both in shape and normalization - perhaps suggesting that this scale choice is no longer appropriate (\rightarrow new gg processes).

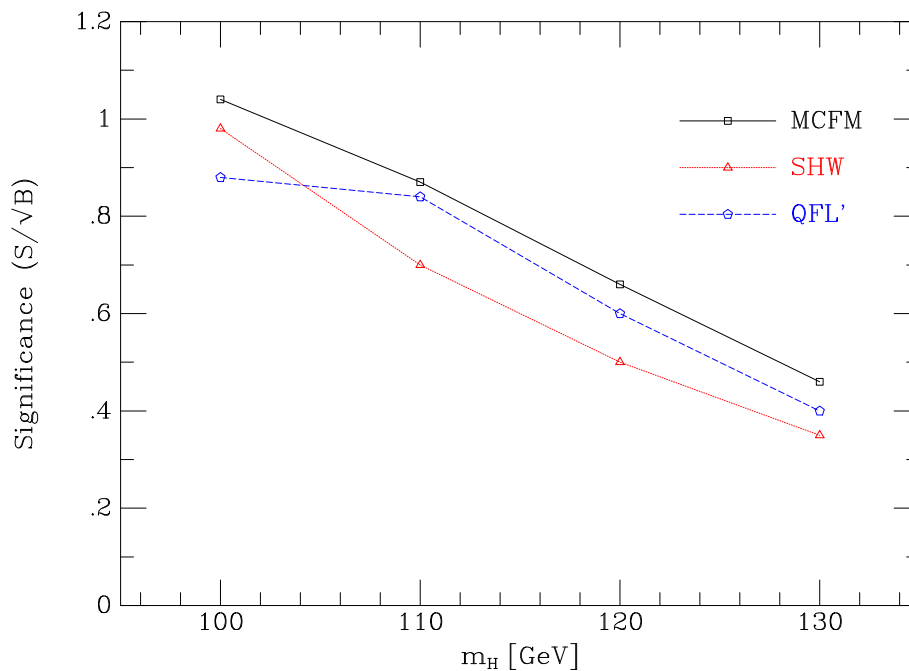


Z Channel: Signal and backgrounds

2 TeV				
M_H [GeV]	100	110	120	130
$Z^0(\rightarrow \nu\bar{\nu})H(\rightarrow b\bar{b})$	5.8	4.4	2.9	1.8
$Z^0(\rightarrow \nu\bar{\nu})g^*(\rightarrow b\bar{b})$	12.6	11.6	10.3	8.6
$Z^0(\rightarrow \nu\bar{\nu})Z(\rightarrow b\bar{b})$	9.6	6.1	3.0	1.3
$W^{\pm*}(t(\rightarrow bW^+)\bar{b})$	0.6	0.7	0.7	0.7
$q't(\rightarrow bW^+)$	0.2	0.1	0.1	0.1
Total B	23.0	18.5	14.1	10.7
S/B	0.25	0.24	0.21	0.17
S/\sqrt{B}	1.21	1.02	0.77	0.55



Significance compared with SUSY Higgs study



- MCFM overestimates significance because of lack of detector effects (misidentification of leptons).
- Susy Higgs study overestimates significance because of underestimate of $Zb\bar{b}$ background. Our **guess** is that this could be a 20% effect.



Higher mass Higgs

Basic idea (Han, Turcot, Zhang): exploit the $H \rightarrow W^*W$ decay in the higher mass region.

- **Signal processes**

$$gg \rightarrow H \rightarrow WW^*$$

$$ZH \rightarrow ZWW^*$$

$$WH \rightarrow WWW^*$$

- **Experimental signatures**

$$l^+l^- + \text{missing energy } (\nu\bar{\nu})$$

$$l^\pm l^\pm jj$$

$$l^\pm l^\mp l'^\pm$$

- **Main backgrounds** $WW, WZ, t\bar{t}$



MCFM Results on High-mass Higgs

- The only experimental signature which can be usefully analysed in MCFM is $l^+l^- + \text{missing energy}$. (This is the most significant according to HTZ)

Signal (m_H)	140	150	160	170	180	190
σ_{cuts} (fb)	4.36	5.32	6.12	5.15	3.90	2.47
σ_{total} (pb)	0.181	0.206	0.215	0.180	0.143	0.0976
Background	WW	$t\bar{t}$	ZZ	WZ		
σ_{cuts} (fb)	185	9.55	2.48	10.4		
σ_{total} (pb)	13.0	6.82	1.56	3.96		

Table 1: Signal and background cross-sections for a Higgs search with di-lepton final states. The total cross sections (without leptonic decays from the vector bosons) are also shown.



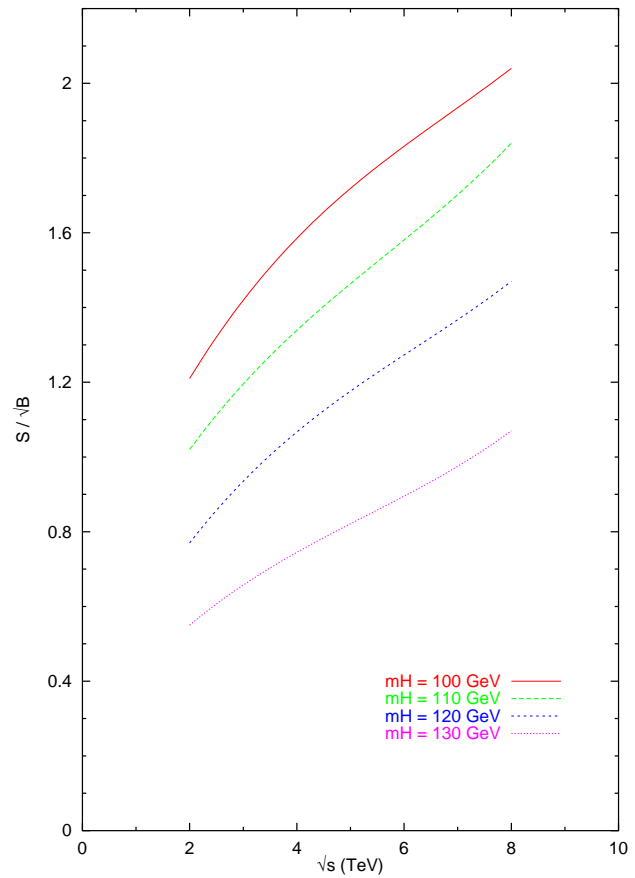
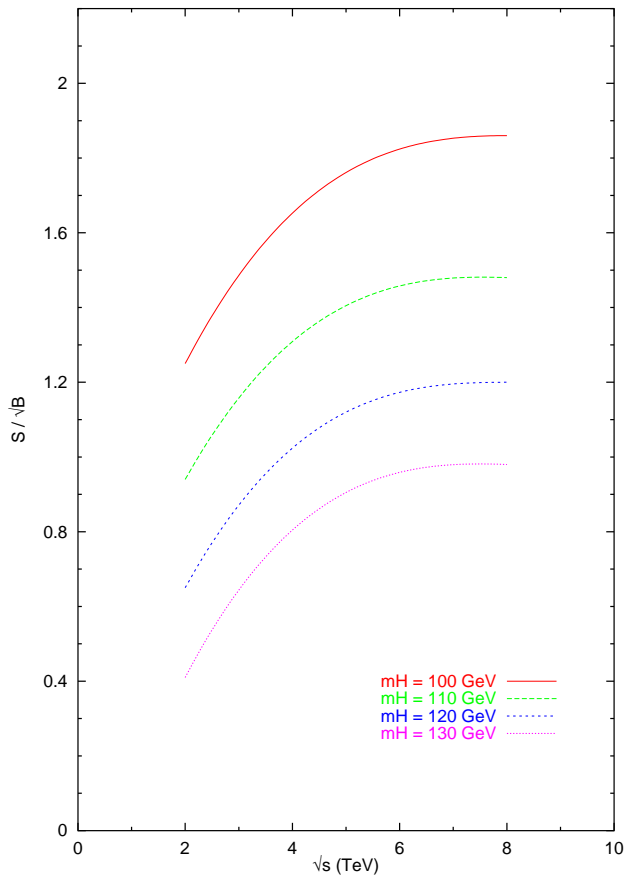
Comparison with HTZ

- We find that all the backgrounds due to the di-boson processes are larger than in **Han, Turcot and Zhang** (HTZ) because we have normalized to the $O(\alpha_s)$ cross-sections which are about 30% bigger than the Born cross-sections..
- Our signal cross sections are also larger than HTZ by about 20%; the net effect on S/\sqrt{B} may be small.
- We do not know exactly what the $gg \rightarrow H$ rate is (in the light of estimate of NNLO corrections by **Harlander**).
- The background from the WZ -class of events is twice as big as in HTZ because of our inclusion of the $W\gamma^*$ contribution.



Signal and backgrounds beyond 2 TeV

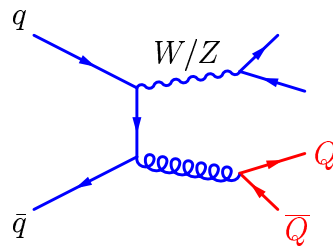
- S/\sqrt{B} as a function of collider energy for the WH (left) and ZH (right) processes



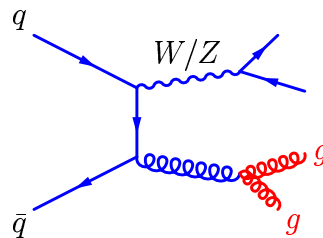
$W + 2$ jets: work in progress

- View the $W + 2$ jets process as an extension of the $Wb\bar{b}$ and $Zb\bar{b}$ calculations already performed:

- $Wb\bar{b}$ – part of $q\bar{q} \rightarrow W + q'\bar{q}'$



- $Zb\bar{b}$ – contains $gg \rightarrow Z + q\bar{q}$ + crossings

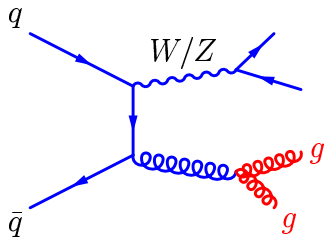


- There are extra parton configurations that we must count.
- The real radiation diagrams have extra singularities due to more configurations of soft/collinear gluons and collinear quark pairs.

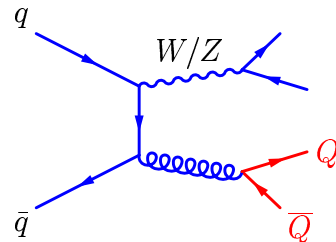


W + 2 jets at lowest order

- Separate the diagrams by colour structure:

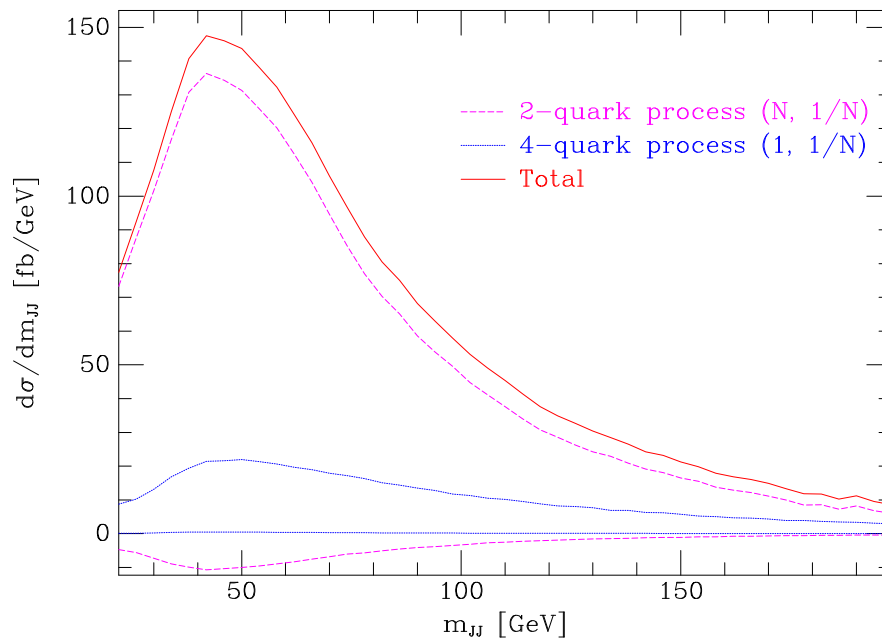


$$\propto N, 1/N$$



$$\propto 1, 1/N \times \delta_{qQ}$$

- Typical distribution, standard set of cuts:



$W + 2$ jets: strategy

$$|\mathcal{M}_{NLO}(Vq\bar{q}gg)|^2 \sim \quad 1 \quad \leftarrow \text{Near completion}$$
$$+ \frac{1}{N^2}$$
$$+ \frac{1}{N^4}$$

$$|\mathcal{M}_{NLO}(Vq\bar{q}Q\bar{Q})|^2 \sim \quad \frac{1}{N} \quad \leftarrow \text{Next target}$$
$$+ \frac{1}{N^3}$$
$$+ \frac{1}{N^2} \times \delta_{qQ}$$
$$+ \frac{1}{N^4} \times \delta_{qQ}$$

- Emphasis on $W + 2$ jet first



Conclusions

- $Wb\bar{b}$ background has approximate K factor already included in Susy-Higgs analysis.
- $Zb\bar{b}$ background is modified significantly (new class of diagrams). We believe that the significance will be less than the Susy-Higgs working group estimate.
- Our work underscores the importance of having an experimental determination of the $Zb\bar{b}$ and $Wb\bar{b}$ background, by relating them to $Vb\bar{b}$ events at lower $m_{b\bar{b}}$ or $V + 2\text{jet}$ events.
- We find that the rates used in the (HTZ) analysis too low, but the radiative corrections lead to little net change in S/\sqrt{B} .
- **More luminosity** is the panacea for all background ills.

