
Experimental Benchmarking of the Magnetized Friction Force

A.V. Fedotov¹, B. Galnander², V.N. Litvinenko¹, T. Lofnes²

A.O. Sidorin³, A.V. Smirnov³, V. Ziemann²

¹Brookhaven National Lab, Upton, NY 11973

²The Svedberg Laboratory, S-75121, Uppsala, Sweden

³JINR, Dubna, Russia

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High-energy cooling: need for accurate predictions of cooling times

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Cooling times for relativistic energies are much longer than for typical coolers:

$$\tau = \frac{A}{Z^2} \frac{\gamma^2}{4\pi r_p r_e n_e c \eta \Lambda_c} \left(\frac{\gamma \mathcal{E}_{in}}{\beta_{ic}} \right)^{3/2}$$

- standard (order of magnitude) estimate of cooling times for Au ion at RHIC storage energy of 100 GeV gives τ of the order of **1000 sec**, compared to a typical cooling time of the order of **0.1-1 sec** in existing coolers
- while an order of magnitude estimate was sufficient for typical coolers it becomes unacceptable for RHIC with a store time of a few hours and fast emittance degradation due to Intra Beam Scattering (IBS)



We need computer simulations which will give us cooling times estimates with an accuracy much better than an order of magnitude.

Motivation for comparison with formulas: accurate description of the Cooling Force

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Cooling Force studies

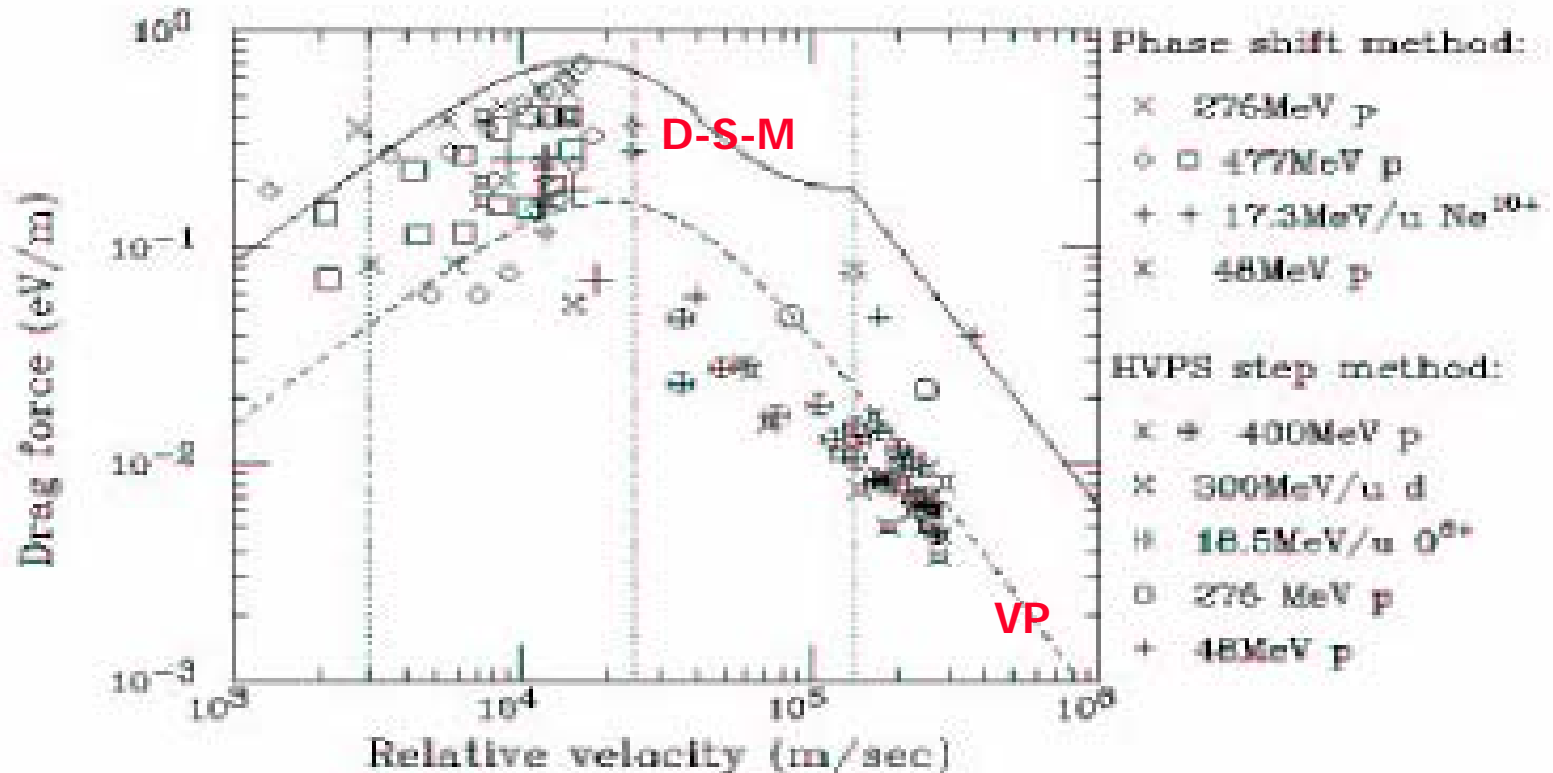


1. Benchmarking of available formulas vs VORPAL code (direct simulation of friction force) for various regimes.
 - D. Bruhwiler et al., AIP Conf. Proc. 773 (Bensheim, Germany, 2004), p.394.
 - A. Fedotov et al.; Bruhwiler et al., Proceedings of PAC'05 (Knoxville, TN, 2005).
 - A. Fedotov et al., "Detailed studies of Friction Force", this conference.
2. Experimental benchmarking:
 - (CELSIUS, December 2004 and March 2005)

Example of some previous comparison of experimental data with Derbenev-Skrinsky-Meshkov (D-S-M) and V.Parkhomchuk (VP) formulas.

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Y-N. Rao et al.: CELSIUS, Sweden'2001:



Motivation for our own data

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“One can compare formulas with simulations – since all the parameters used in simulations are known.”

“One cannot compare formulas with experiments – since many parameters in the experiments are unknown.”



This statement becomes especially true when one wants to use somebody’s else data without knowing all the details/conditions under which this data was taken.

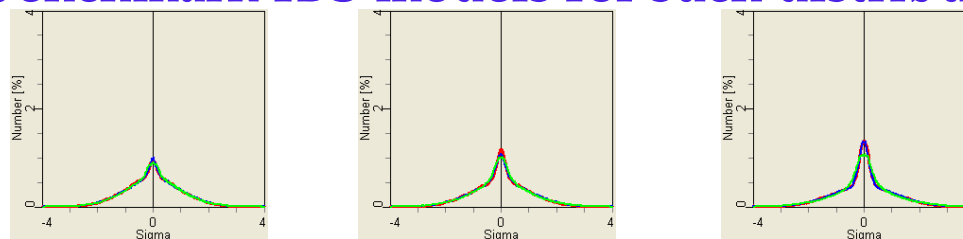
The way out is to do “well controlled” experiments – measure all the parameters which you need. And if you have uncertainty of some unknown parameters try to make an experiment which minimizes such uncertainty.

Major goals

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1. With well controlled experiments – systematically study friction force dependence on various parameters such as current, alignment angle, magnetic field.
2. Using low-energy cooler try to reproduce conditions possible at high-energy cooling:
 - 2.1) Different magnetization regimes – possible transition from good to bad magnetization
 - 2.2) Transient cooling – when as a result of slow cooling one first has clear formation of beam core with subsequent cooling of tails – need to benchmark IBS models for such distributions.

**very important
for collider**



Accuracy of Phase Shift method: important since it allows us to find exact location of the force maximum

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1. One needs to introduce small velocity difference between electrons and ions - **typically, voltage step is used to change energy of electrons.**
 2. One needs accurate measurement of the phase difference between the bunch and RF signal.

In our experiment at CELSIUS:

1. **Changing RF frequency** - allowed very fine steps in velocity difference (done before, for example, at IUCF).
2. Instead of network analyzer without phase lock loop the phase was measured by phase discriminator.

As a result, very accurate experimental data was obtained !

(see B. Galnander' presentation for more details)

Experiments at CELSIUS

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1. $B=0.1T$, current dependence: ($I_e=500mA, 250mA, 100mA, 20 mA$)

Measure all needed parameters, including parameters of ion distribution.

2. Dependence on $V_{\text{effective}}$:

- measured for several values of tilt in both horizontal and vertical direction - both negative and positive directions.
- always recorded longitudinal and transverse sigmas to perform accurate convolution over distributions. Measured values are close to those predicted by BetaCool simulations
- did calibration of tilt angle with both BPM's and H^0 monitor

Check with available theory.

3. Measured “transient cooling”
(IBS+COOLING) both for longitudinal and
transverse profiles:

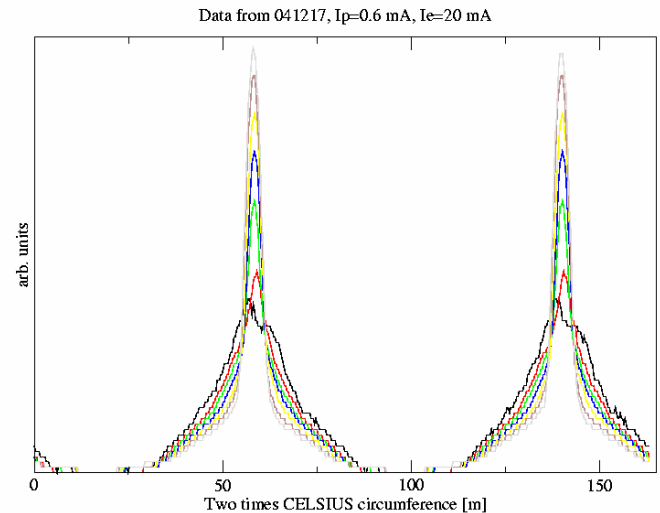
Test models of IBS for non-Gaussian
distribution -needed for high-energy
cooling.

4. Various values of B with various currents:
 $I_e=500\text{mA}$, 300mA , 100mA , 50 mA ($B=0.03$,
 0.04 , 0.05 , 0.06 , 0.08 , 0.1 , 0.12T)

Study various regimes of magnetization -
needed for high-energy cooling.

5. Effects of solenoid errors.

Study description via $V_{\text{effective}}$.



Magnetized logarithm:
 $L_M=1.5 \rightarrow 0.7$

V. Parkhomchuk's (VP) empiric formula

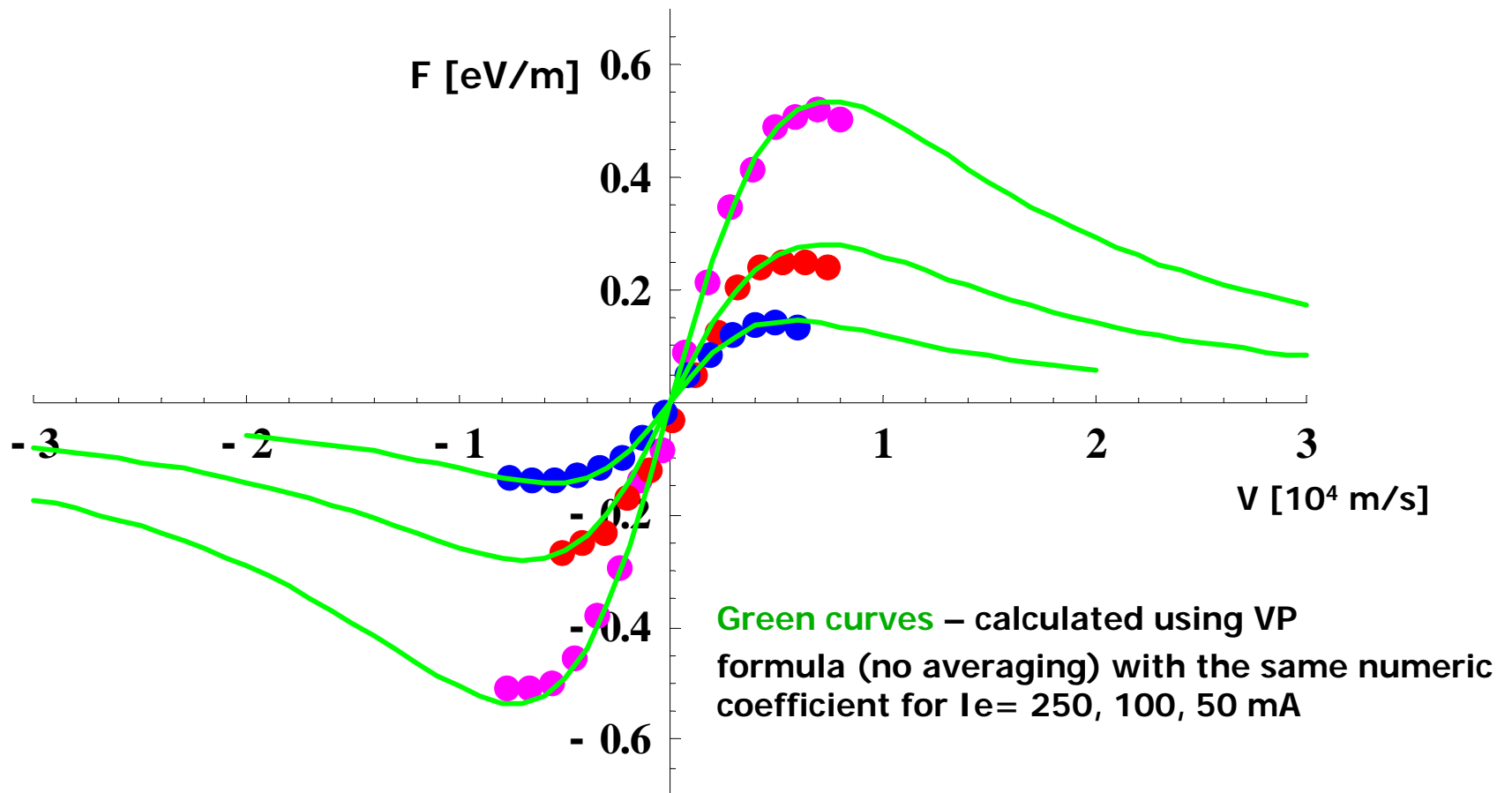
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empiric formula (VP) – single-particle formula

$$\mathbf{F} = -\frac{1}{\pi} \omega_{pe}^2 \frac{(Ze)^2}{4\pi\epsilon_0} \ln\left(\frac{\rho_{\max} + \rho_{\min} + r_L}{\rho_{\min} + r_L}\right) \frac{\mathbf{V}_{ion}}{(V_{ion}^2 + V_{eff}^2)^{3/2}}$$

March 2 data: $B=0.1\text{T}$, electron current $I_e=250$ (pink color), 100 (red), 50 (blue) mA

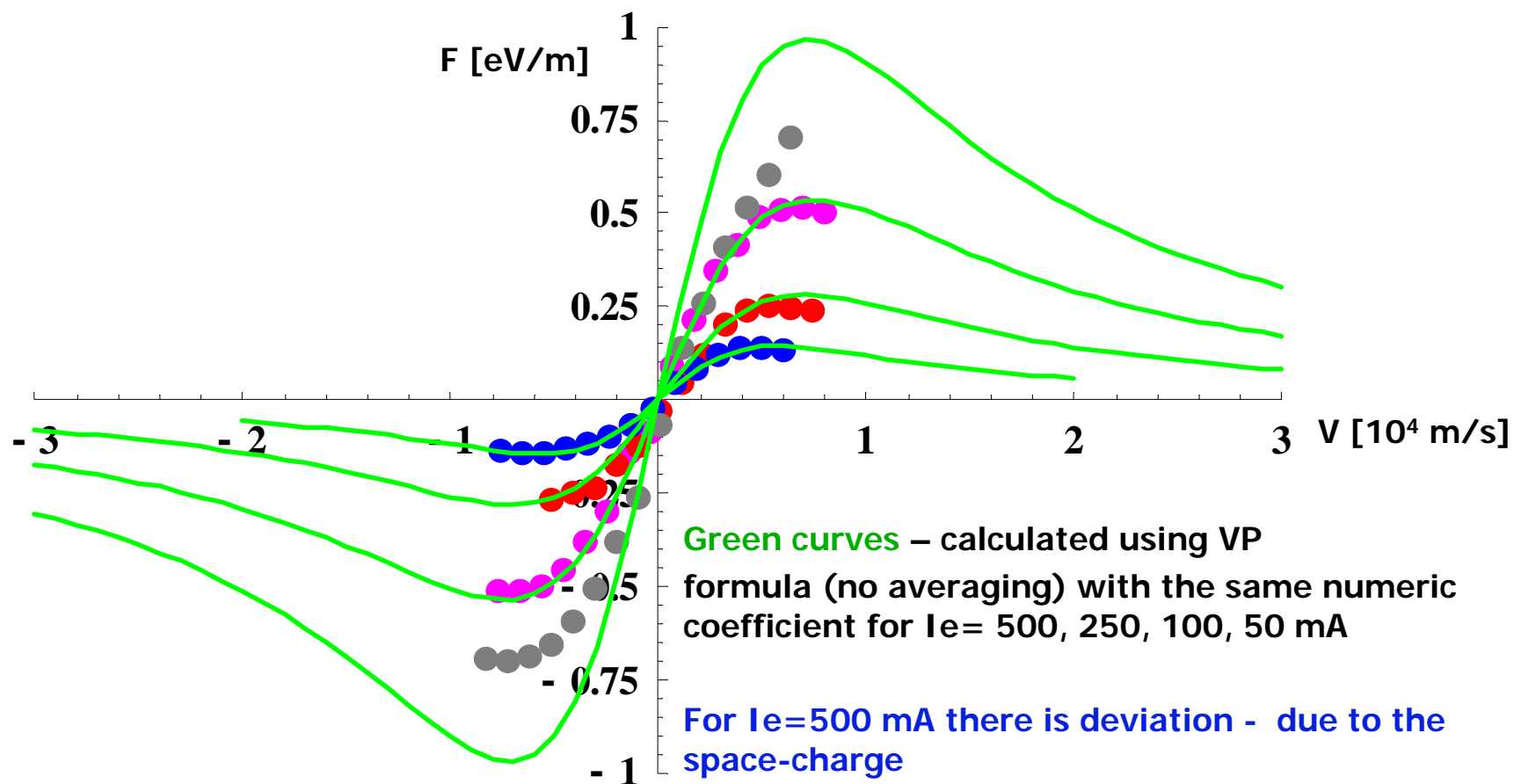
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March 2 data: $B=0.1T$

$I_e=500$ (gray), 250 (pink), 100 (red), 50 (blue) mA

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Electron current $I_e=500\text{mA}$

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For high currents of the electron beam the space-charge of the electron beam becomes important:

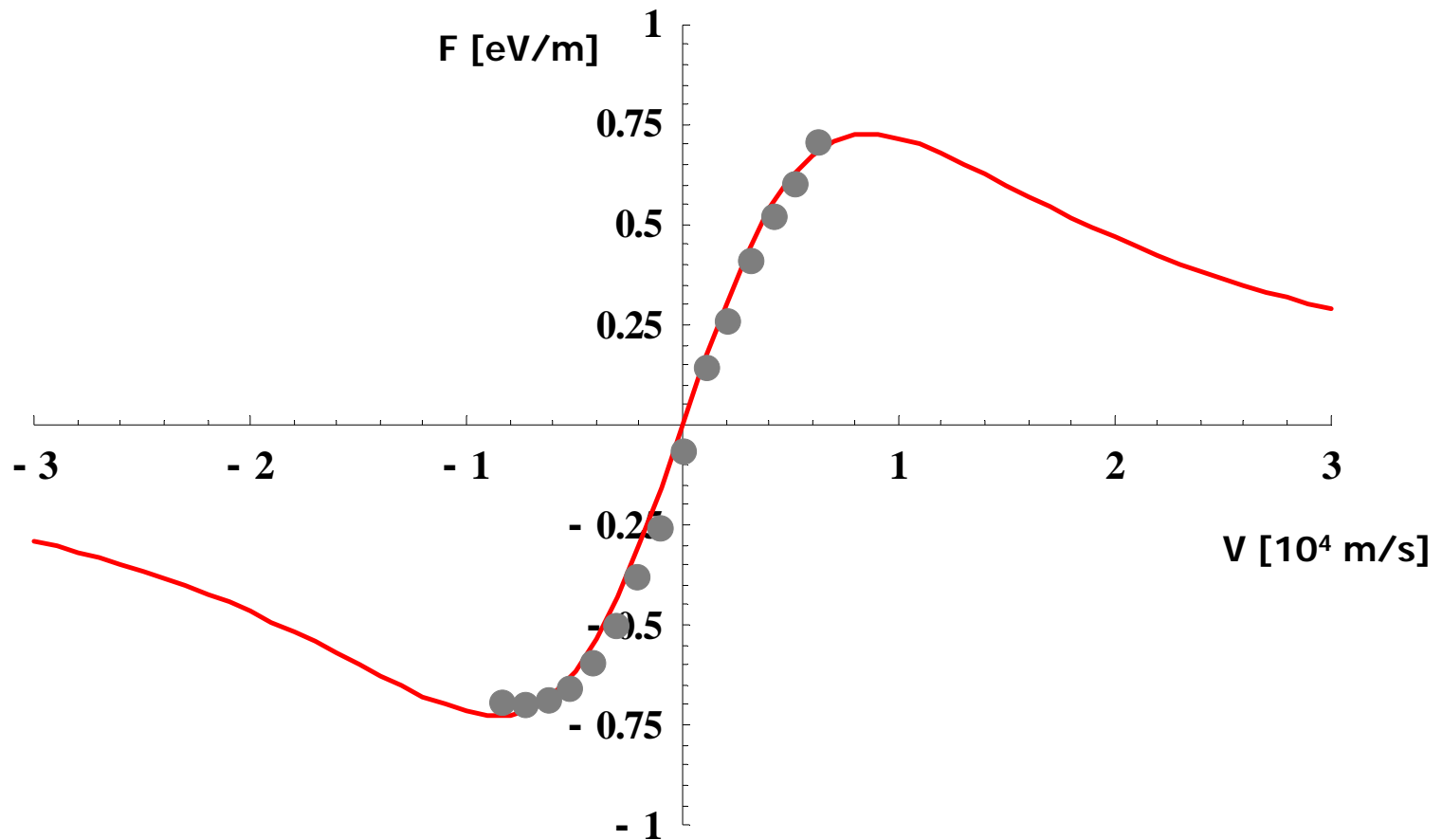
The electron drift in crossed fields – the electric and magnetic fields of the electron beam and longitudinal magnetic field of the cooler:

$$v_d = \frac{2I}{B\beta\gamma^2} \frac{r}{a^2}$$

For measured distribution of the proton beam for the case under comparison (March 2, set#23, $B=0.1\text{T}$, $I_e=500\text{mA}$) - $V_{\text{drift}}=6-7 \cdot 10^3\text{m/s}$ – which is an additional contribution to $V_{\text{effective}}$ in the cooling force formulas.

March 2 data: $I_e=500\text{mA}$, $B=0.1\text{T}$ - formula vs experiment with additional contribution to $V_{\text{effective}}$ from V_{drift}

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Fits with single-particle formulas

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1. Current dependence - friction force scales linearly with current/density - as expected from formula.
2. Numeric coefficient for the force is in agreement with the one in Parkhomchuk's formula. Also, it can be adjusted to agree with Derbenev's coefficient (which results in only slightly different effective velocity) - the coefficients are similar for the region of low relative velocities ($1/\pi$ vs $1/(2\pi)^{1/2}$).
3. Note that Coulomb logarithm depends on relative ion velocity and $V_{\text{effective}}$ - fitting was done with such velocity-dependent logarithm.
4. Fitted $V_{\text{effective}}$ has very weak current dependence:

$0.74\text{-}0.78 \cdot 10^4$ m/s

Observations

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- Using single-particle formula allows to fit experimental data and extract $V_{\text{effective}}$.
 - However, since rms velocity spreads of cooled proton beam are significant (for our measurements, we would need to have $dp/p=1e-5$ and $\epsilon=1e-9$ μm to neglect this effect, while parameter of the proton beam with which we did measurements typically had about $dp/p=5e-5$ and $\epsilon=5e-8$ μm), fitted $V_{\text{effective}}$ has contribution from this effect.

The accurate procedure is then to measure rms velocities of the distribution and average single-particle formulas over the proton distribution.



This was done for all 10's of friction force curves which were measured for various parameters

Detailed comparison: Averaging over ion distribution

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$$\langle F \rangle = C \frac{4\pi Z^2 e^4 n_e}{m\sqrt{2\pi}\Delta_{\perp}^2\Delta_{\parallel}} \int_0^{\infty} \int_{-\infty}^{\infty} \frac{v_{\parallel} L_M(v_{\perp}, v_{\parallel}, v_{eff})}{(v_{\perp}^2 + v_{\parallel}^2 + v_{eff}^2)^{3/2}} \exp\left(-\frac{v_{\perp}^2}{2\Delta_{\perp}^2} - \frac{(v_{\parallel} - v_0)^2}{2\Delta_{\parallel}^2}\right) v_{\perp} dv_{\parallel} dv_{\perp}$$

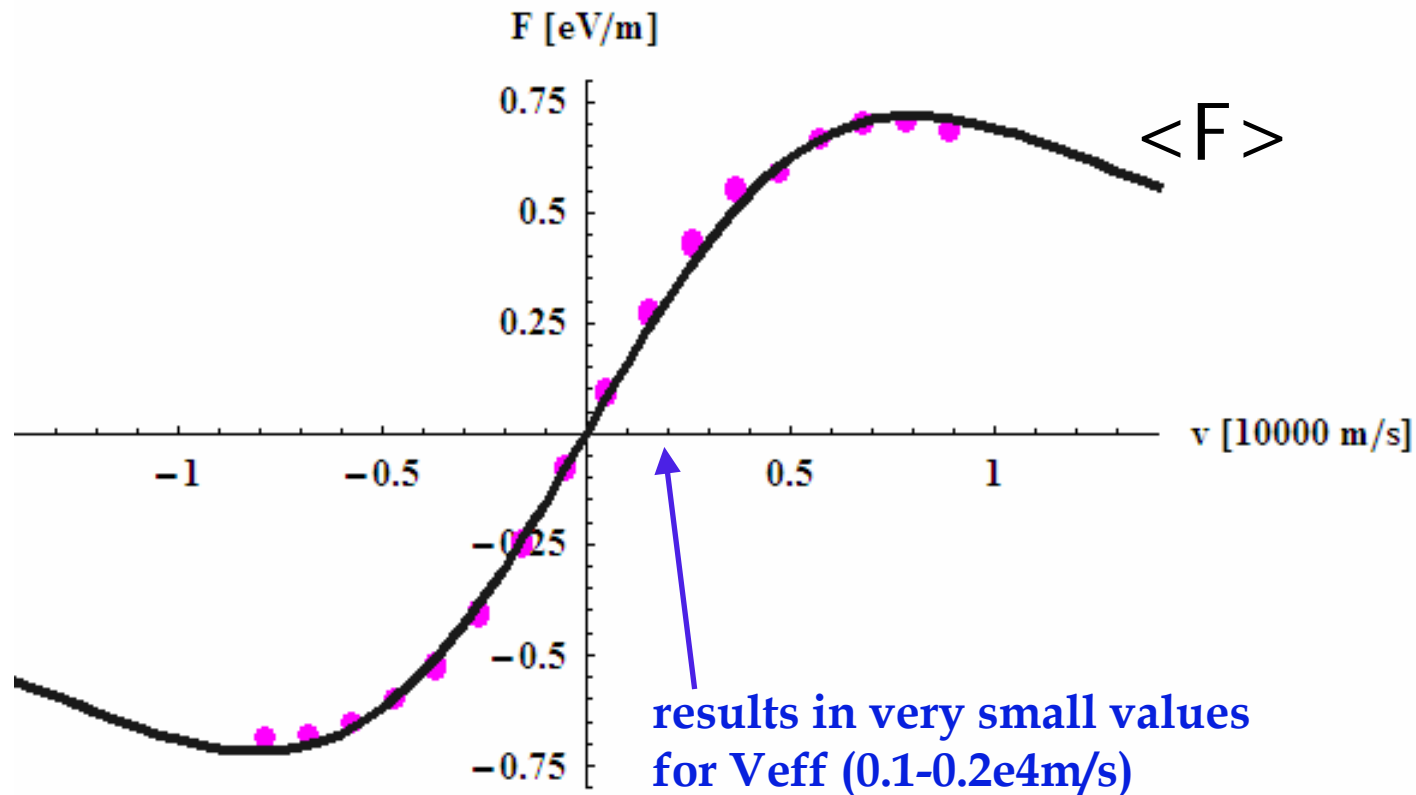
rms parameters of proton beam were measured for each measurement of friction force curve.

1. **First approach:** assume C is known and treat V_{eff} as fitting parameter.
2. **Second approach:** assume V_{eff} is known from measurements and treat C as fitting parameter.

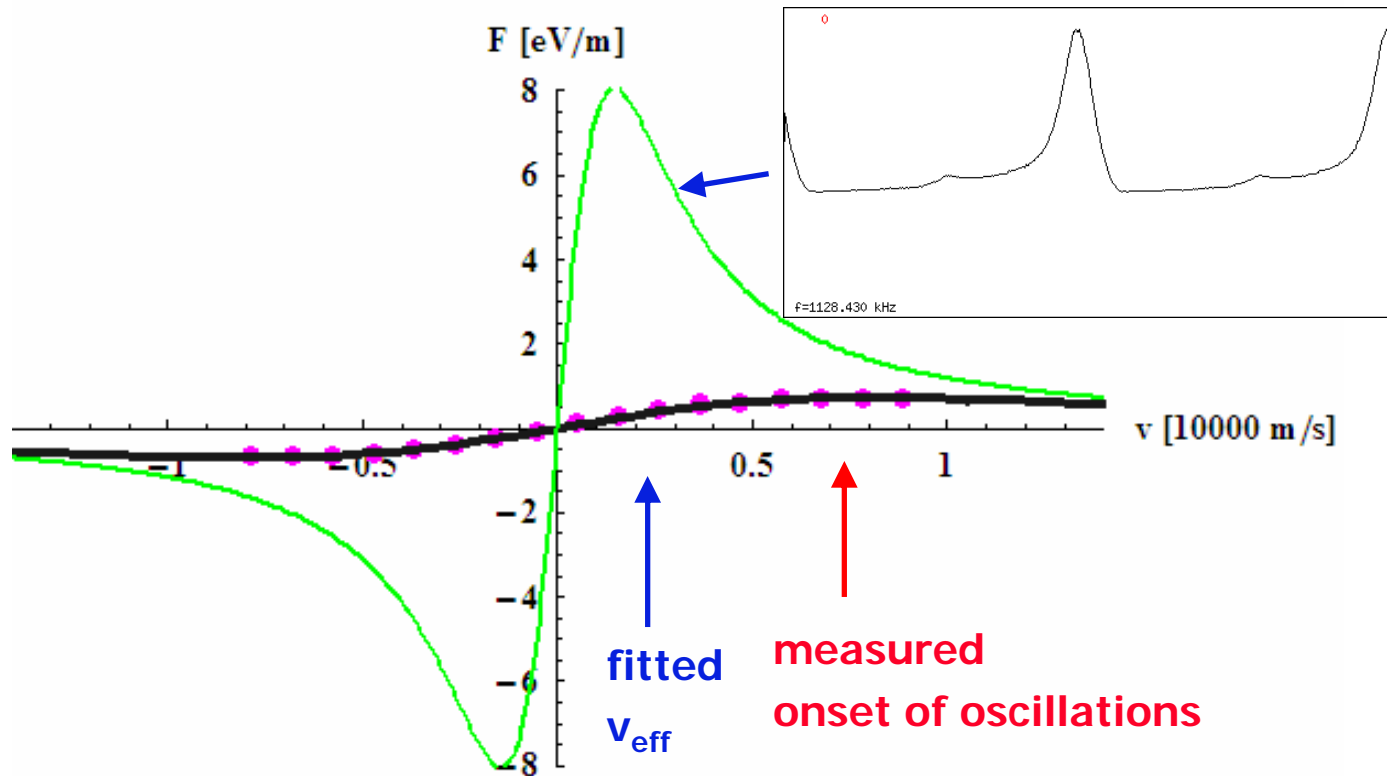
$B=0.12\text{T}$, $I_e=300\text{mA}$

Friction force averaged over proton distribution with measured rms velocity spread

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First approach - one fitting parameter V_{eff}



Longitudinal profiles:
expected onset of oscillations for small v_{eff}

fitted v_{eff}
measured onset of oscillations

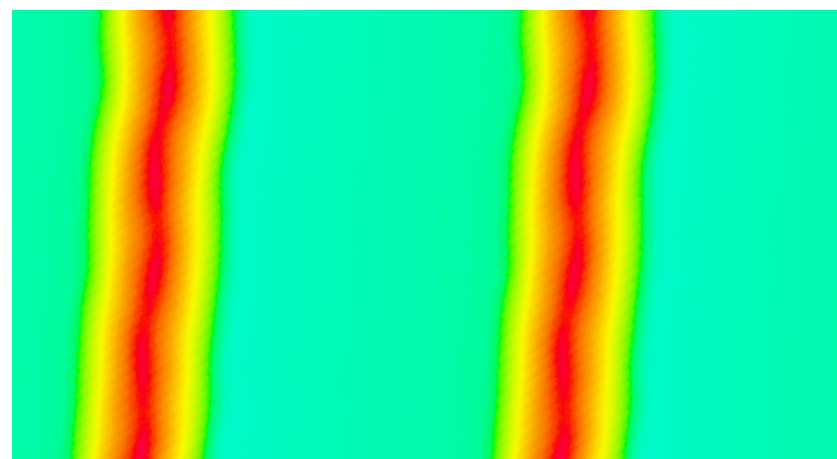
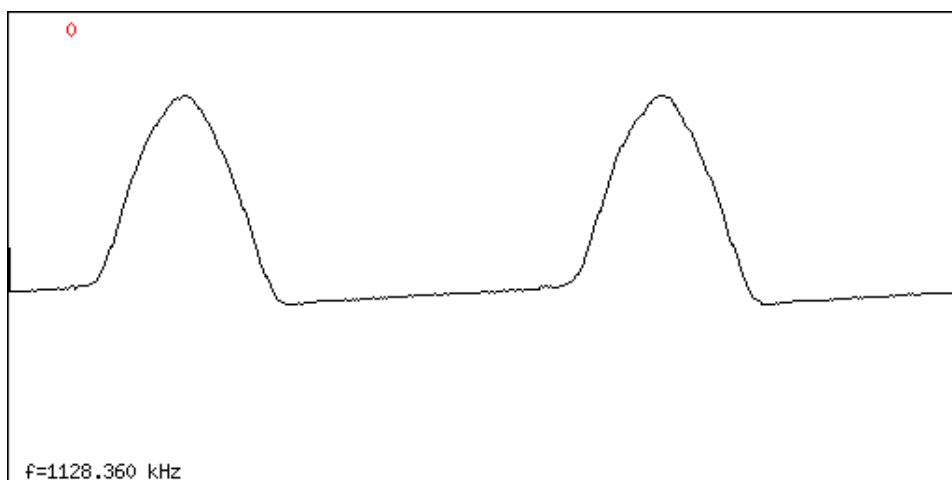
$$\langle F \rangle = C \frac{4\pi Z^2 e^4 n_e}{m\sqrt{2\pi}\Delta_{\perp}^2\Delta_{\parallel}} \int_0^{\infty} \int_{-\infty}^{\infty} \frac{v_{\parallel} L_M(v_{\perp}, v_{\parallel}, v_{eff})}{(v_{\perp}^2 + v_{\parallel}^2 + v_{eff}^2)^{3/2}} \exp\left(-\frac{v_{\perp}^2}{2\Delta_{\perp}^2} - \frac{(v_{\parallel} - v_0)^2}{2\Delta_{\parallel}^2}\right) v_{\perp} dv_{\parallel} dv_{\perp}$$

C = 1/π

Measurements of longitudinal friction force maximum

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Approaching friction force maximum

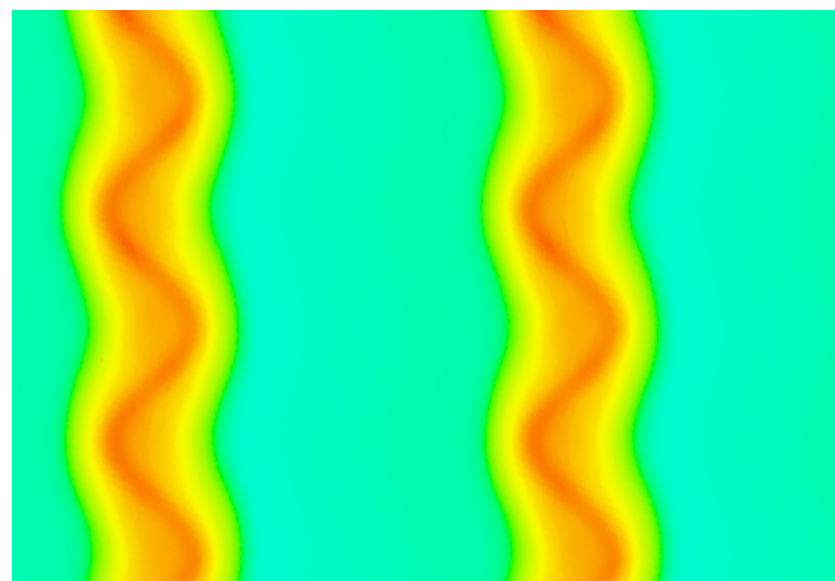
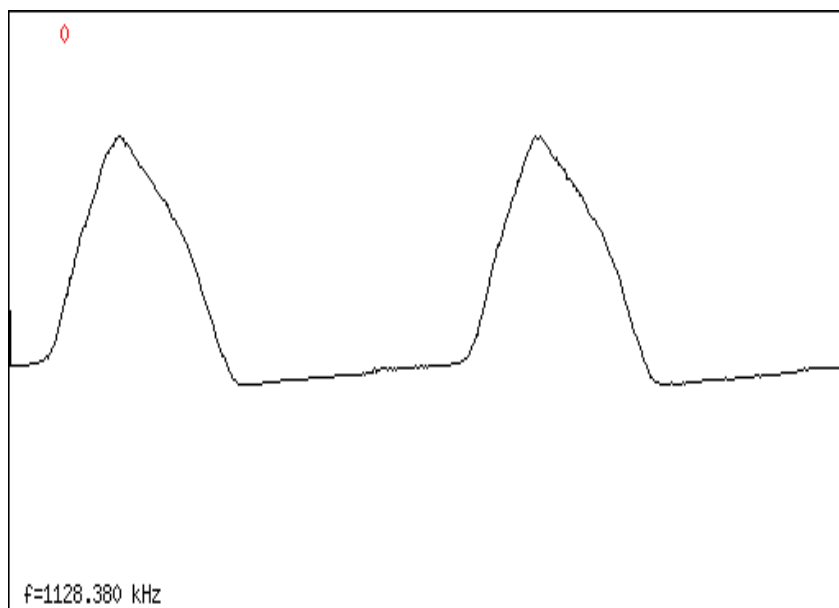


Longitudinal profiles

Measurements of longitudinal friction force maximum

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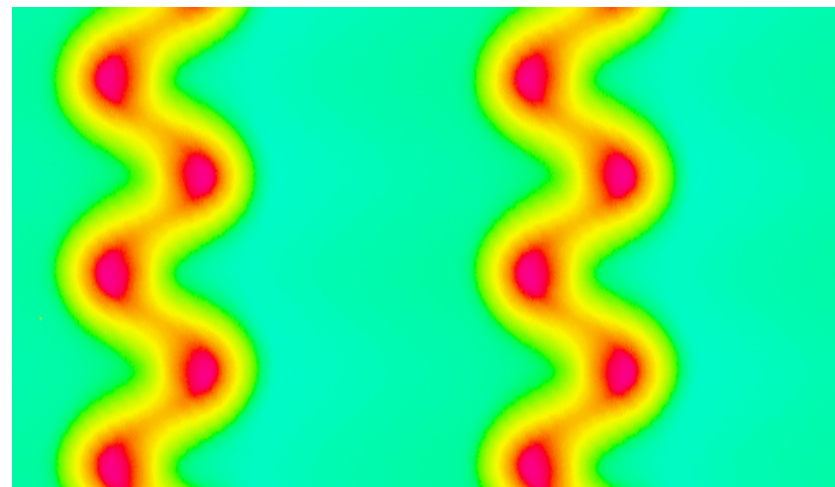
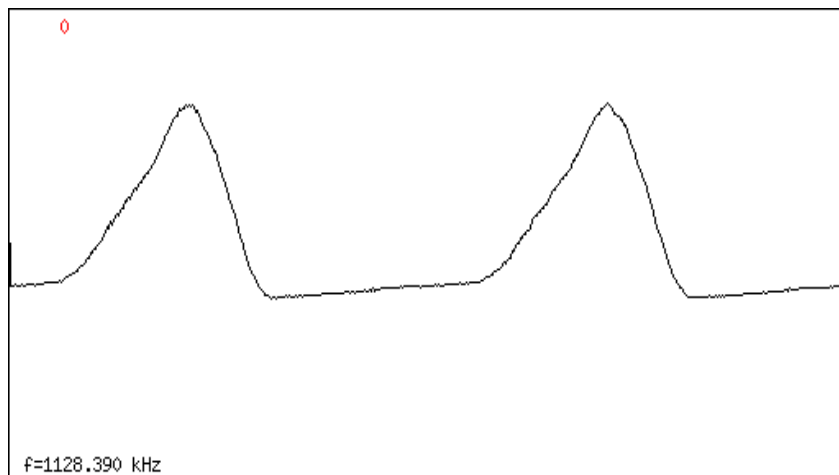
just past the maximum



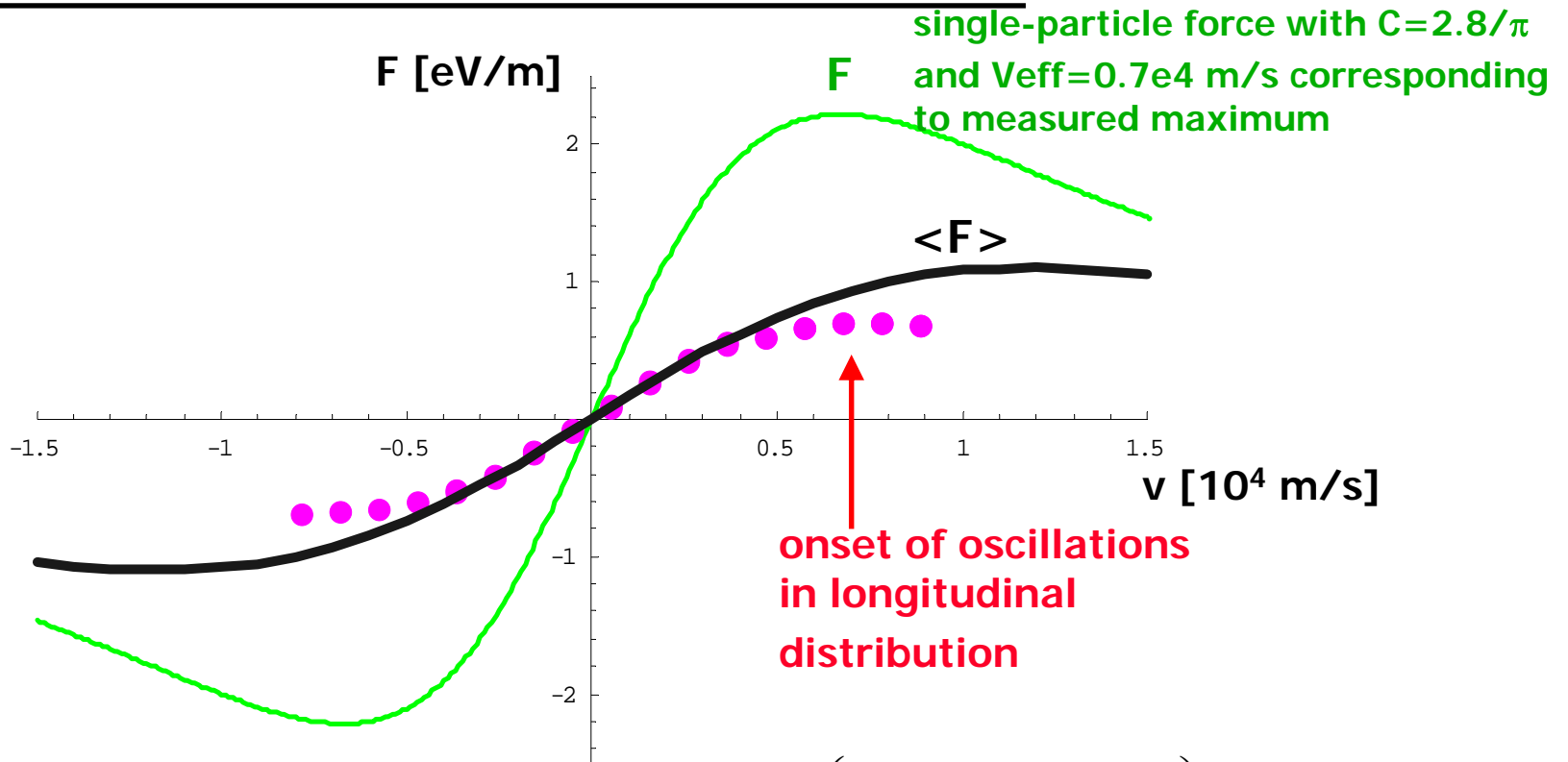
Measurements in non-linear part of the friction force

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far past the maximum

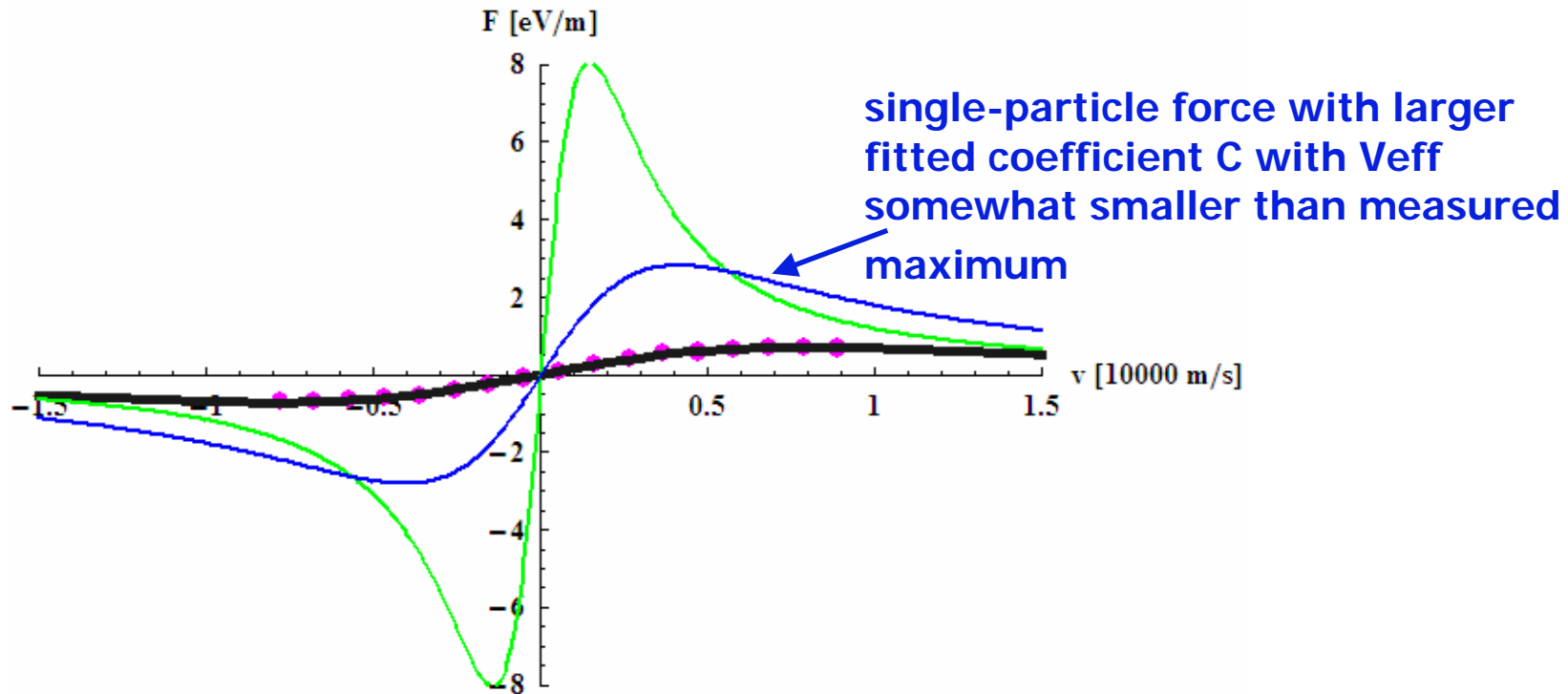


Second approach - one fitting parameter C (with measured V_{eff})



$$\langle F \rangle = C \frac{4\pi Z^2 e^4 n_e}{m\sqrt{2\pi}\Delta_{\perp}^2 \Delta_{\parallel}} \int_0^{\infty} \int_{-\infty}^{\infty} \frac{v_{\parallel} L_M(v_{\perp}, v_{\parallel}, v_{eff})}{(v_{\perp}^2 + v_{\parallel}^2 + v_{eff}^2)^{3/2}} \exp\left(-\frac{v_{\perp}^2}{2\Delta_{\perp}^2} - \frac{(v_{\parallel} - v_0)^2}{2\Delta_{\parallel}^2}\right) v_{\perp} dv_{\parallel} dv_{\perp}$$

Second and 1/2 approach - basically, both C and Veff are fitting parameters (plus averaging)



$$\langle F \rangle = C \frac{4\pi Z^2 e^4 n_e}{m\sqrt{2\pi}\Delta_{\perp}^2\Delta_{\parallel}} \int_0^{\infty} \int_{-\infty}^{\infty} \frac{v_{\parallel} L_M(v_{\perp}, v_{\parallel}, v_{eff})}{(v_{\perp}^2 + v_{\parallel}^2 + v_{eff}^2)^{3/2}} \exp\left(-\frac{v_{\perp}^2}{2\Delta_{\perp}^2} - \frac{(v_{\parallel} - v_0)^2}{2\Delta_{\parallel}^2}\right) v_{\perp} dv_{\parallel} dv_{\perp}$$

Summary - benchmarking of experiments

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At CELSIUS, we were able to measure longitudinal friction force with very good precision which allows us to use experimental data for accurate benchmarking of theory and simulations.

A careful experimental study of various parameters was performed:

- 1) Current dependence
- 2) Dependence of tilt between electron and proton beams
- 3) Dependence on solenoid errors
- 4) Various degrees of magnetization
- 5) Transient cooling

Benchmarking of experimental data for each of the experiments is presently in progress.

Acknowledgements

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