

*Major!*

# Progress in lattice QCD ?

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*Howard Trottier*

*Simon Fraser University  
(HPQCD Collaboration)*

*Aspen Winter Conference  
February 2006*

# (1) Precision *Unquenched* LQCD

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Progress has been made: “Greatest Hits”

Some outstanding theoretical issues

# (2) Role of Perturbation Theory

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# (3) Recent phenomenology

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See Ian Shipsey’s talk for more theory vs. experiment

# (4) Other Approaches

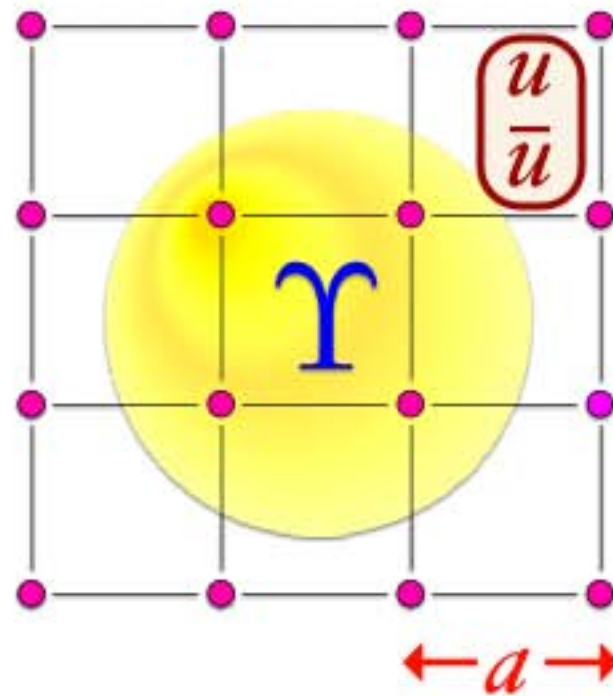
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# (1) LQCD: *Progress has been made*

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For LQCD to be relevant to experiment:

(i) Include the effects of virtual  $u, d, s$  quarks: *Unquench*  
(and do so at sufficiently small quark masses  $\rightarrow$  physical point)

(ii) Get sufficiently close to **continuum limit**

(iii) **Quantify** the statistical & systematic **errors**



# HPQCD & Friends: *Greatest Hits*

$$\langle \mathcal{O} \rangle = \int [dU_\mu(x)][d\bar{\psi}d\psi] \mathcal{O} e^{-\beta(S_{\text{gluon}} + S_{\text{quark}}^{\text{stagg}})} \leftarrow \text{!}$$

▶ Only **5** input parameters (same as in continuum QCD)

▶  $m_u (= m_d), m_s, m_c, m_b, a$  ( $\Leftrightarrow \alpha_s$ )

▶  $m_{u/d} \leftarrow m_\pi^2$

▶  $m_s \leftarrow 2m_K^2 - m_\pi^2$

▶  $m_c \leftarrow m_D$

▶  $m_b \leftarrow m_\Upsilon$

▶  $a \leftarrow m_{\Upsilon'} - m_\Upsilon$

Each experimental quantity roughly  $\propto$  the one  $m_{\text{quark}}$  and roughly independent of the other masses

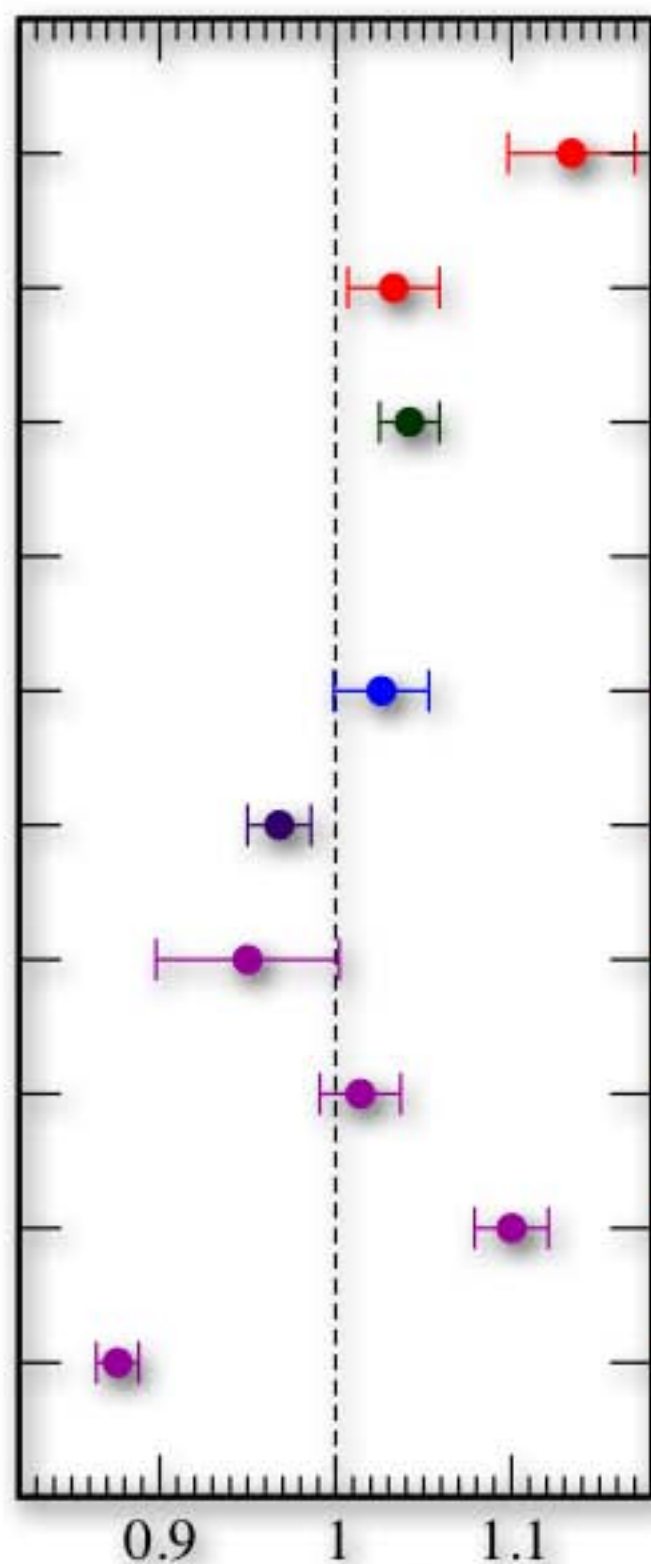
This mass difference roughly independent of all  $m_{\text{quark}}$ 's

▶ Same  **$a$**  using *any* input:  $f_\pi, m_K, m_{B_s}, \dots$

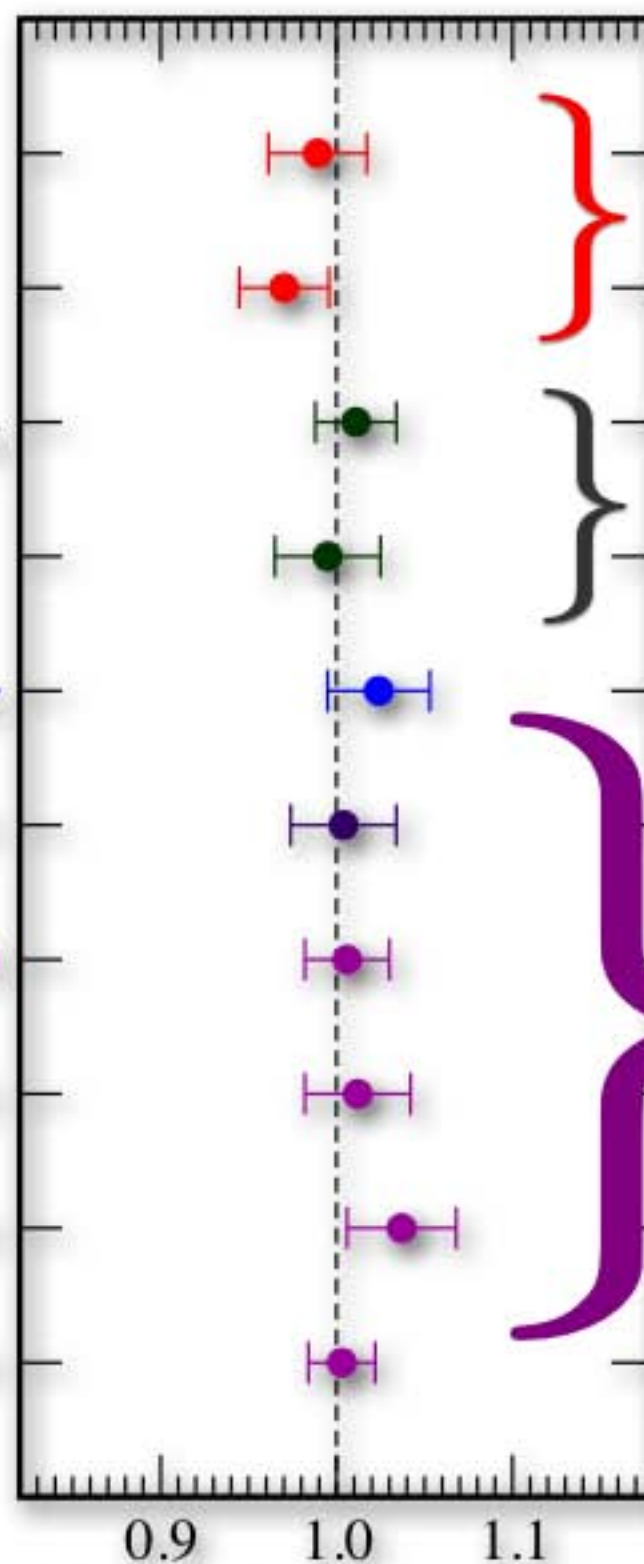


# LQCD / Experiment

HPQCD,  
MILC,  
UKQCD,  
Fermilab,  
Collab'ns  
PRL 92,  
2004



$f_\pi$   
 $f_K$   
 $3m_\Xi - m_N$   
 $m_\Omega$   
 $2m_{B_s} - m_\Upsilon$   
 $\psi(1P-1S)$   
 $Y(1D-1S)$   
 $Y(2P-1S)$   
 $Y(3S-1S)$   
 $Y(1P-1S)$



Light-quarks  
&  
form-factors

Baryons

Heavy quarks:

Note "unified"  
treatment of  
e.g.  $B$  &  $\Upsilon$

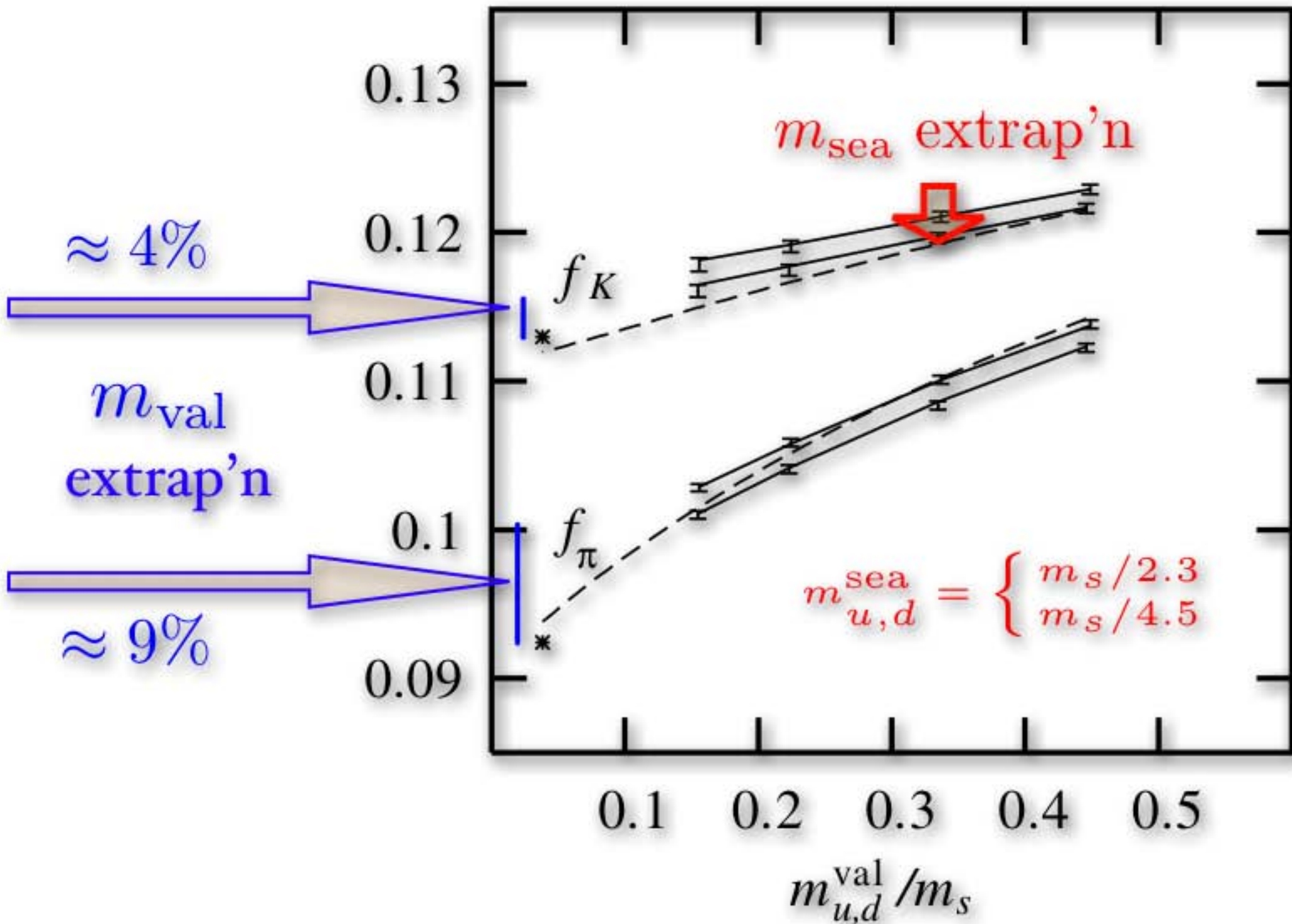
Errors  
 $\sim 10-15\%$  Quenched/Experiment

Errors  
 $< 3\%$  ( $N_f=2+1$ )/Experiment

# Partially Quenched $\chi$ PT

Lee & Sharpe  
Bernard & Aubin

Run at several  $m_{u/d,s}^{\text{valence}}$  not necessarily equal to  $m_{\text{sea}}$



Sea quark masses  
 $\sim 3$ - $10$  times smaller than in previous unquenched

Recent  $m_u/m_s \sim 1/10$



# LQCD can't do everything

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(not for now!)

- ▶ Unstable hadrons e.g.  $\rho \rightarrow \pi\pi \rightarrow \rho$ 
  - ▶ intermediate particles propagate to lattice boundaries, induces large finite volume errors
- ▶ Hadrons near decay thresholds e.g.  $\psi' \rightarrow D\bar{D} \rightarrow \psi'$
- ▶ At most one hadron in initial and final states

## “Gold-plated” LQCD quantities

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- ▶ Narrow/Stable hadrons sufficiently below threshold
  - ▶ e.g.  $\pi, K, p, D, D_s, B, B_s, J/\psi, \Upsilon, \Upsilon', \dots$  but not  $\rho, D^*, \psi', \dots$
- ▶ Semileptonic ✓ while nonleptonic ✗
- ▶ CLEO-c offers crucial validation (Ian Shipsey's talk)

# “Gold-plated” LQCD meets CKM

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$V_{ud}$

$\pi \rightarrow \ell\nu$

$V_{cd}$

$D \rightarrow \ell\nu$

$D \rightarrow \pi\ell\nu$

$\langle B_d | \overline{B}_d \rangle$

$V_{us}$

$K \rightarrow \ell\nu$

$K \rightarrow \pi\ell\nu$

$V_{cs}$

$D_s \rightarrow \ell\nu$

$D \rightarrow K\ell\nu$

$\langle B_s | \overline{B}_s \rangle$

$V_{ub}$

$B \rightarrow \pi\ell\nu$

$V_{cb}$

$B \rightarrow D\ell\nu$

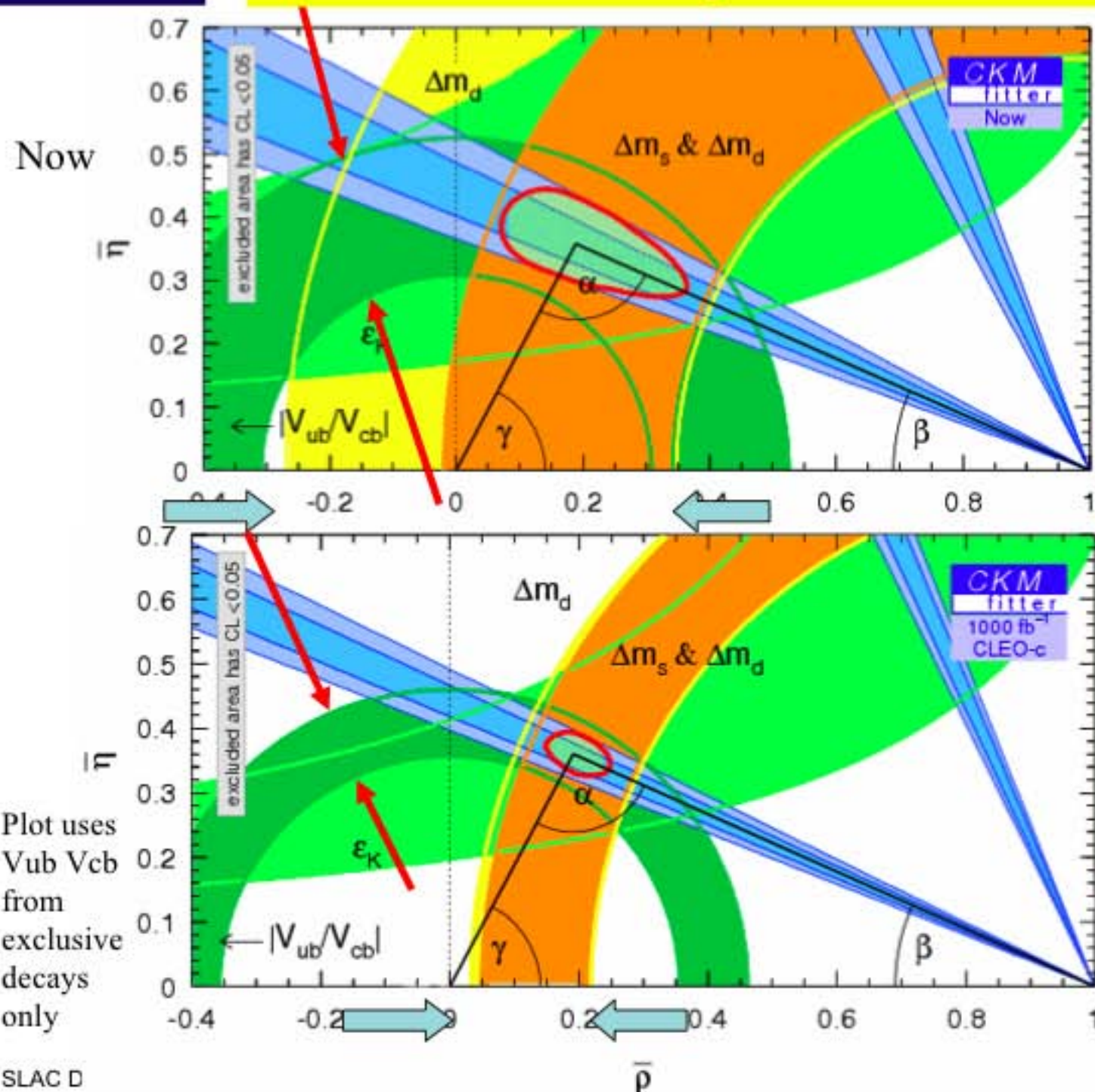
$V_{tb}$



# CKM: Present & Future (Ian Shipsey)



Precision theory + charm = large impact



Plot uses  $V_{ub} V_{cb}$  from exclusive decays only

Theoretical errors dominate width of bands

precision QCD calculations tested with precision charm data at threshold  
 → theory errors of a few % on B system decay constants & semileptonic form factors

500 fb<sup>-1</sup> @ BABAR/Belle



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# A potential pit-fall for HPQCD

- ▶ Staggered quarks **cheap to simulate**: remnant  $\chi$  symmetry
  - ▶ Staggered quarks plagued by **flavour doubling**:

$$\bar{\psi}(x) \gamma \cdot D \psi(x) \Rightarrow \sin(p_\mu a) \times \bar{\psi}(p) \gamma_\mu \psi(p)$$

$\Rightarrow$  low-energy modes at  $p_\mu = 0, \pi/a$

$\therefore 2^4$  copies (“tastes”) (reduce to 4 tastes by “staggering”)

To get desired **(2+1)**-flavours instead of 4  
 $\det(\gamma \cdot D + m) \rightarrow \det(\gamma \cdot D + m)^{1/4}$

☂ potentially worrisome **non-localities**

$$\det(\gamma \cdot D + m) \rightarrow \det(\gamma \cdot D + m)^{1/4}$$

**Is this a local effective theory ?**

- ☺ Correct to all orders in PT (Batrouni et al 1985)
- ☺ Chiral anomalies correctly handled (Sharatchandra et al 1981; Smit & Vink 1988)

**☺ Controlled by short-distance interactions**

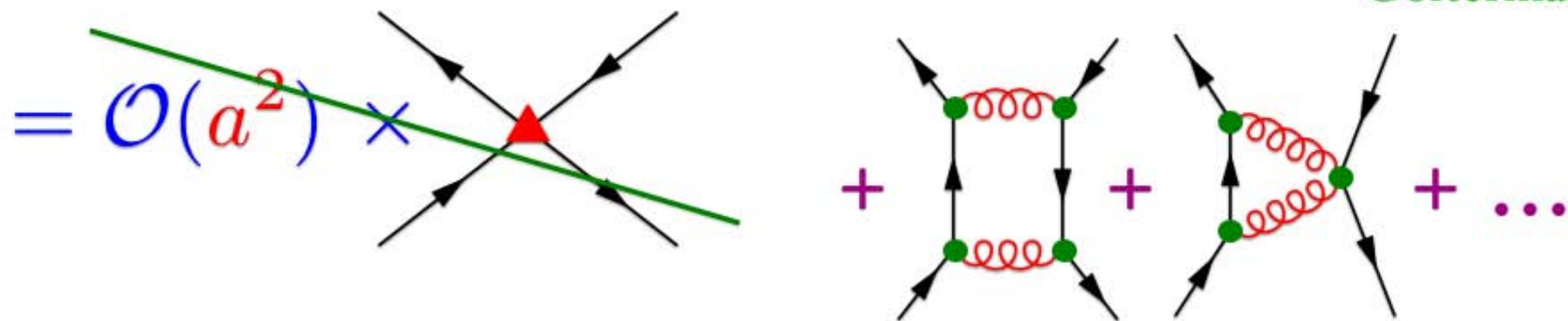
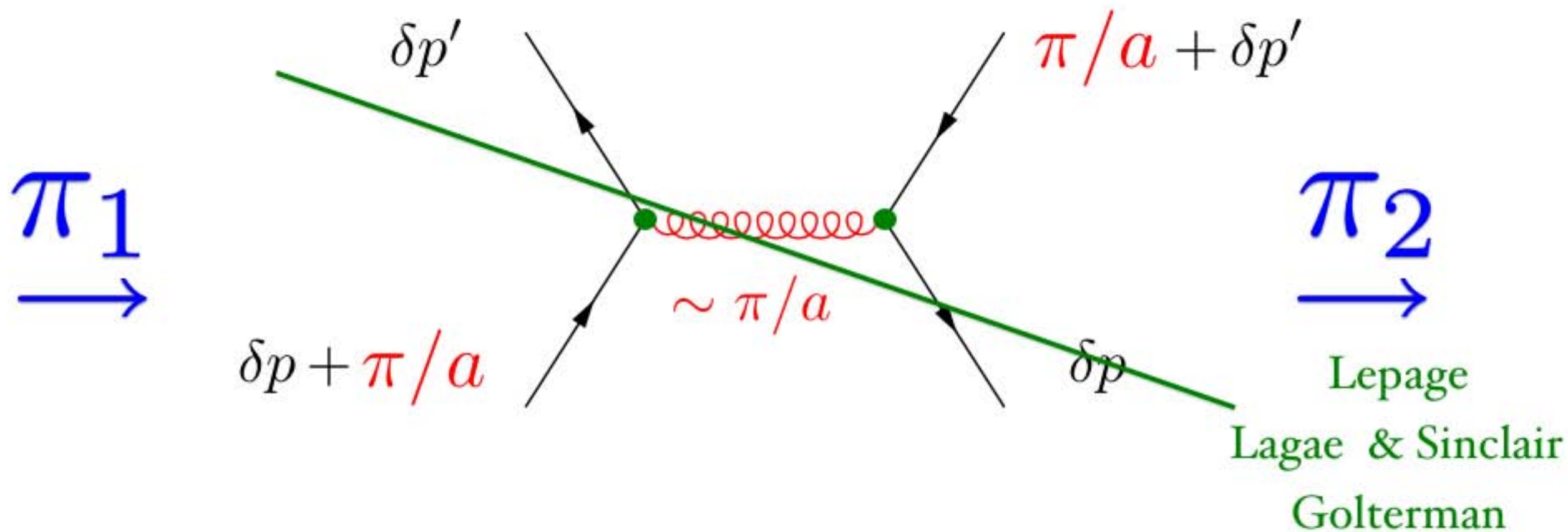
$$[\det(\gamma \cdot D + m)]^{1/4} = [\prod_n (\lambda_n + m)]^{1/4}$$

- ▶ 1/4 root OK if eigenvalues are  $\approx$  quadruply degenerate
- ▶ eigenvalue- or “taste-” splittings controlled by short-distance interactions





# Taste-changing interactions

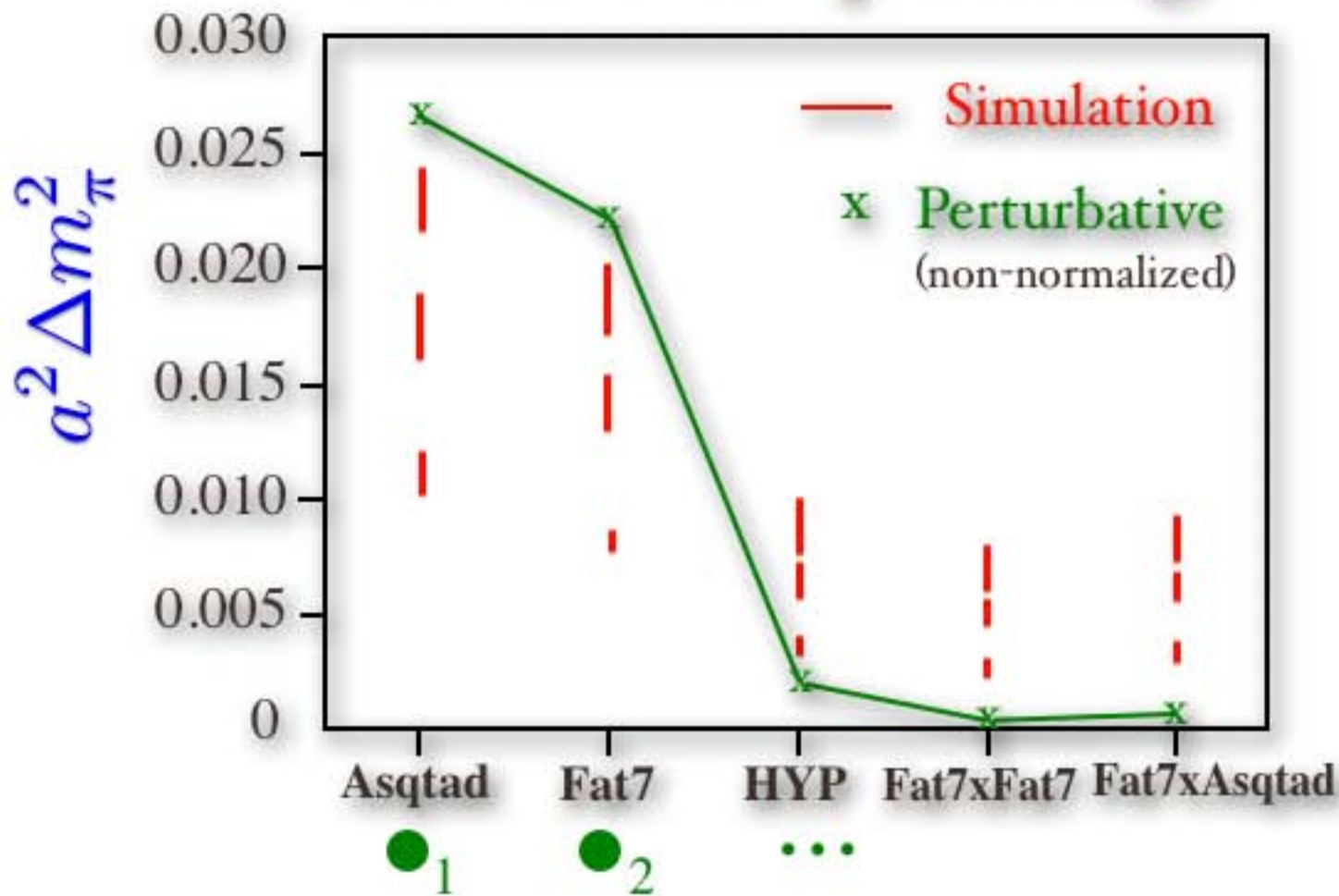


Minimize taste-changing interactions by perturbative analysis of the effective interaction ●

Gave us the “improved” staggered action in current use ... do more!

# Pion "taste" splittings

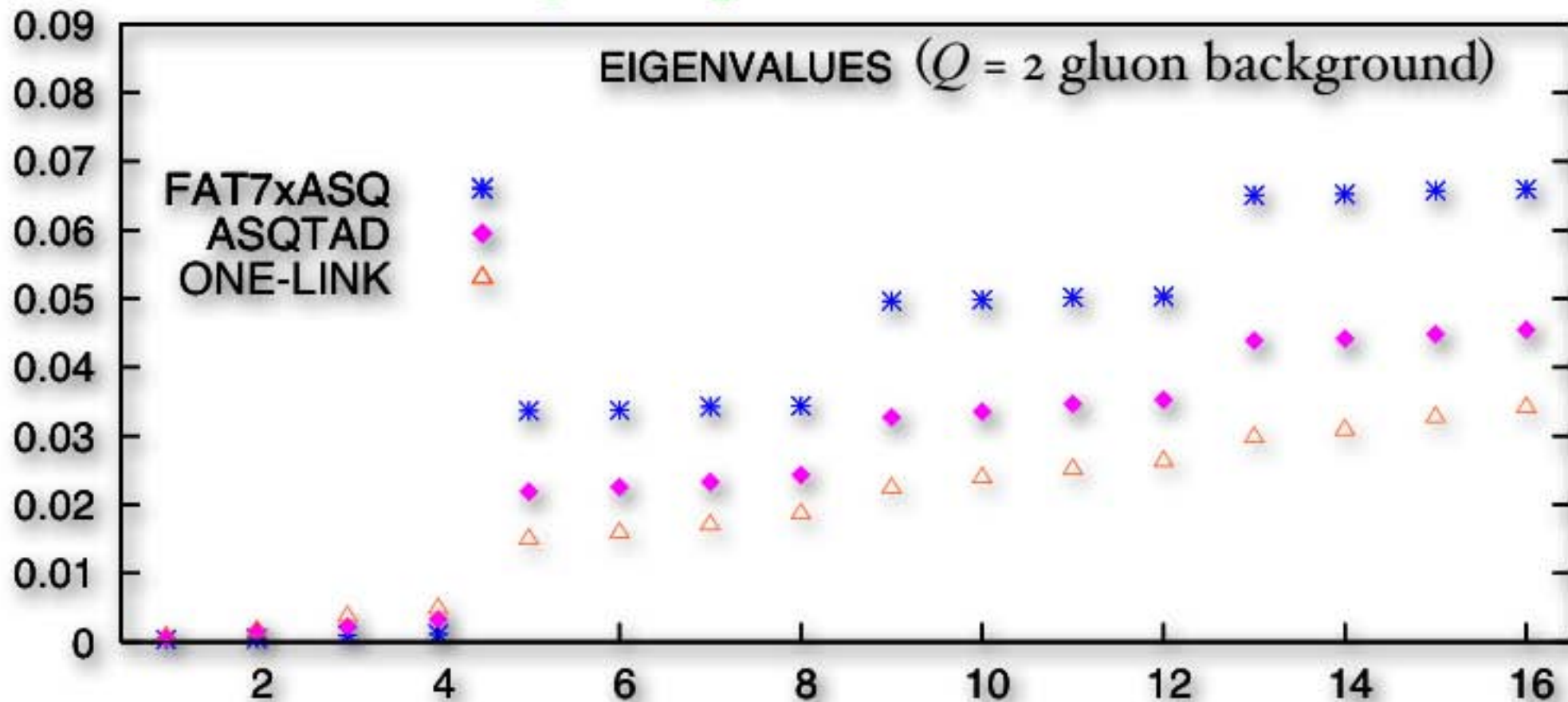
I/4 root  
  
 (so far)



Follana,  
 Davies,  
 Hart,  
 Lepage,  
 Mason,  
*HDT*  
 Lattice  
 2003

∅

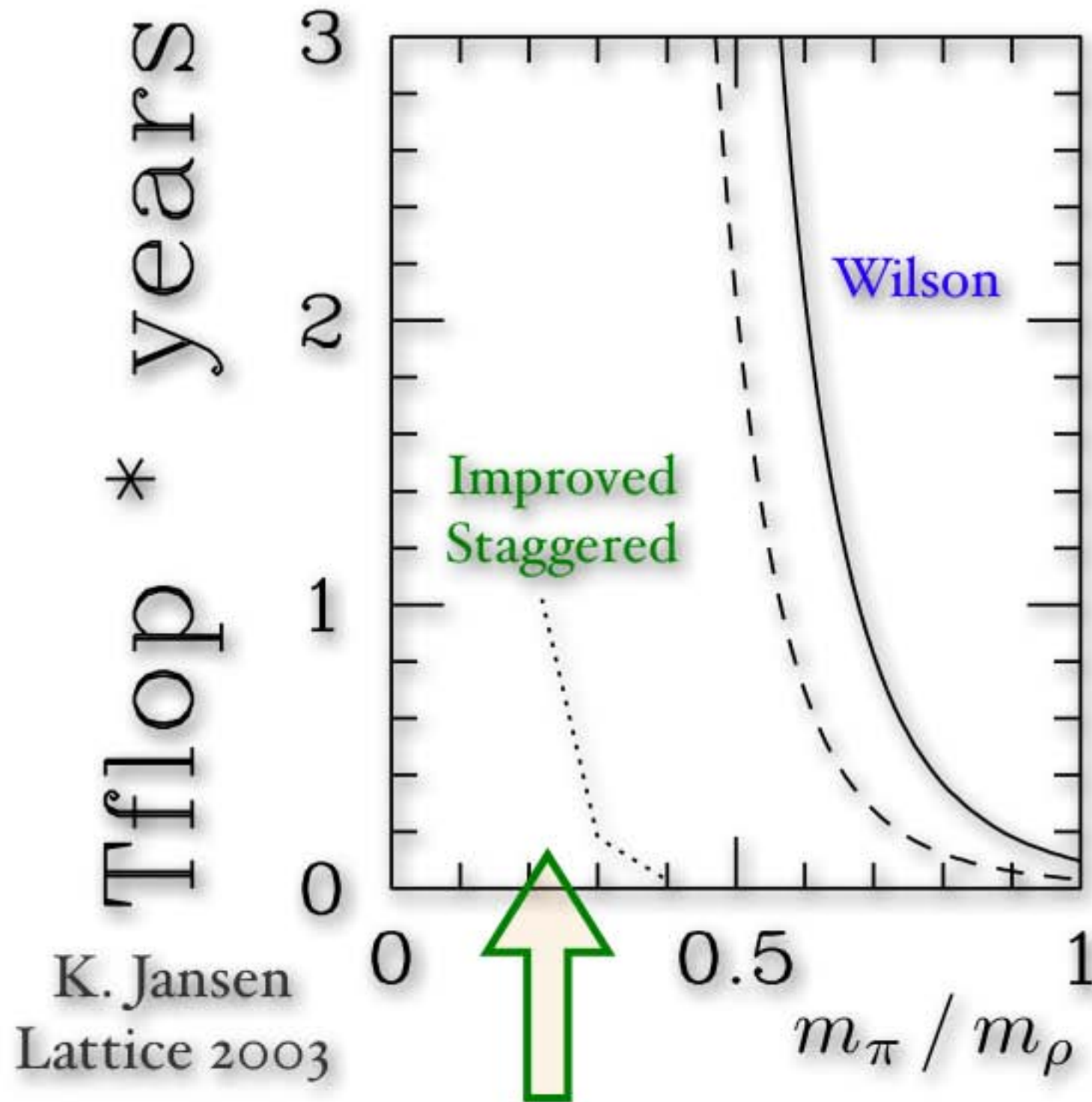
Also:  
 Woloshyn  
 & Wong;  
 Durr,  
 Hoelbling,  
 & Wenger



Follana,  
 Hart,  
 Davies,  
 Mason  
 PRL  
 2004



# The price we pay for speed



Why now?

Discretization effects finally understood using perturbation theory:  
Lepage & Mackenzie

New algorithm  
Wilson quarks:

Appears to radically lower cost:  
M. Luscher,  
Lattice 2005

No surprise: staggered quarks have  $\chi$  symmetry which prevents zero modes  $\Rightarrow$  matrix inversion

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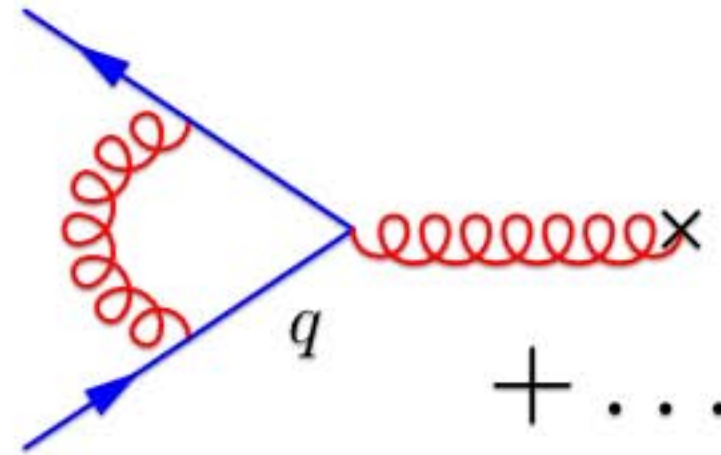
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## (2) Lattice PT key to HPQCD progress

Why PT when lattice is nonperturbative?

Scattering amplitude  $\sim$

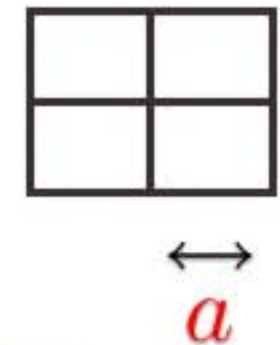


(e.g. off external chromo-B field)

$$\text{Amp}|_{\text{Continuum}} \sim \int_{-\infty}^{\infty} d^4 q \dots$$

How to make up difference?

$$\text{Amp}|_{\text{Lattice}} \sim \int_{-\pi/a}^{\pi/a} d^4 q \dots$$

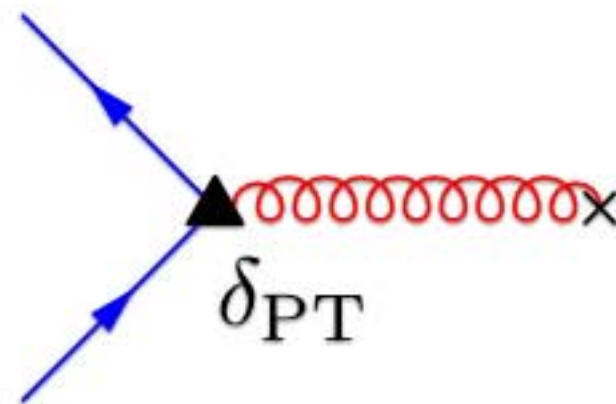


(i) Send lattice spacing  $a \rightarrow 0$

► but lattice simulation cost  $\sim 1/a^6$

## (ii) Add new contact interactions

$$\mathcal{L}_{\text{lattice}} = \mathcal{L}_{\text{Dirac}} + \delta_{\text{PT}} \bar{\psi} a \vec{\sigma} \cdot \vec{B} \psi + \dots$$



Perturbative:  
Fills-in high-energy modes

$$\delta_{\text{PT}} = \text{“Lattice - Continuum”} = c_1 \alpha_s + c_2 \alpha_s^2 + \dots$$

► Few % precision demands ambitious program:

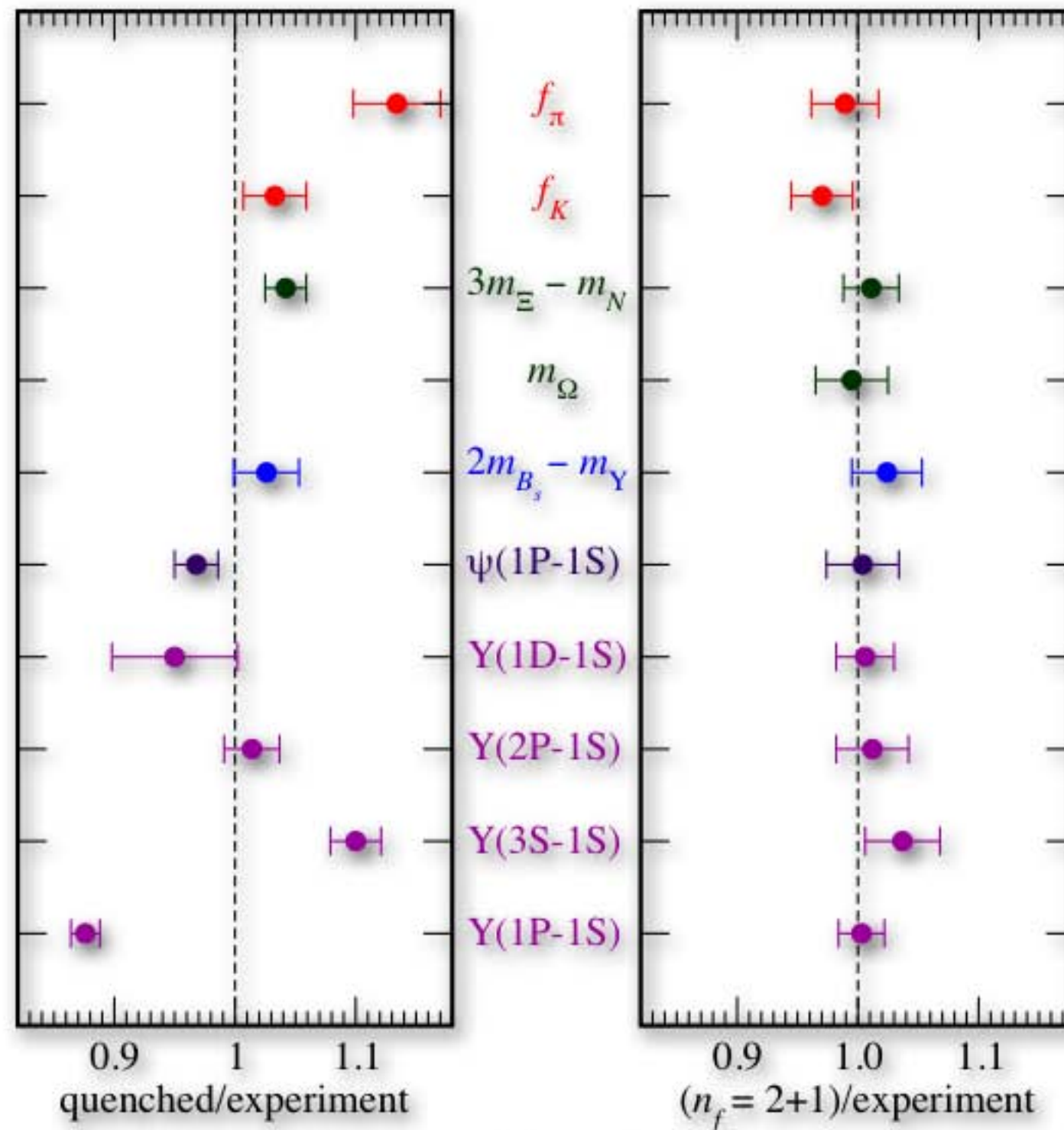
$$\alpha_s(1/a) \approx 0.2-0.3 \approx a \Lambda_{\text{QCD}} \quad (a \approx 0.1 \text{ fm})$$

► PT matching through 2nd-order (NNLO) hard, but  
in progress!

►  $O(a^2)$ -accurate discretizations (HPQCD ✓ but  
not so all other approaches)



# Not much “explicit” PT needed earlier



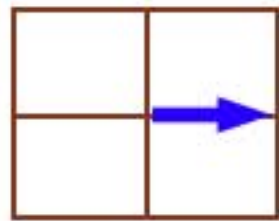
[*Except* for design of “improved” discretizations]

- ▶ “Absolutely” normalized hadron masses and currents
- ▶ However most all other quantities *will* require PT

# Lattice regulator makes PT ☹

$$-4\pi C_F \frac{\alpha_V(q^2)}{q^2}$$

Complication is due to link field



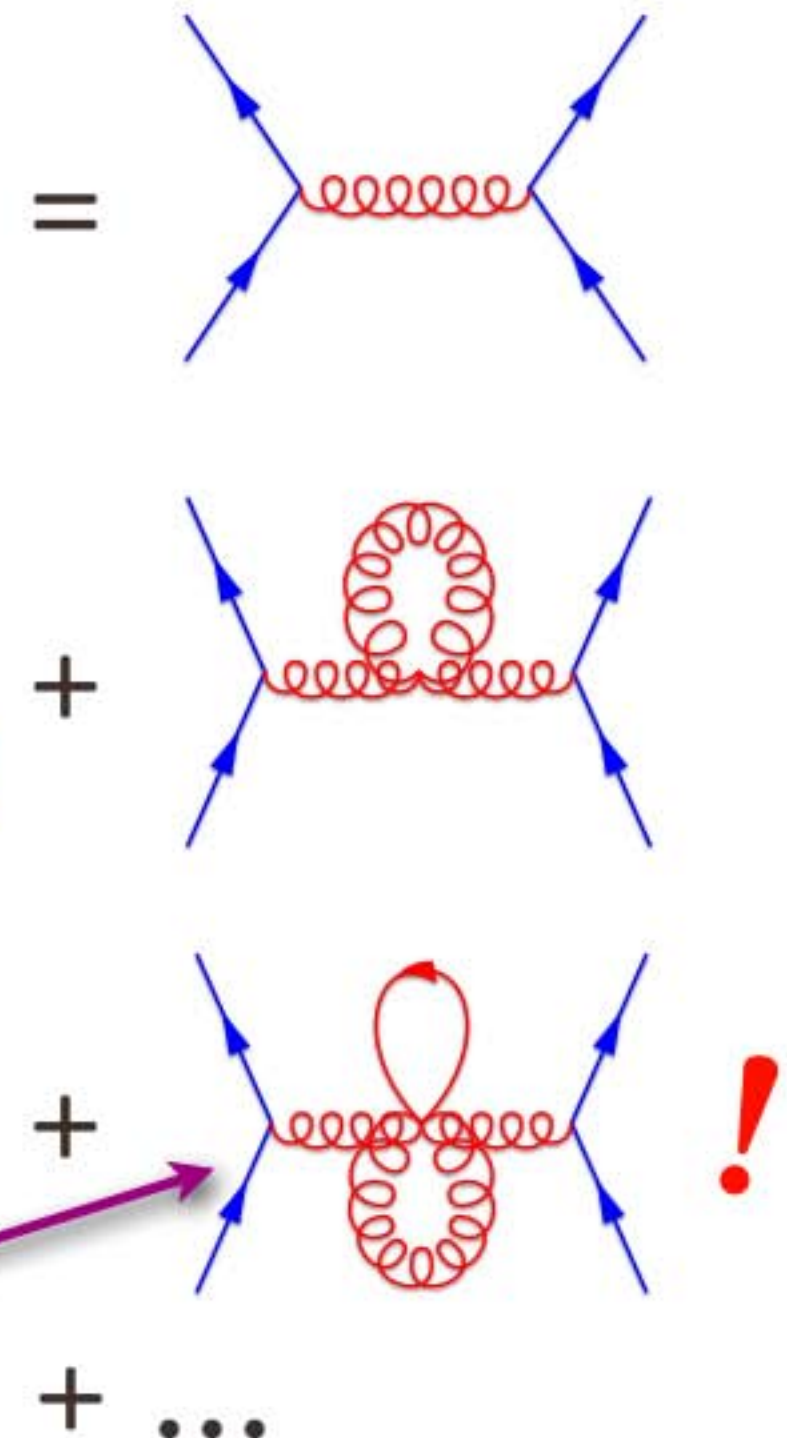
$$U_\mu(x) = e^{igaA_\mu(x)}$$

☺ Lattice PT can be automated!

► M. Lüscher and P. Weisz,

Nucl. Phys. B266, 309 (1986)

10,000's terms





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$\alpha_{\overline{\text{MS}}}(M_Z)$      $m_{u,d,s}$      $D$  &  $B$  decays

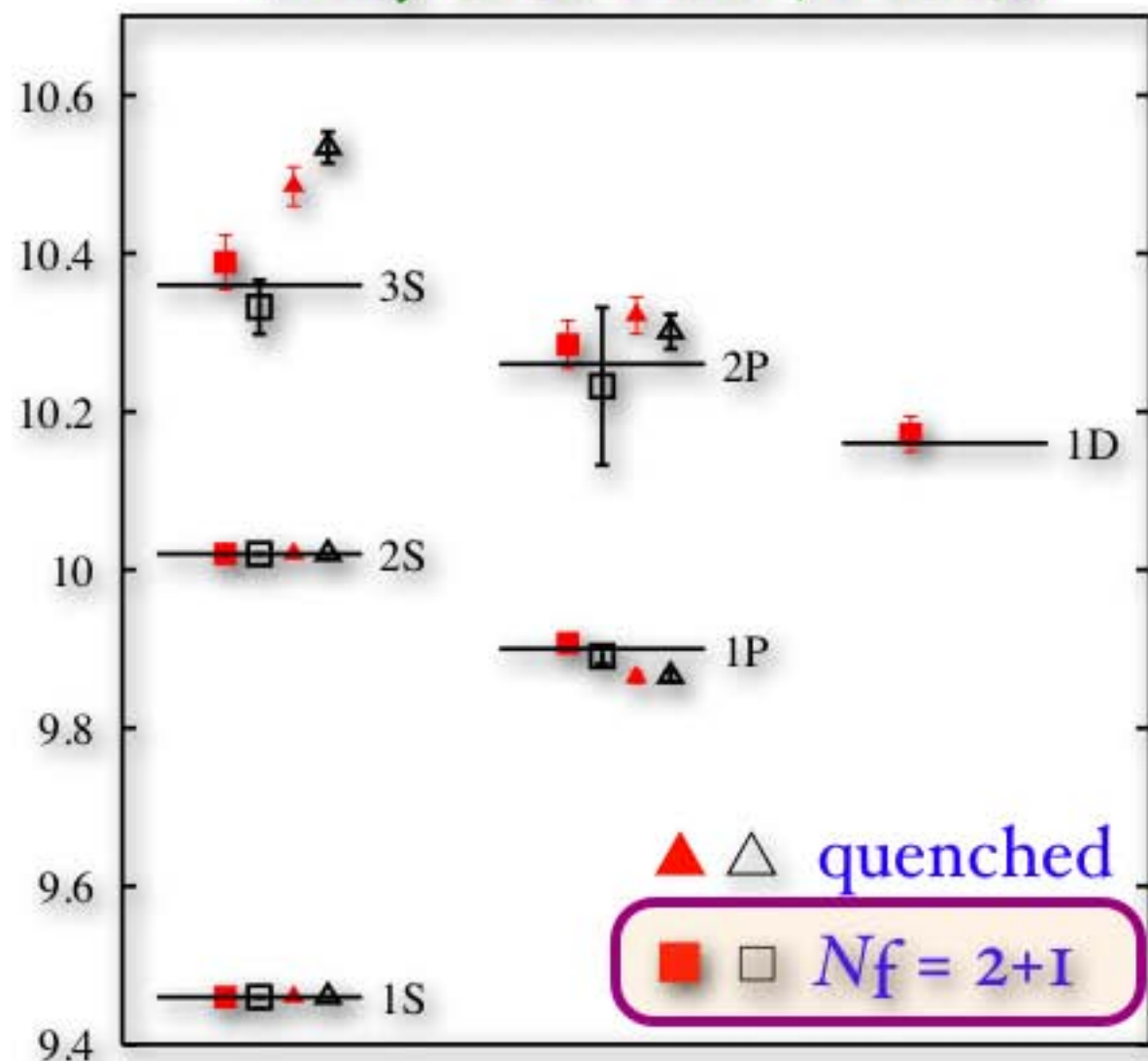
# (4) Other Approaches

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# (3a) Determination of $\alpha_{\overline{\text{MS}}}(M_Z)$

“Novel” use of *both* long- & short distance QCD (NRQCD Collab'n 1997)

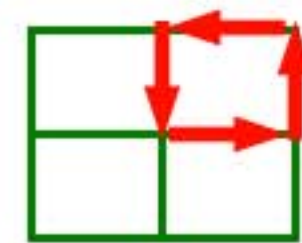
$\Upsilon$  spectra (GeV)  
(Gray et al. PRD 72 2005)



What's new?

(i) NPT input e.g.  $\Upsilon' - \Upsilon \Rightarrow a$

(ii) Measure short-distance quantity



Wilson loop

► Characteristic scale  $q^* \propto 1/a$

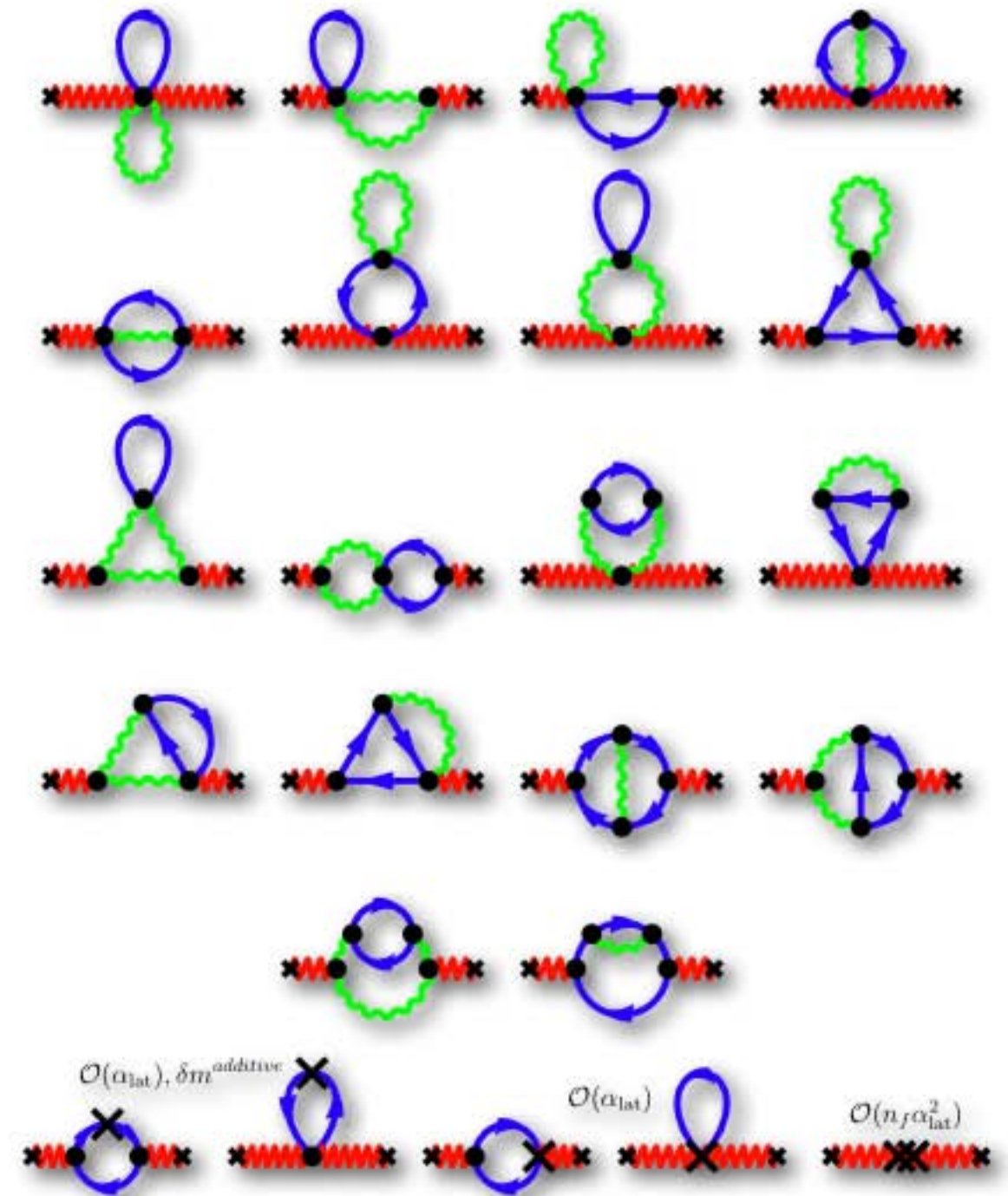
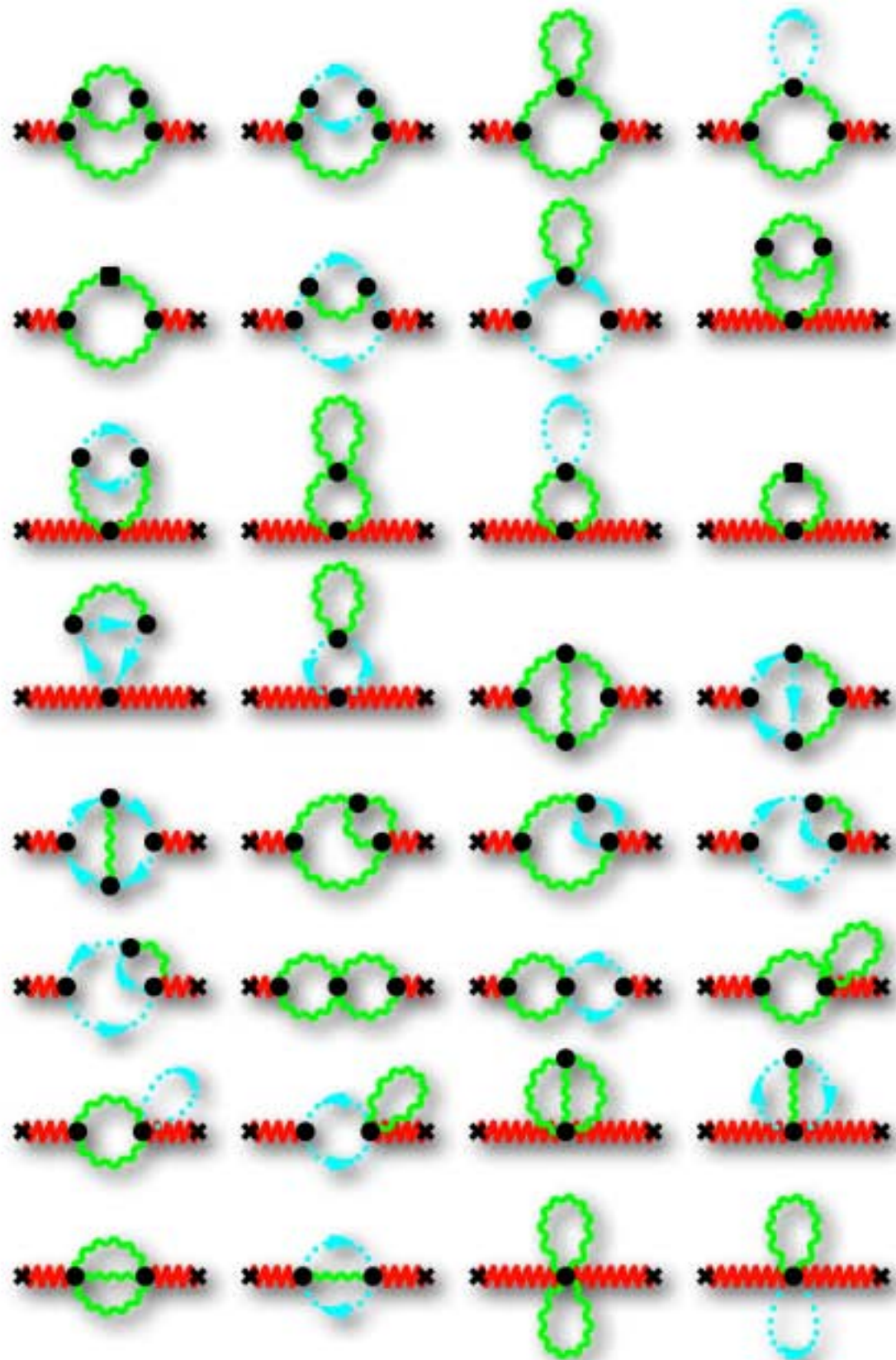
(iii) Use perturbation theory:

$$\langle \mathcal{O} \rangle = c_1 \alpha(q^*) + c_2 \alpha^2(q^*) + c_3 \alpha^3(q^*) + \dots$$

(iv) Evolve  $\alpha(q^*)$  to  $\alpha_{\overline{\text{MS}}}(M_Z)$

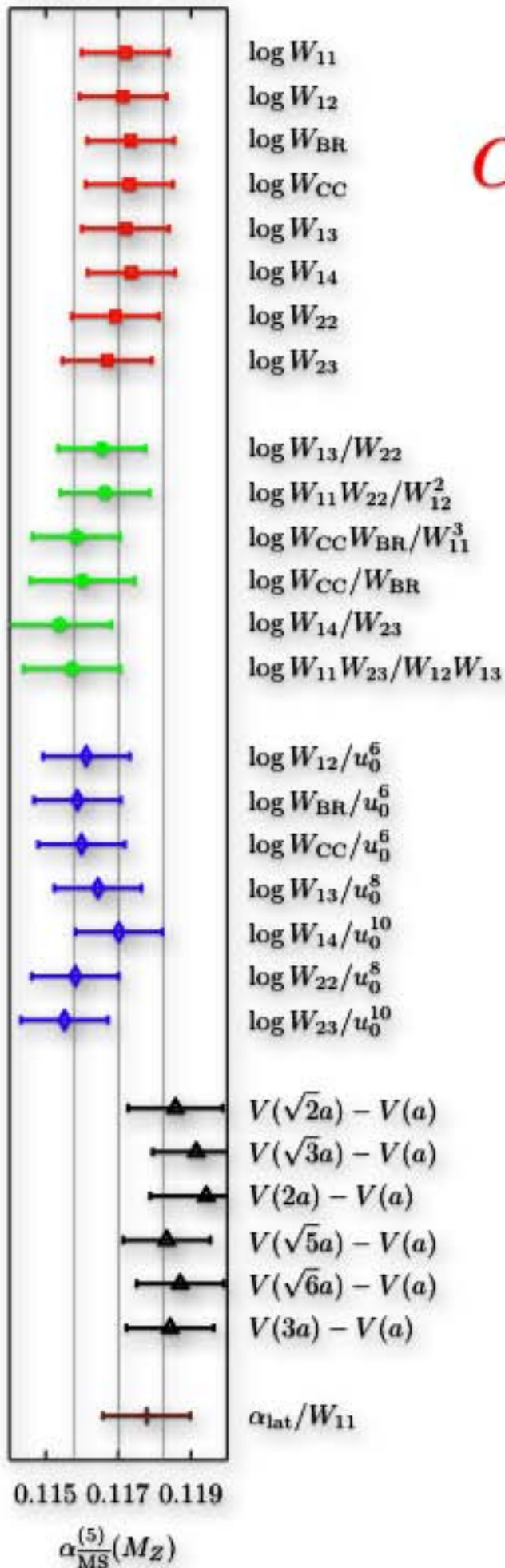


# Some of the NNLO Diagrams





0.115 0.117 0.119



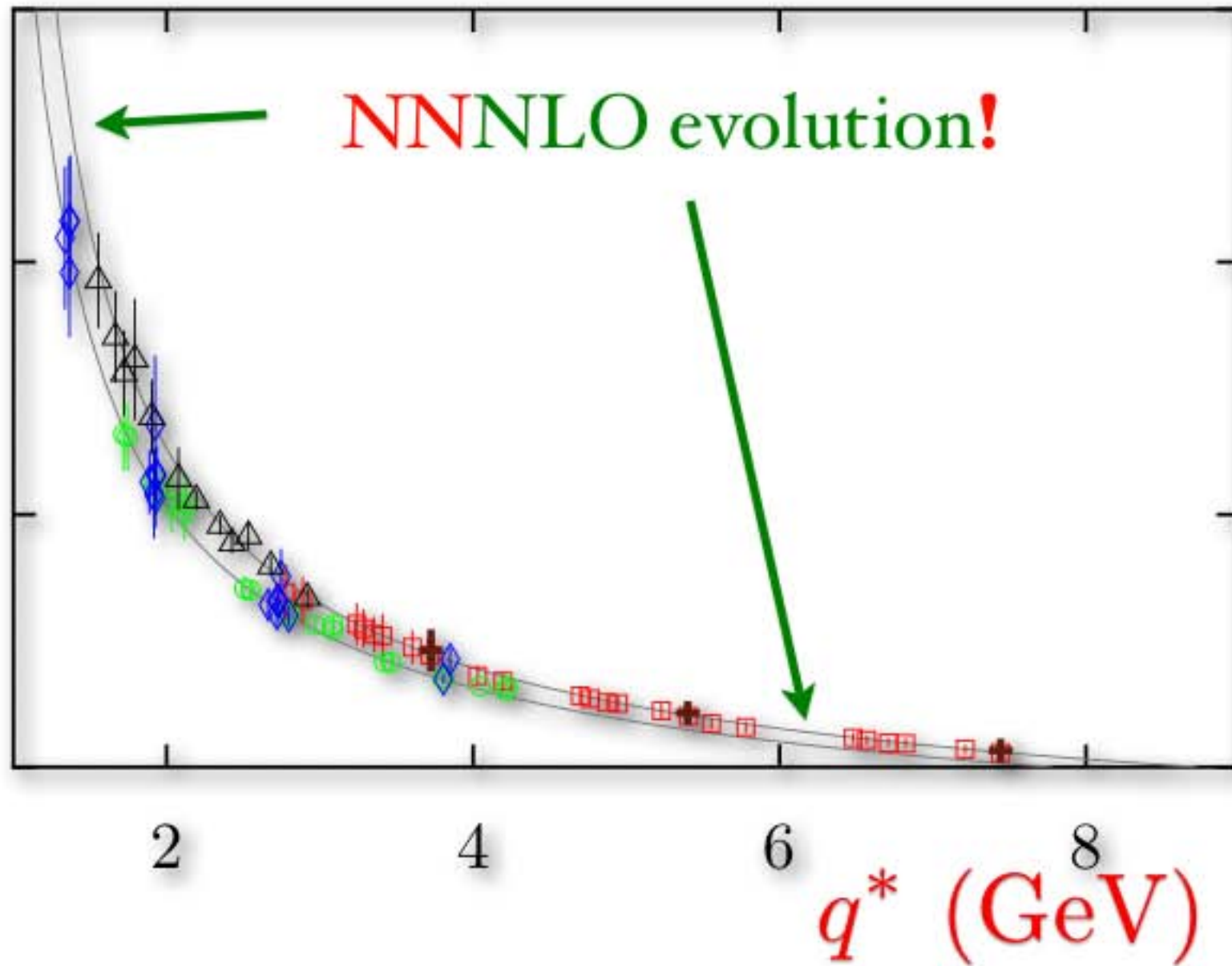
$\alpha_V(q^*)$

0.8

0.6

0.4

0.2

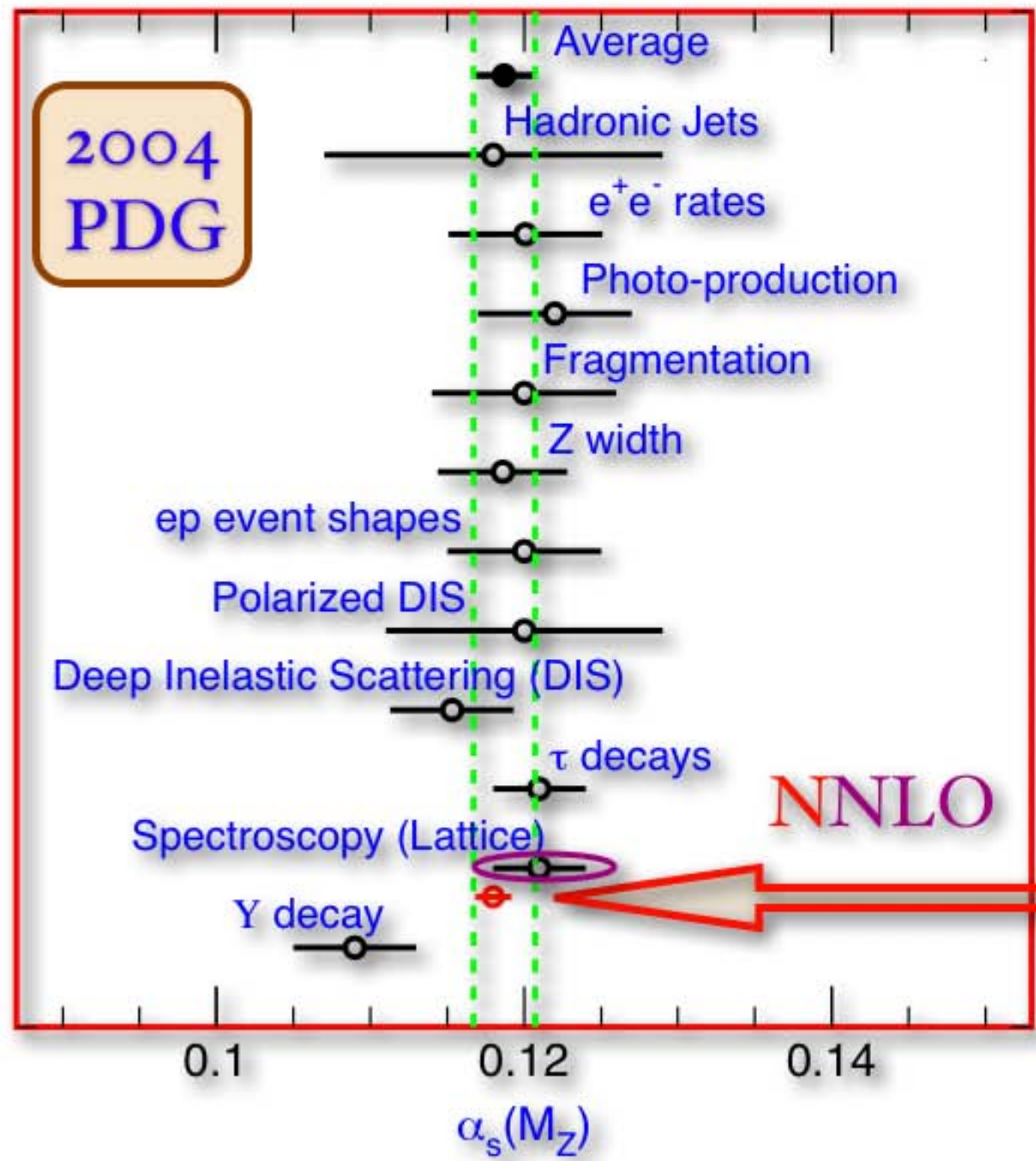


**Error Budget**

	$\log W_{11}$	$\log W_{13}/W_{22}$	$V(\sqrt{2}a) - V(a)$
$a^{-1}$	0.0007	0.0010	0.0010
$c_1 \dots c_3$	0.0001	0.0004	0.0004
$c_n$ for $n \geq 4$	0.0008	0.0005	0.0004
$V \rightarrow \overline{\text{MS}} \rightarrow M_Z$	0.0001	0.0001	0.0001
condensate	0.0002	0.0001	0.0001
$m_u, m_d, m_s$	0.0004	0.0001	0.0001
$m_c, m_b$	0.0002	0.0002	0.0002
simulation errors	0.0000	0.0000	0.0002
total uncertainty	0.0012	0.0012	0.0012



# Result of new NNLO analysis



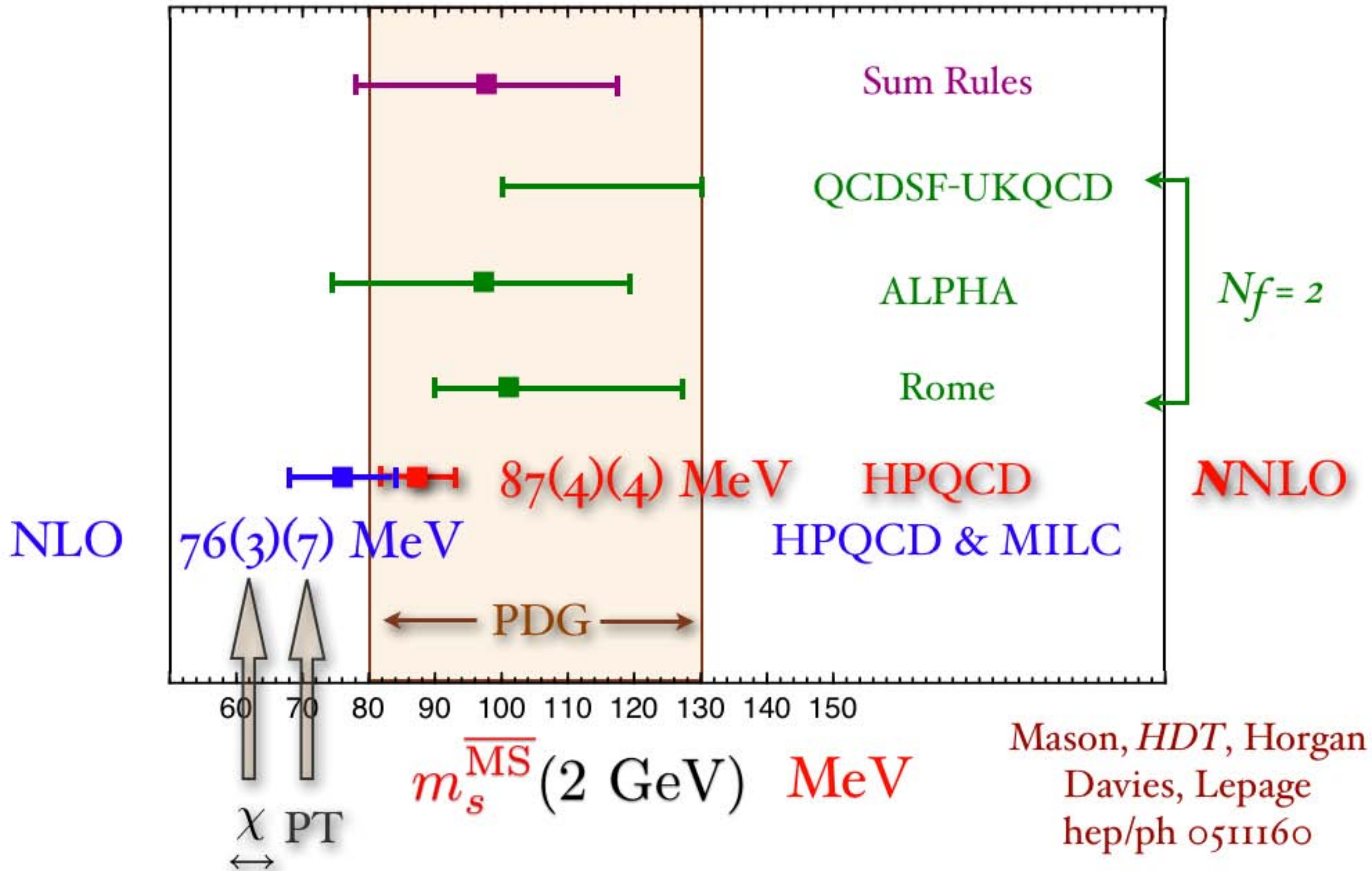
Mason, *HDT*,  
 Davies, Foley,  
 Gray, Lepage,  
 Nobes,  
 Shigemitsu  
 PRL 95,  
 2005

Previous  
 NLO result

$$\alpha_{\overline{MS}}(M_Z) = 0.1170(12)$$

PDG world avg = 0.1187(20)

# (3b) NNLO light quark masses

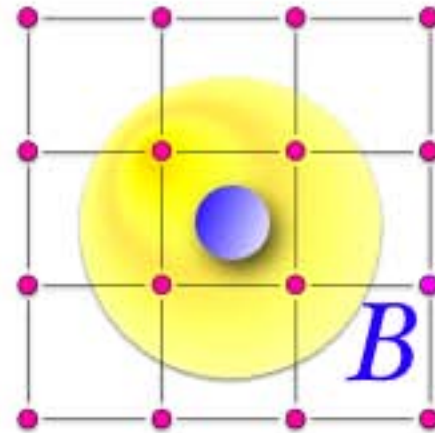




# (3c) D & B meson decays

Heavy quarks on the lattice present another set of challenges

$b$  quarks:  $\frac{1}{m_b} \ll a$



Fortunately:

$$\frac{\Lambda_{\text{QCD}}}{m_b} \ll 1$$

Non-relativistic QCD (NRQCD)

Ex:  
(NLO)

$$|V_{ub}| = 4.22(30)(51) \times 10^{-3}$$

Gulez et al.  
hep-lat/0601021  
(HPQCD:  $N_f=2+1$ )

$c$  quarks:

Fermilab heavy-Q formalism

sums all  
orders  
in  $m_Q a$

Ex:  
(NLO)

$$f_D = 201(3)(17) \text{ MeV}$$

$$= 223(17)(3) \text{ MeV}$$

Aubin et al., PRL 2005  
(FNAL, MILC,  
HPQCD:  $N_f=2+1$ )

CLEOc

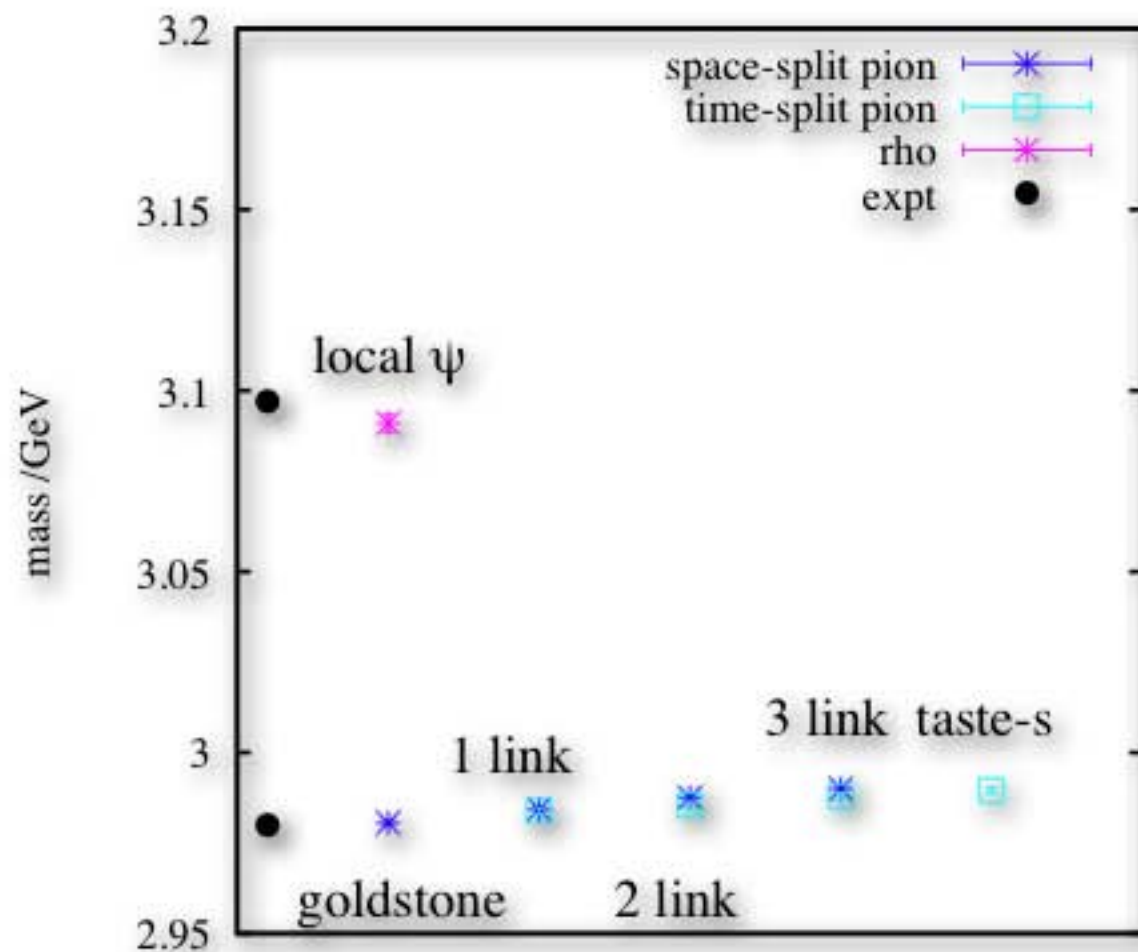


# Relativistic approach to charm

- Use highly-improved staggered action (HISQ) developed for light quarks:  $O(a^4)$ - and  $O(\alpha_s a^2)$ -accurate

- $m_c a \sim 0.4$  on MILC fine lattices:  $\sim 5\%$  accuracy!

Davies, Lepage, Follana, Wong *NNLO*  
*fd: in progress!*



- taste-splittings

(lattice artifacts)  $\sim 10$  MeV

- $2m_{D_s} - m_{\eta_c} =$

0.984(15)(43) GeV (HISQ)

0.956(1) GeV (Expt)

*preliminary*



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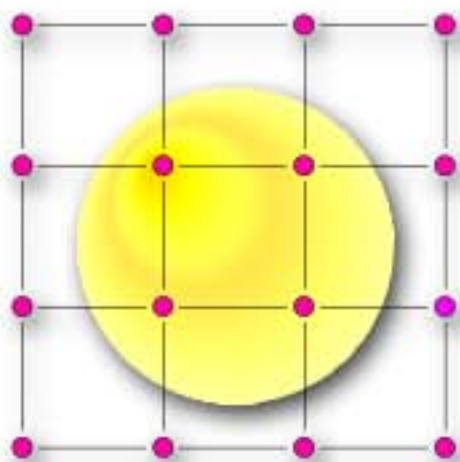
# (4) Other Approaches

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# (4) “Chiral Symmetry vs. Lattice”

Creutz



fermion

Staggered Quarks

HPQCD,  
FNAL, MILC

Accept doublers: work hard to live with them

Wilson Fermions

CPU barrier

Kill doublers, but break chiral symmetry

Luscher,  
Rome Group

- ▶ New simulation algorithm: “*Domain decomposition*”
- ▶ UV & IR mode separation: dramatic performance gain
- ▶  $m_\pi/m_\rho \sim 0.38$  ~ half previous Wilson quark mass !

Still quenched  
&  $O(a)$  →



# Chiral Fermions

Ginsparg-Wilson  
(1982)

Two related schemes  
under development for sometime:

When realistic QCD?

Start with vector-like theory, decouple mirror modes

Domain wall: 5th dimension

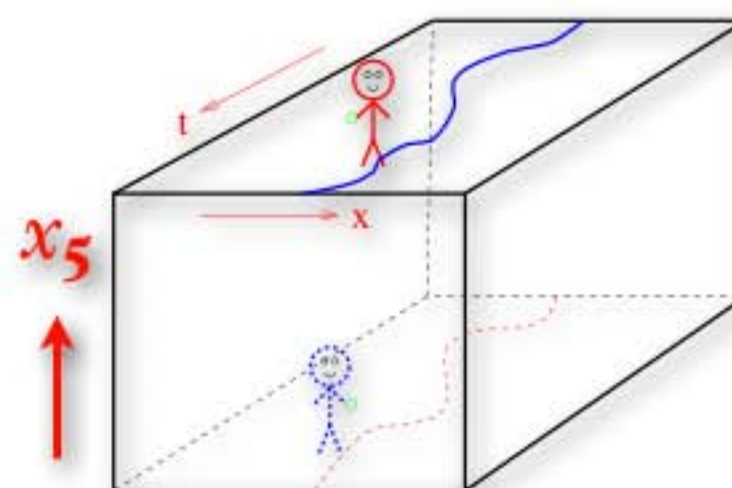
(Kaplan 1992)

Dynamical Simulations:  $N_f = 2+1$

RBC & UKQCD

$m_u / m_s \sim 1/4$

Still early  
days for



Ripped  
from  
Creutz

Overlap fermions: integrate out mirror modes

(Narayanan & Neuberger 1993; Neuberger 1998)

*highly  
nonlinear*

*Exploratory* Dynamical Simulations: Wuppertal; Boulder



# Summary

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- ▶ Unquenched LQCD:  
few-% precision now demonstrated for many quantities
- ▶ HPQCD uses staggered quarks: residual issues
  - ▶ Perturbation theory key to most phenomenology
- ▶ Some important high precision applications done
  - ▶ NNLO  $\alpha_{\overline{\text{MS}}}(M_Z)$  ✓
  - ▶ NNLO  $m_s, m_c, m_b$  ✓
- ▶ Golden opportunity: precision measurements @ CLEO-c
  - ▶ In progress:  $D$ -decays, incl.  $f_D$  to  $\sim 5\%$  this summer
- ▶ Other approaches to realistic LQCD: results *may* be coming?