

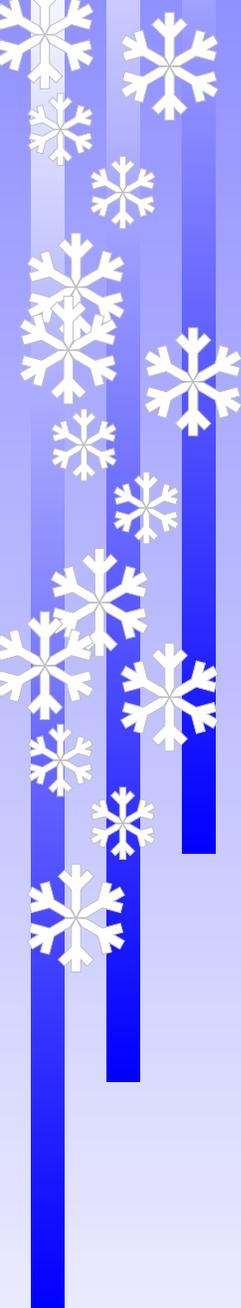
# ELECTROWEAK BARYOGENESIS: *Beyond the Standard Model*

MARIANO QUIROS

ICREA/IFAE

# OUTLINE

- ✓ Introduction and SM results
- ✓ SM extensions with scalars: the MSSM
- ✓ SM extensions with fermions
- ✓ Towards split supersymmetry
- ✓ Conclusions



# INTRODUCTION

Conditions for baryogenesis were stated by Sakharov in 1967<sup>a</sup>

✓ B-violation

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- ✓ B-violation
- ✓ C and CP violation
- ✓ Departure from thermal equilibrium
- ✓ Kuzmin, Rubakov and Shaposhnikov considered in 1985 the possibility of baryogenesis at the electroweak phase transition (EWPT)

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# Introduction

The question that created lot of excitement in the physics community was

**CAN THE SM PRODUCE  
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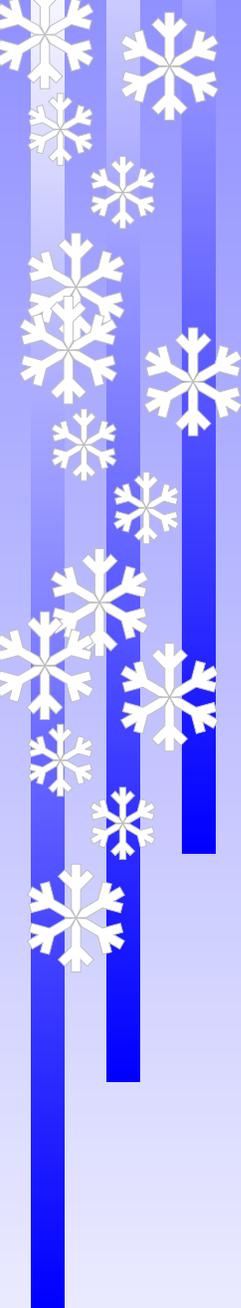
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- ✓ Baryon number is non-perturbatively violated in the SM: sphalerons at finite temperature
- ✓ C and CP violating (CKM) phases are present in the SM
- ✓ The out-of-equilibrium conditions are present in the bubble wall in a **FIRST ORDER PHASE TRANSITION**

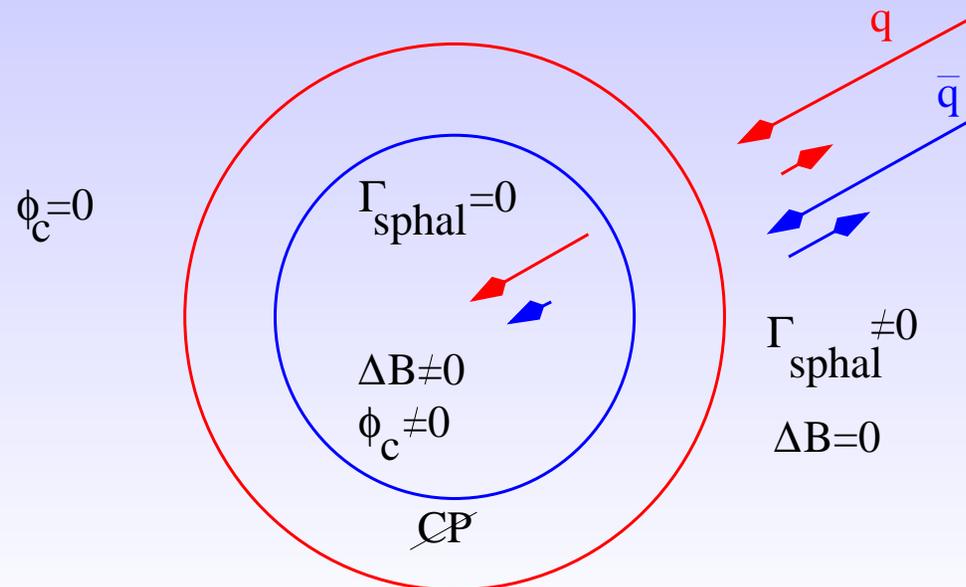


# Electroweak Baryogenesis

A mechanism for the generation of the **BAU** was suggested by **Cohen, Kaplan and Nelson in 1993** using CP violating interactions of fermions with the domain wall of a bubble. The reflection and transmission coefficients of **fermions** and **anti-fermions** scattering off the CP violating wall are different

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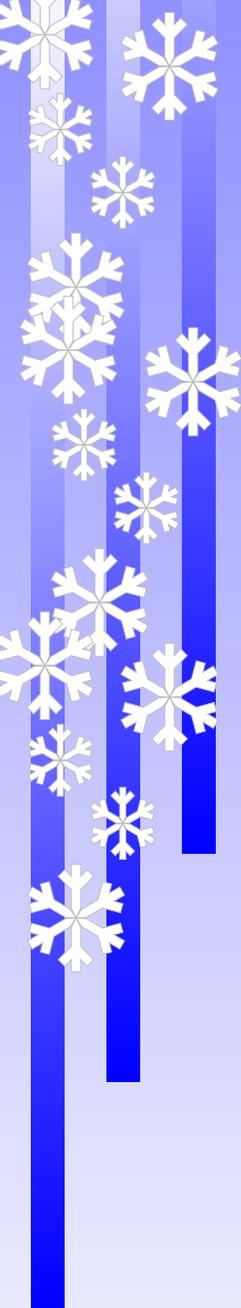


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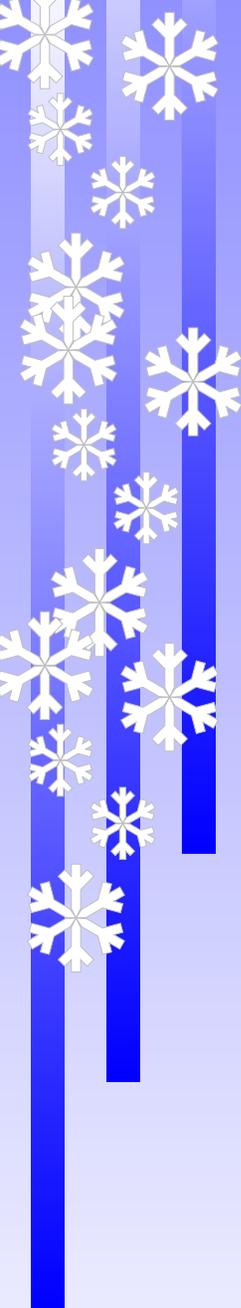
If the phase transition is not strongly enough first order any previously generated BAU is erased by sphalerons in

$$\text{the symmetric phase} \Rightarrow \frac{\phi_c}{T_c} \geq 1$$



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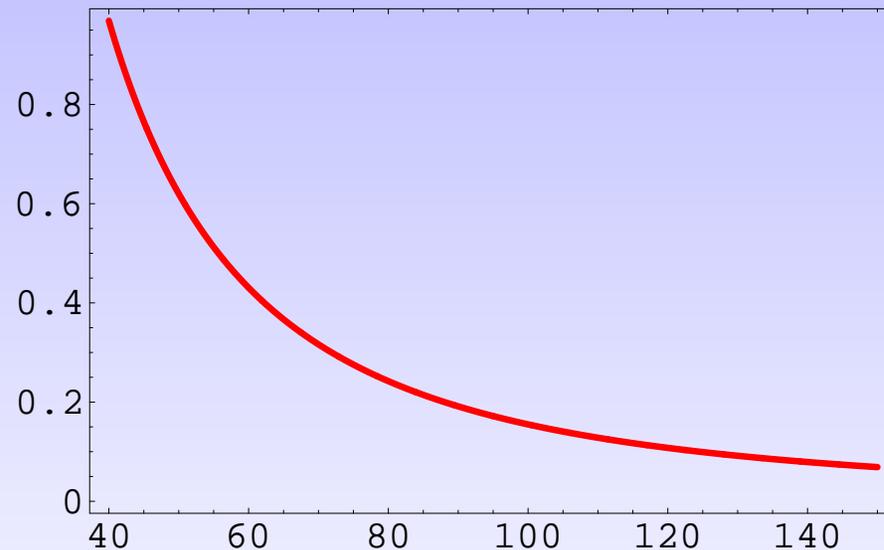
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- The **CP** violation provided by the **CKM** phase is too small to generate the required BAU
- The phase transition is **not strong** enough. Would a BAU be generated it would be **erased** by **weak sphalerons** in the broken phase. In fact the strength of the phase transition strongly depends on the **Higgs mass** and for present experimental limits it is extremely weak. A one-loop (improved by hard thermal loops) result is plotted

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- ✓ Bosons have  $n = 0$  Matsubara modes and thus they contribute to the **cubic terms** in the finite-temperature potential and to create a first order phase transition
- ✓ Bosons appear in **supersymmetric** extensions of the SM: in particular **STOPS**

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  - ★ The SM-like Higgs is light enough

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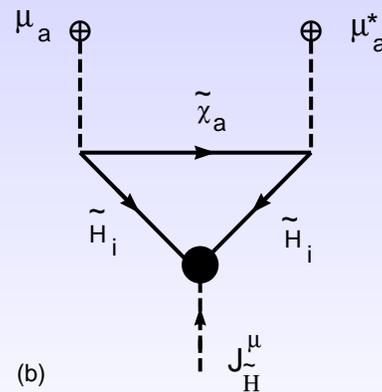
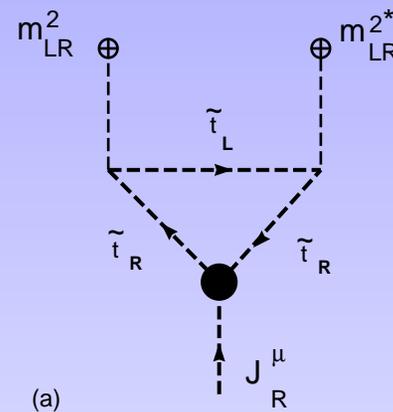
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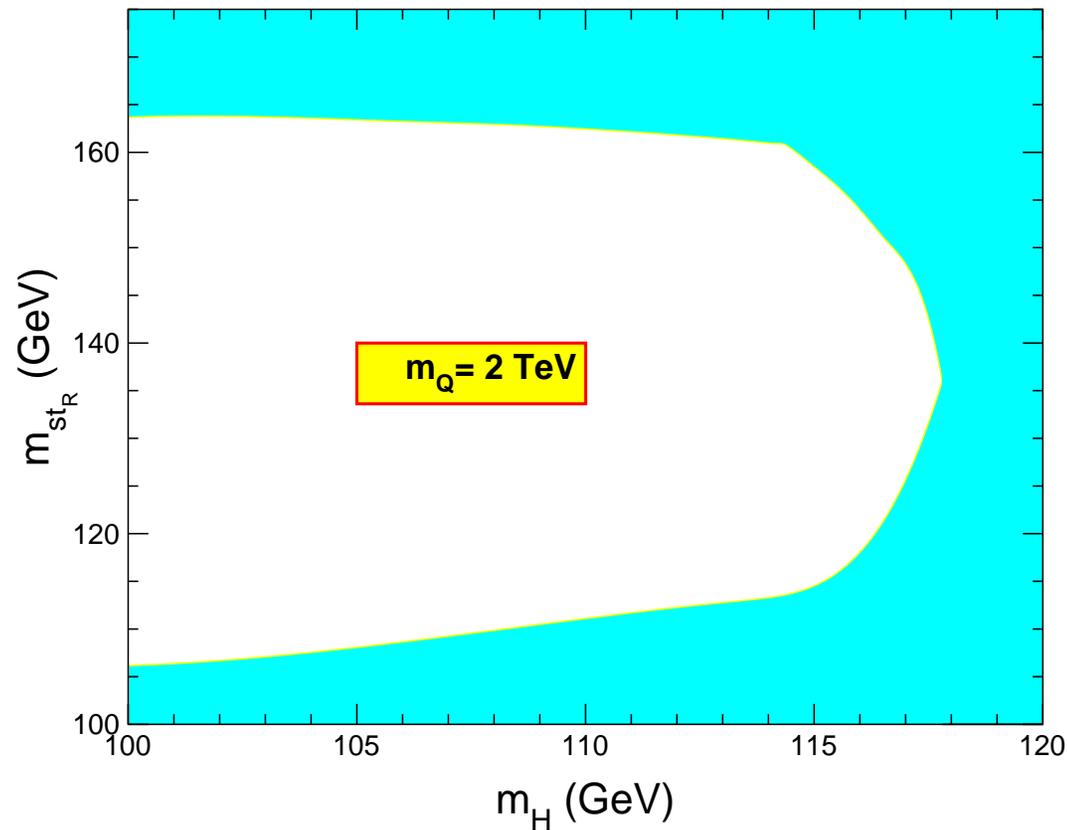
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BAU is barely consistent with WMAP results for  $\mathcal{O}(1)$  phases and light charginos and neutralinos

The lightest neutralino is a candidate to Dark Matter

# MSSM

The baryogenesis window for the MSSM is on the verge of experimental limits



# Conclusion

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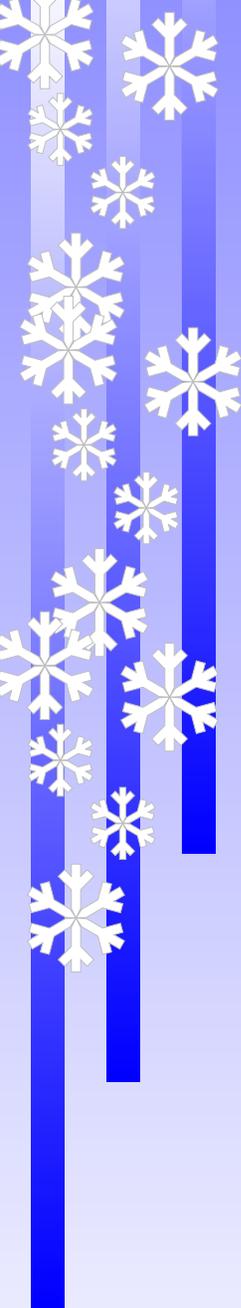
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If the Higgs mass turns out to be (much) larger an economical solution requires that supersymmetry is broken (sfermions masses are) at a much larger scale:

**SPLIT SUPERSYMMETRY**



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Two minimal possibilities arise depending on whether the right-handed stop turns out to be heavy or light:

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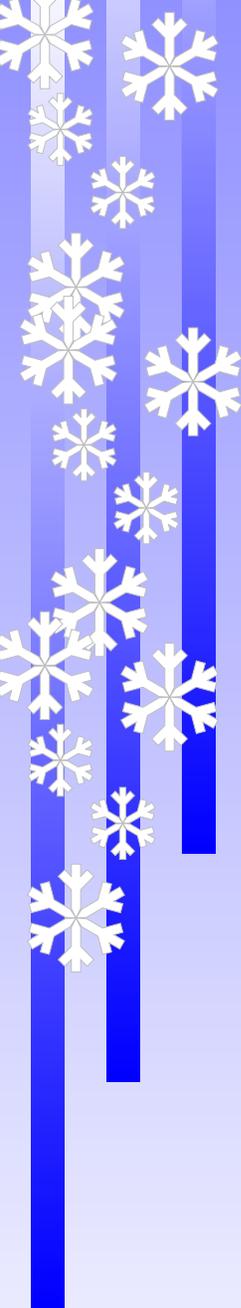
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- If  $\tilde{t}_R$  is heavy ( $\sim m_{susy}$ ) charginos and neutralinos should be strongly coupled and responsible for **both** the strong **phase transition** and **EWBG** [M. Carena, A. Megevand, M.Q. and C.E.M. Wagner, NPB716 (2005) 319]

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- If  $\tilde{t}_R$  is heavy ( $\sim m_{susy}$ ) charginos and neutralinos should be strongly coupled and responsible for **both** the strong **phase transition** and **EWBG** [M. Carena, A. Megevand, M.Q. and C.E.M. Wagner, NPB716 (2005) 319]
- If  $\tilde{t}_R$  is light ( $\sim m_t$ ) charginos and neutralinos can be weakly coupled and only responsible for the **EWBG** [M. Carena, G. Nardini, M.Q. and C.E.M. Wagner, in preparation]: A special version of split supersymmetry.



# SM EXTENSION WITH FERMIONS

We will consider the **SM + Higgsinos** ( $\tilde{H}_{1,2}$ ), **Winos** and **Binos** ( $\tilde{W}^a, \tilde{B}$ ) coupled to the SM Higgs doublet  $H$  with the Lagrangian [M. Carena, A. Megevand, M.Q. and C.E.M. Wagner, NPB716 (2005) 319]

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$$\begin{aligned} \mathcal{L} = & H^\dagger \left( h_2 \sigma_a \tilde{W}^a + h'_2 \tilde{B} \right) \tilde{H}_2 \\ & + H^T \epsilon \left( -h_1 \sigma_a \tilde{W}^a + h'_1 \tilde{B} \right) \tilde{H}_1 + \frac{M_2}{2} \tilde{W}^a \tilde{W}^a \\ & + \frac{M_1}{2} \tilde{B} \tilde{B} + \mu \tilde{H}_2^T \epsilon \tilde{H}_1 + h.c. \end{aligned}$$

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The matching with the MSSM couplings and Higgs fields would be

$$\begin{aligned}h_2 &= g \sin \beta / \sqrt{2}, & h_1 &= g \cos \beta / \sqrt{2} \\h'_2 &= g' \sin \beta / \sqrt{2}, & h'_1 &= g' \cos \beta / \sqrt{2} \\H &= \sin \beta H_2 - \cos \beta \epsilon H_1^*\end{aligned}$$

However we will not match them with the MSSM but instead will consider  $h_i, h'_i$  as **INDEPENDENT PARAMETERS**

# The phase transition

- The phase transition is much stronger than in the SM depending on the values of the Yukawa coupling  $h_1 \simeq h_2 = h$  and the masses  $M_1 = M_2 = M, \mu$

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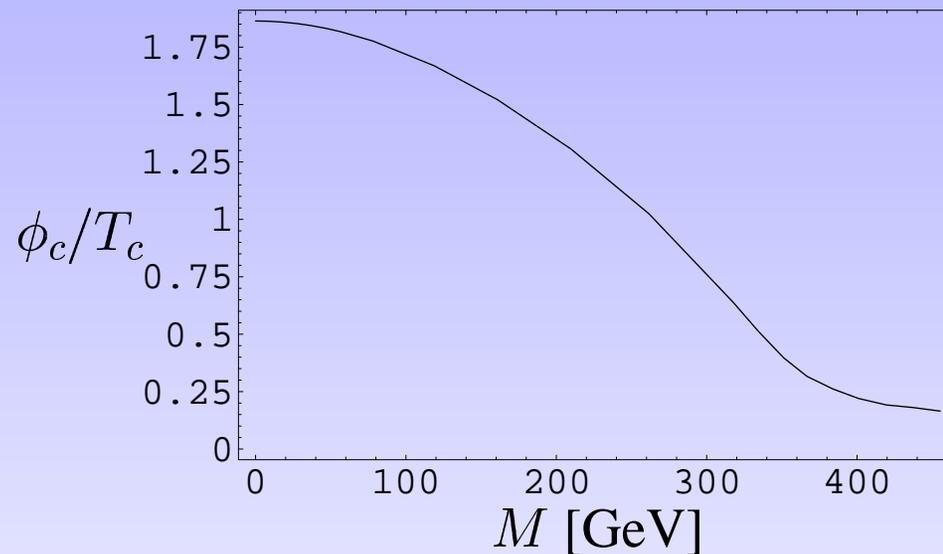


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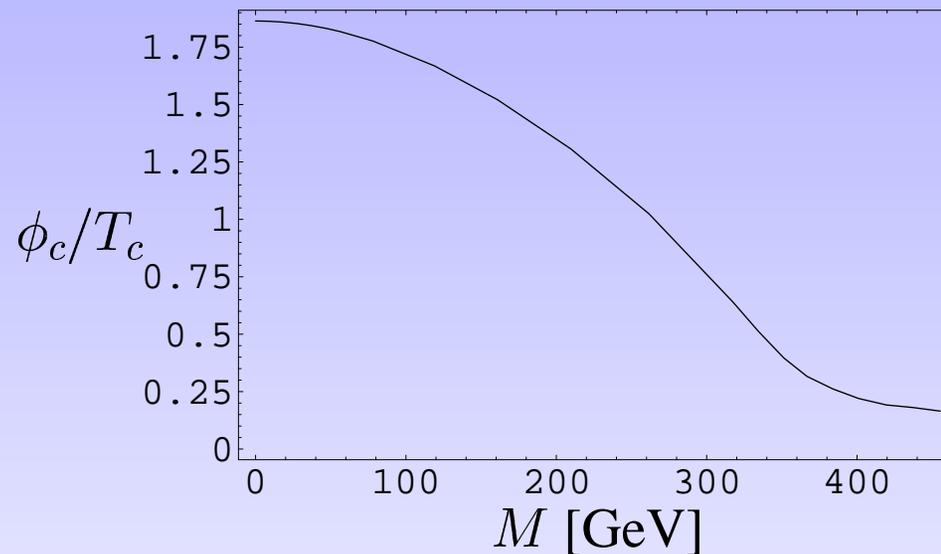


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- In the  $M \rightarrow \infty$  limit the phase transition goes to the Standard Model one

# The Higgs mass

Up to now we have fixed it to a “minimal” value  
 $m_H = 120 \text{ GeV}$

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Our mechanism of strengthening the phase transition, although certainly sensitive to the Higgs mass, permits to go to higher values

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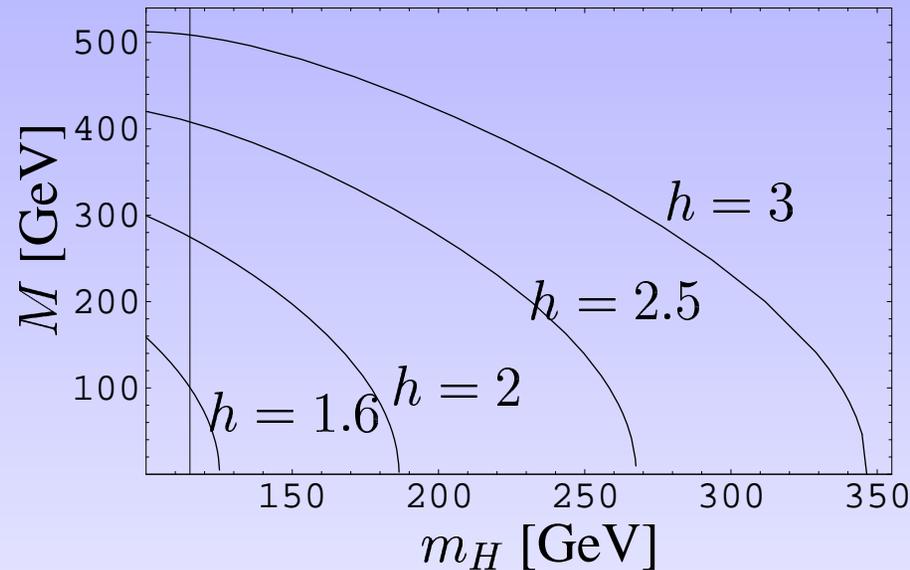
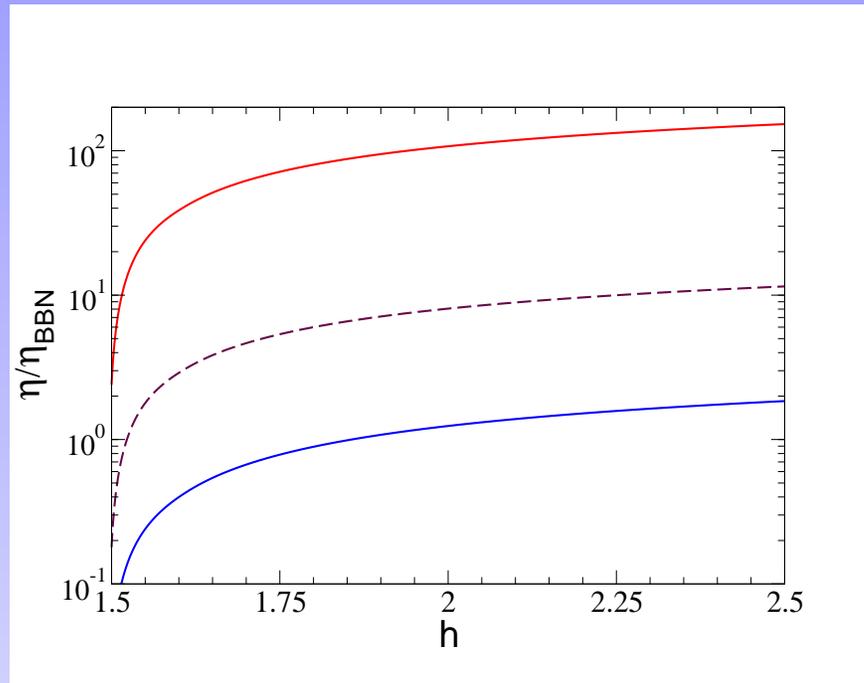


Figure 2: *Contours of  $\phi_c/T_c = 1$  in the  $(M, m_H)$ -plane for  $h = 1.6, 2, 2.5, 3$  and  $M = -\mu$*

# Baryon asymmetry



$\eta/\eta_{\text{BBN}}$  for  $\mu = -M_2 \exp(i\varphi)$ ,  $M_2 = 50 \text{ GeV}$ ,  $\sin \varphi = 1$ , and bubble parameters  $L_\omega = 10/T_c$ ,  $v_\omega = 0.1$ . The lower (upper) solid line corresponds to heavy (light) right-handed stops,  $m_T > 1 \text{ TeV}$  ( $m_T \simeq 100 \text{ GeV}$ ). Dashed line corresponds to right-handed stops with  $m_T \simeq 500 \text{ GeV}$ . WMAP  $\eta_{\text{BBN}} = (8.7 \pm 0.3) \times 10^{-11}$

# Dark Matter

The charginos and two of the neutralinos acquire masses of about  $h_+ v$ . The mass of the lightest neutralino is close to  $|\mu|$  and the lightest neutralino (Dark Matter candidate) is therefore an almost pure Higgsino state

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The annihilation cross section is governed by the coupling of the lightest neutralino to the  $Z$ -gauge boson

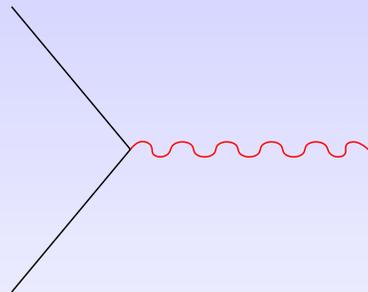


Figure 2: **Annihilation**  $\chi\bar{\chi} \rightarrow Z$

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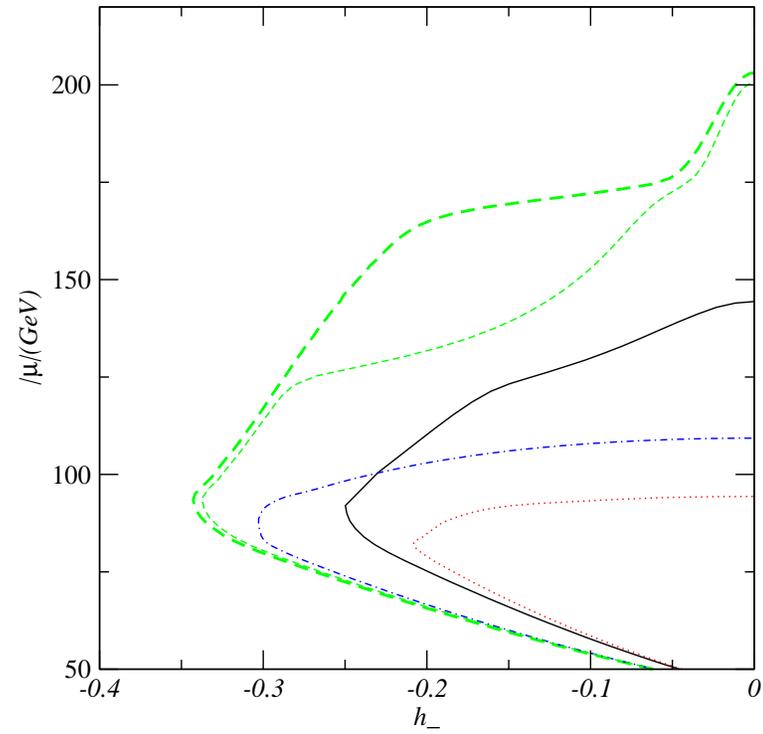
The coupling of a neutralino state to the  $Z$ -gauge boson is proportional to the difference of the square of the components  $N_{\tilde{\chi}\tilde{H}_i}$  of the neutralino into the two weak Higgsino states  $\tilde{H}_i$

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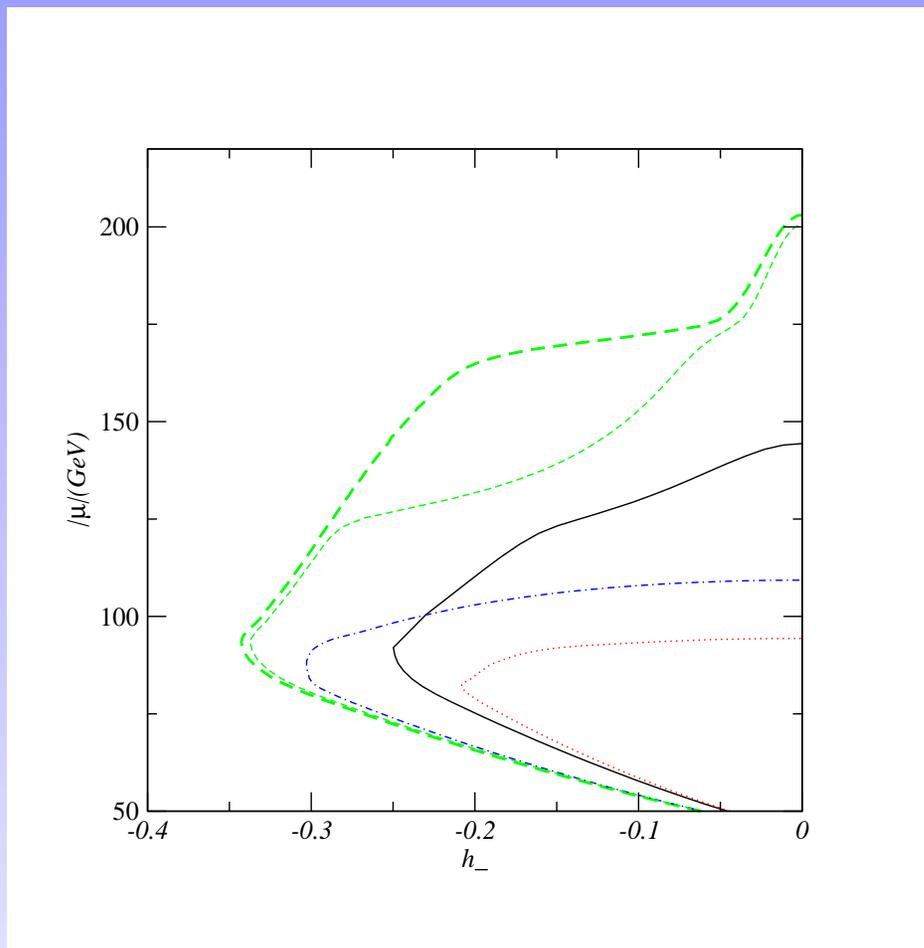
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$$g_{\tilde{\chi}Z} \propto \frac{h_2^2 - h_1^2}{h_2^2 + h_1^2}$$

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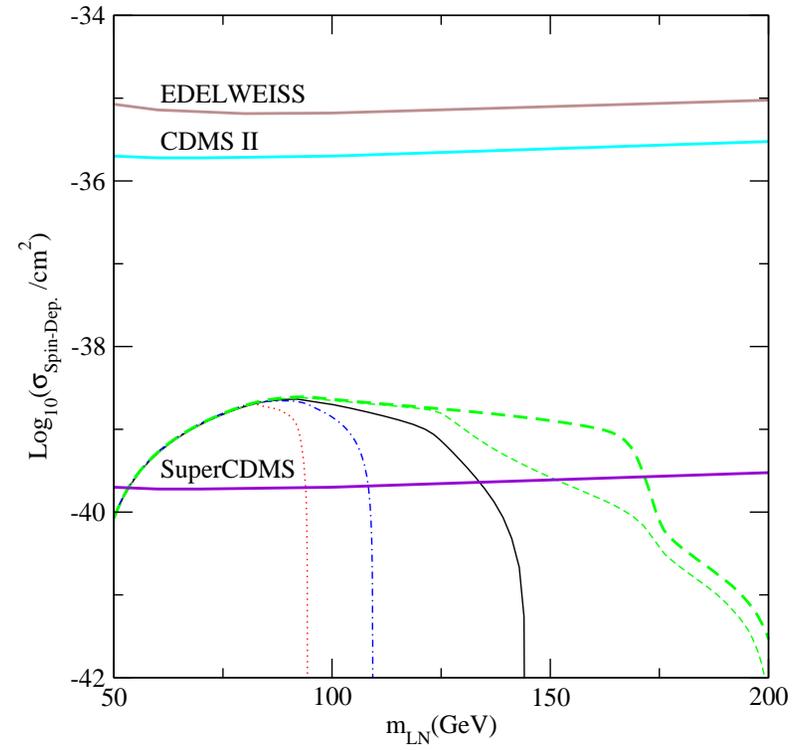
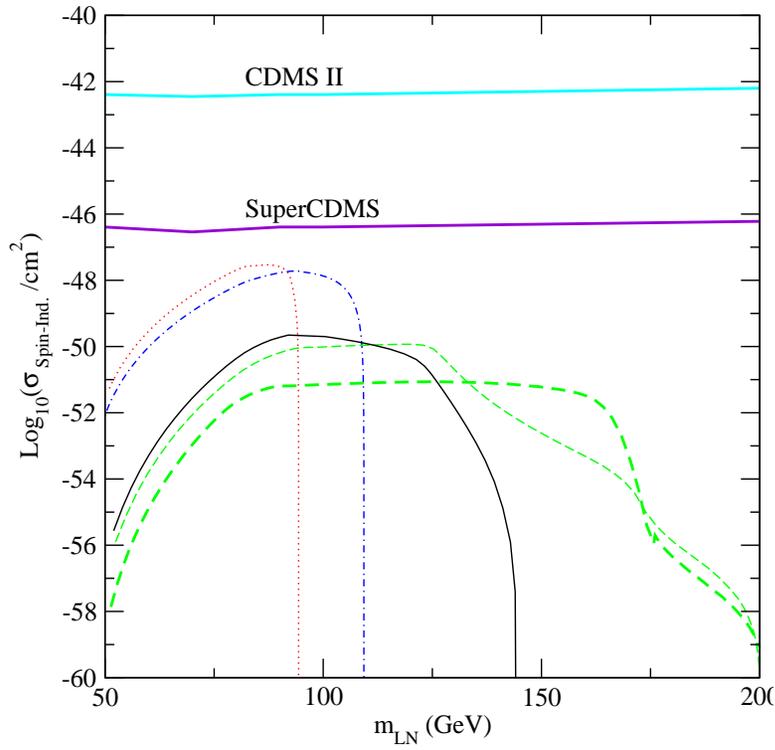


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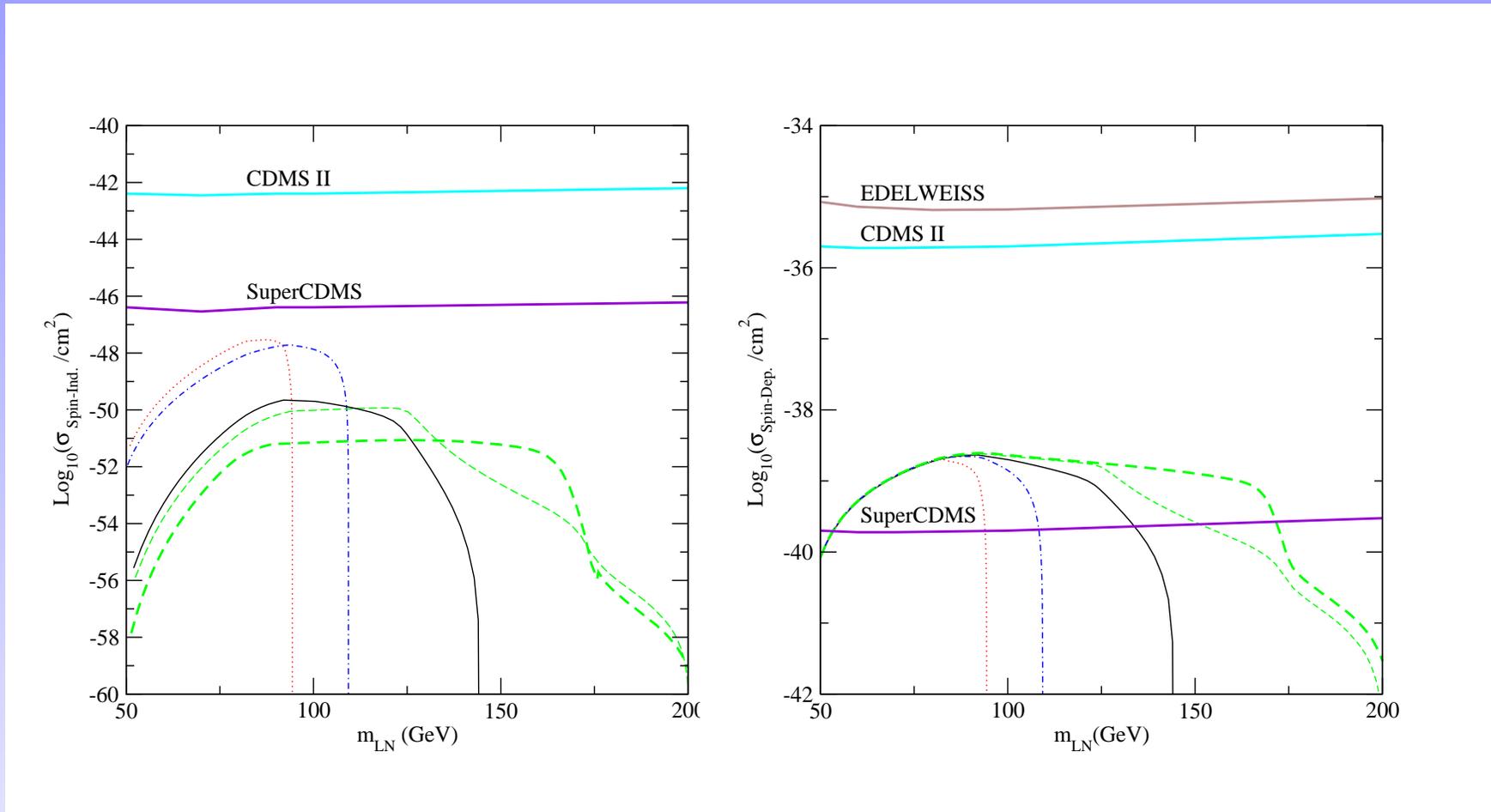


$\Omega_{LN} h^2 = 0.113$  for  $h$ ,  $\varphi$  and  $m_H/\text{GeV} = (2,0,300)$   
[green dashed],  $(2,0,150)$  [green dotted],  $(1.5,0,150)$   
[black],  $(2,\pi/2,150)$  [blue],  $(1.5,\pi/2,150)$  [red]

# Dark Matter detection

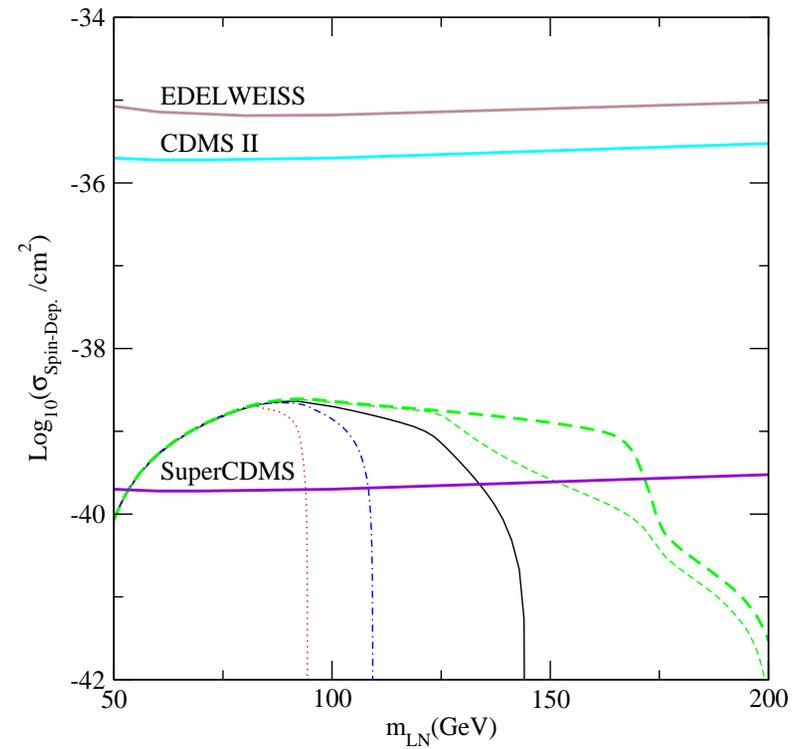
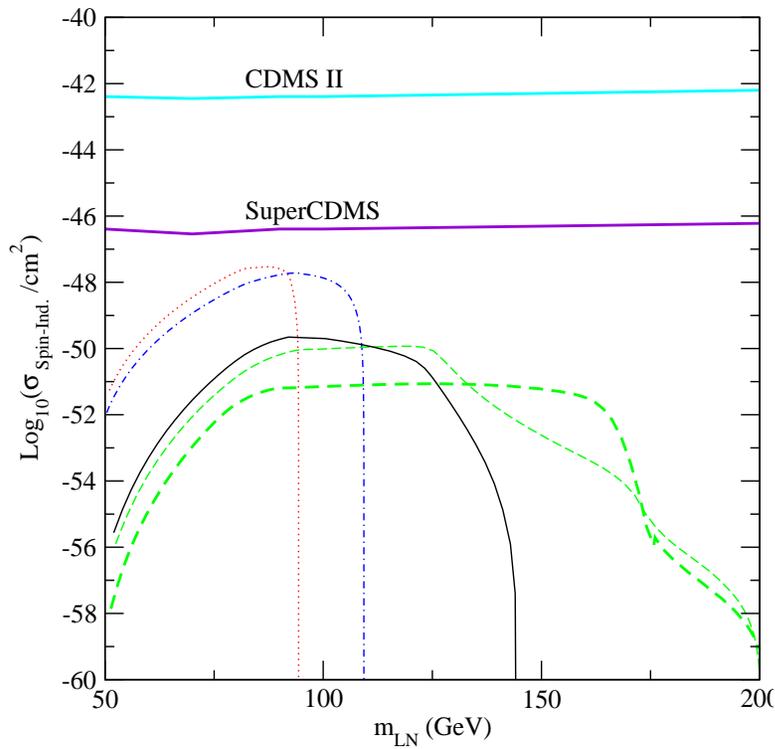


# Dark Matter detection



Proton spin-independent (left panel) and neutron spin-dependent (right panel) cross-sections Vs. current (CDMS II) and projected (SuperCDMS) bounds

# Dark Matter detection



A. Provenza, P. Ullio, M.Q., arXiv:hep-ph/0507325

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- ✓ This theory requires two fine tunings (unlike split supersymmetry)

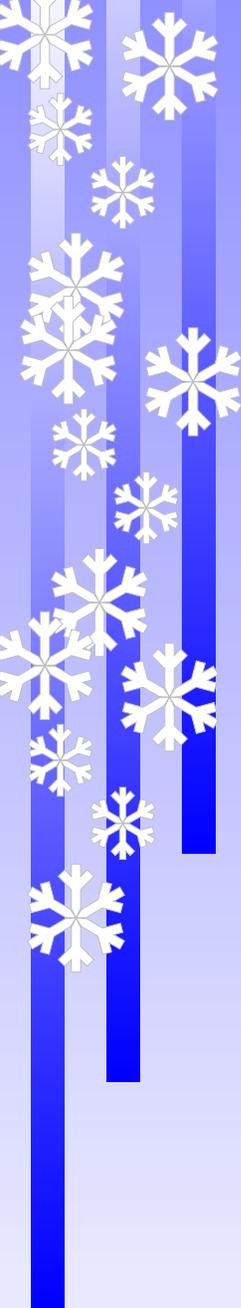
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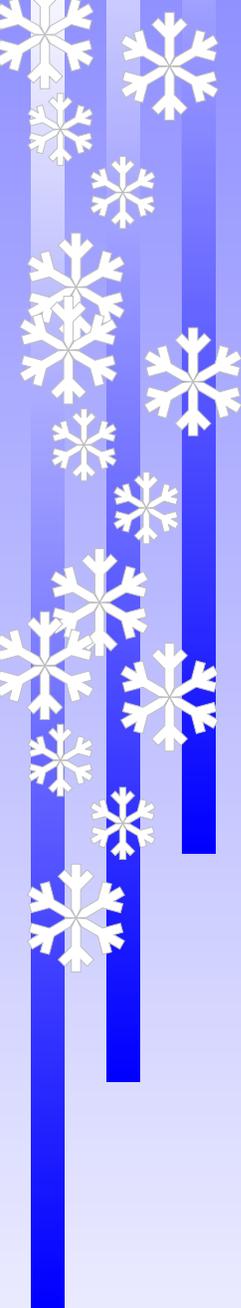
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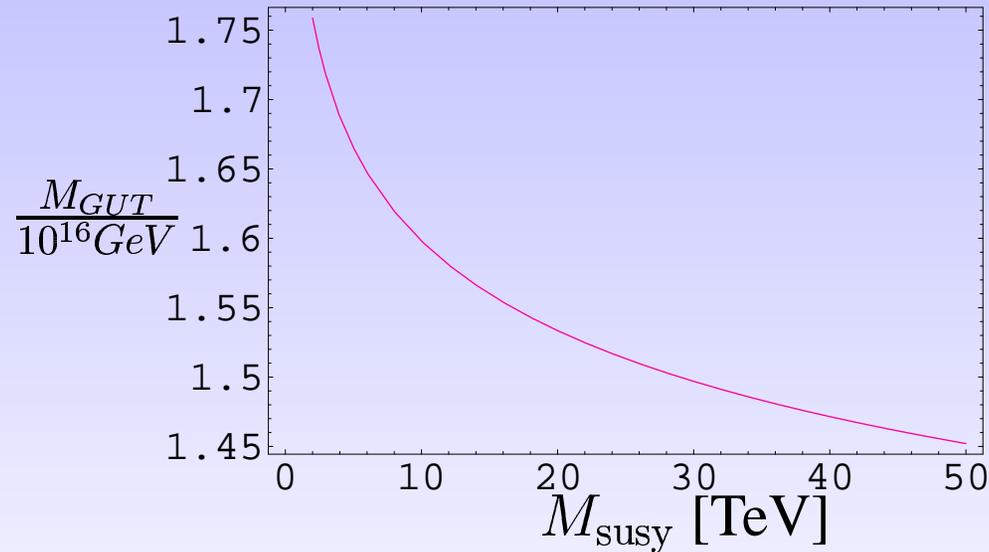
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    - ★ Be consistent with unification
    - ★ Have a Dark Matter candidate
    - ★ Describe EWBG

# Unification

Unification is consistent with scale of supersymmetry breaking around 10 – 100 TeV. [Including two-loop corrections amounts to  $\Delta_{2-loop}\alpha_3(M_Z) \sim 0.01$ .] The one-loop prediction is

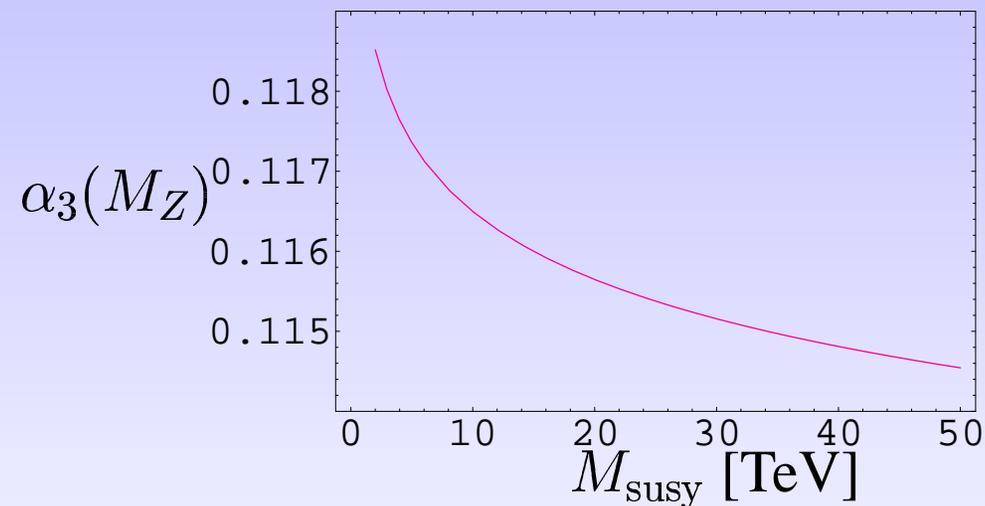
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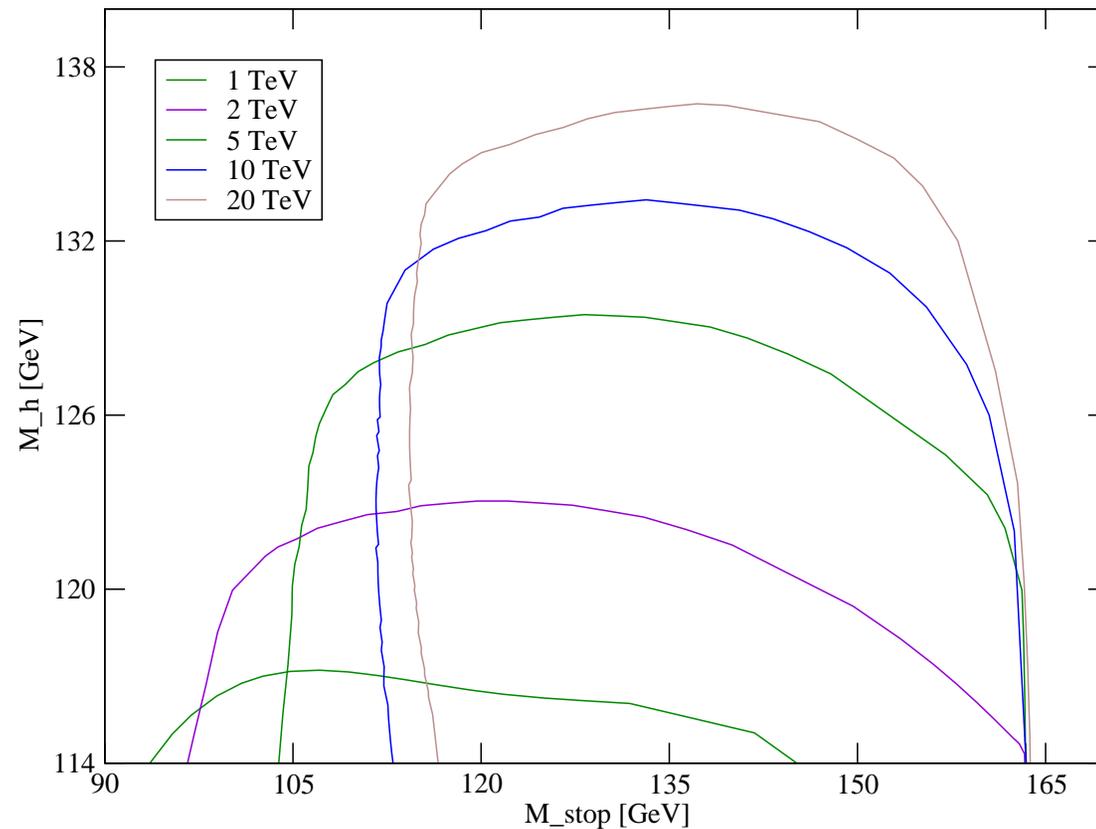


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# Phase transition



Level contours of  $\frac{\phi_c}{T_c} \simeq 1$  for different values of  $M_{\text{susy}}$

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EWBG can be tested at present and future accelerators that are useful to probe models

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- ✓ If the Higgs mass turns out to be below  $\sim 200$  GeV then probably some sort of split supersymmetry can do the job

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- ✓ Present LEP data already exclude the SM and thus require New Physics
- ✓ If the Higgs mass was on the verge of experimental detection at LEP and the right-handed stop turns out to be light ( $\sim m_t$ ) then the MSSM could be responsible for the baryon asymmetry
- ✓ If the Higgs mass turns out to be below  $\sim 200$  GeV then probably some sort of split supersymmetry can do the job
- ✓ If the Higgs mass is much heavier ( $\gg 300$  GeV) then we should (most probably) abandon the idea of SUPERSYMMETRY